Precision Higgs Physics

the Next Step

short version

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Ben, Daniel and Nigel

Outline: there are three vital steps or stages one must climb

- Theoretical precision: Missing Higher Orders (MHO)

- On – Off Shell: the Dalitz sector

- BSM: SM $\oplus$ $d = 6$ operators

Chiara Mariotti, Reisaburo Tanaka
Ansgar Denner and André David

Assertion

Precision Physics : *restricting our attention to the relative merits of realism and instrumentalism.*

Do we have a way of knowing whether "unobservable" theoretical entities really exist, or that their meaning is defined solely through measurable quantities?

Leplin (1984), **Sokal** (2001)

✳ **If you believe that $\lambda(M_{\text{plank}}) = 0$, please skip the rest of the talk** ✳

Here we go

Prolegomena

From my Logbook:
now we must move on to the next step

melting BSM-*physics with high-precision* SM-*technology* The question has been repeated many times

But:

- **Answers** converging around *Not yet*

WELL, SEVERAL YEARS AGO WE AVOIDED THAT FATE, MAY BE THE HISTORY WILL REPEAT ITSELF?

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But:

- **Answers** converging around *Not yet*
- *Meanwhile, it came dangerously close to realizing a nightmare, of Physics done by sub-sets of diagrams instead of cuts*.

WELL, SEVERAL YEARS AGO WE AVOIDED THAT FATE, MAY BE THE HISTORY WILL REPEAT ITSELF?

Missing Higher Orders

What is THU?

The traditional way for estimating THEORETICAL UNCERTAINTIES associated to collider physics is based on the notion of **QCD** scale variation

We introduce the concept of

- **MHO**(**MHOU**), missing higher order (uncertainty), which has to do with the TRUNCATION ERROR IN THE PERTURBATIVE EXPANSION;

In the past 30 years the commonly accepted way for estimating MHOU has been based on scale variations.

Consider an observable $\sigma(Q,\mu)$ where

- Q is the typical scale of the process and
- $\mu \equiv \{\mu_R, \mu_F\}$ are the renormalization and factorization scales. The *conventional strategy* defines

$$
\begin{aligned}
\sigma_\xi^- &= \min\left\{\sigma\left(Q, \frac{\mu}{\xi}\right), \sigma(Q, \xi\mu)\right\}, \\
\sigma_\xi^+ &= \max\left\{\sigma\left(Q, \frac{\mu}{\xi}\right), \sigma(Q, \xi\mu)\right\},
\end{aligned}
$$

- selects a value for ξ (typically $\xi = 2$) and predicts $\sigma^- \leqq \sigma \leqq \sigma^+$

There is an open and *debatable* question on how to assign a probability distribution function (**pdf) to the MHOU**

- the *generally accepted* one is based on a Gaussian (or log-normal) distribution centered at $\sigma(Q,Q)$. What to use for the standard deviation, remains an *open problem*.
- Alternatively, it can be assumed that the pdf is a FLAT-BOX

Recently, there has been a proposal by **Cacciari and Houdeau**, based on a flat (**uninformative**) *Bayesian prior* for the MHOU.

More generally, dependence on scales is only part of the problem: indeed, the MHO problem is based on the following fact: given an observable $\mathscr{O}$, related to a perturbative series

$$\mathscr{O} \asymp \sum_{n=0}^{\infty} c_n g^n$$

➾ how should we interpret the relation?

- The perturbative expansion is unlikely to converge, **Simon,** 1972
- the asymptotic behavior of the coefficients is expected to be

$$c_n \quad \sim \quad K\, n^{\alpha}\, \frac{n\,!}{S^n}, \quad n \to \infty$$ **Vainshtein** 1994

The requirement ($\asymp$) is not a formal one, it has a physical content: it means that there is a smooth transition between the system with interaction and the system without it, **Fischer** 1995. Furthermore, Borel and Carleman proved that there are analytic functions corresponding to arbitrary asymptotic power series.

We (**A. David and I**) did not answer general questions (e.g. to prove uniqueness) but concentrated (**arXiv:1307.1843 [hep-ph]**) on

predicting higher orders

using the well-known concept of " series acceleration", i.e. one of a collection of **sequence transformations (ST)** for improving the rate of convergence of a series.

- If the original series is divergent, the **ST** acts as an extrapolation method
- in the case of **infinite sums**, **ST**s have the effect that sums that formally diverge may return a result that can be interpreted as evaluation of the **analytic extension of the series for the sum.**
- the relation between **Borel summation** (usual method applied for summing divergent series) and these extrapolation methods is known Note that the definition of a sum of a **factorially divergent** series, including those with non-alternating coefficients, is always **equivalent** to Borel's definition, **Suslov** 2005

Example

$$S_\infty = \sum_{n=0}^{\infty} n!\, z^{n+1} = e^{-1/z}\,\mathrm{E}i\left(\frac{1}{z}\right)$$

where the **exponential integral** is a single-valued function in the plane cut along the negative real axis.
However, for $z > 0$ $\mathrm{E}i(z)$ can be computed to great accuracy using several Chebyshev expansions. Note that the r.h.s. is the **Borel sum of the series**.

Levin τ-transform, given the partial sum

$$S_n = \sum_{i=0}^{n} \gamma_i z^i, \quad \text{define the } \tau\text{-transform as}$$

$$\tau_k = \frac{N_k}{D_k},$$

$$N_k = \sum_{i=1}^{k} W(k,i)\, S_i, \qquad D_k = \sum_{i=1}^{m} W(k,i),$$

$$W(k,i) = (-1)^i \begin{pmatrix} k \\ i \end{pmatrix} \frac{(i)_{k-1}}{\Delta S_{i-1}}$$

where $(z)_a = \Gamma(z+a)/\Gamma(z)$ is the Pochhammer symbol and Δ is the usual forward-difference operator, $\Delta S_n = S_{n+1} - S_n$.

Weniger δ-transform

$$\delta_k(\beta) = \frac{\sum_{i=0}^{k} W^\delta(k,i,\beta)\, S_i}{\sum_{i=0}^{k} W^\delta(k,i,\beta)}$$

$$W^\delta(k,i,\beta) = (-1)^i \begin{pmatrix} k \\ i \end{pmatrix} \frac{(\beta+i)_{k-1}}{(\beta+k)_{k-1}} \frac{1}{\gamma_{i+1}\, z^{i+1}}$$

ST Examples & Details

The whole strategy is based on the fact that one can predict the *coefficients* by

- constructing an approximant with the known terms of the series $(\gamma_0, \ldots, \gamma_n)$ and
- expanding the approximant in a Taylor series. The first n terms of this series will exactly agree with those of the original series and

the subsequent terms may be treated as the *predicted coefficients*. i.e. if $S_1, \ldots, S_k$ are known, one computes

$$\therefore \; \tau_k - S_k = \bar{\gamma}_{k+1} z^{k+1} + \mathcal{O}\left(z^{k+2}\right)$$

➾ $\bar{\gamma}_{k+1}$ is the prediction for γ_{k+1}

How to use it?

Consider a specific example, $\mathbf{gg} \to \mathbf{H}$. Define

$$\sigma_{\mathrm{gg}}\left(\tau, M_{\mathrm{H}}^2\right) = \sigma_{\mathrm{gg}}^0\left(\tau, M_{\mathrm{H}}^2\right) K_{\mathrm{gg}}\left(\tau, M_{\mathrm{H}}^2, \alpha_{\mathrm{s}}\right)$$

where $\boldsymbol{\tau = M_{\mathrm{H}}^2/s}$ and $\boldsymbol{\sigma_{\mathrm{gg}}^0}$ is the LO cross section. The K-factor admits a formal power expansion in $\boldsymbol{\alpha_{\mathrm{s}}(\mu_{\mathrm{R}})}$

$$K_{\mathrm{gg}}\left(\tau, M_{\mathrm{H}}^2, \alpha_{\mathrm{s}}\right) = 1 + \sum_{n=1}^{\infty} \alpha_{\mathrm{s}}^n(\mu_{\mathrm{R}})\, K_{\mathrm{gg}}^n$$

Known coefficients are **11.879** and **72.254**

In their recent work, **Ball et al, (Ball:2013bra)** computed (at $\sqrt{s} = 8$ ***TeV***)

$$\alpha_s^3\left(\frac{M_H}{2}\right) K_{gg}^3\left(\mu = \frac{M_H}{2}\right) = 0.323 \pm 0.059$$

$$\alpha_s^3\left(M_H\right) K_{gg}^3\left(\mu = M_H\right) = 0.527 \pm 0.043$$

$$\alpha_s^3\left(2\,M_H\right) K_{gg}^3\left(\mu = 2\,M_H\right) = 0.729 \pm 0.032$$

Warming up with two coefficients $\tau_2 - S_2 = \frac{\gamma_2^2}{\gamma_1} z^3 + \mathscr{O}\left(z^4\right)$

☛

$$\mathbf{346.42} \leqq \gamma_3\left(\mu = M_H\right) \leqq \mathbf{407.48} \qquad \text{(Ball:2013bra)}$$

$$\bar{\gamma}_3\left(\mu = M_H\right) = \mathbf{439.48} \quad \Lleftarrow \quad \text{predicted}$$

✳ which has the correct sign and the right order of magnitude.

Introducing

$$S_{\mathrm{N},n} \;=\; \sum_{k=0}^{n} \gamma_k\, z^k + \sum_{k=n+1}^{\mathrm{N}} \bar{\gamma}_k\, z^k,$$

and $\boldsymbol{\delta}_{\mathbf{N},\boldsymbol{n}}$ etc, constructed accordingly, our strategy for estimating MHO and MHOU can be summarized as follows:

- we select a scale, $\mu = M_{\mathrm{H}}$ for $\mathbf{gg}$-fusion
- ESTIMATE THE UNCERTAINTY DUE TO HIGHER ORDERS AT THAT SCALE, I.E. THE (SCALE VARIATION) UNCERTAINTY AT THE CHOSEN SCALE IS PART OF THE UNCERTAINTY DUE TO HIGHER ORDERS AND SHOULD NOT BE COUNTED TWICE

∴ we compare

$$\sigma^{S,n}_{gg} = \sigma^{0}_{gg}\left(\mu = M_H\right) S_{n,3}\left(\mu = M_H\right)$$
$$\sigma^{\delta,n}_{gg} = \sigma^{0}_{gg}\left(\mu = M_H\right) \delta_{n,3}\left(\mu = M_H\right)$$

Our conclusion is that, to a very good accuracy,

$$\sigma_{gg} \in \left[\sigma^{S,3}_{gg}\, ,\, \sigma^{\delta,5}_{gg}\right]$$

with a flat interval of **16.37%**.

The uncertainty on the width, induced by the error on the coefficient $\gamma_3\left(\mu = M_H\right)$ brings it to 26.01%

✂ N^{3}LO & QCD scales var.

$\sigma_{gg} \in [18.90\, ,\, 21.93]\ pb$

completion & MHO ★★★

$\sigma_{gg} \in [20.13\, ,\, 23.42]\ pb$

NNLO $\twoheadrightarrow$ $\boxed{+17\%}$ $\twoheadrightarrow$ N^{3}LO $\twoheadrightarrow$ $\approx$ $\boxed{+7\%}$ $\twoheadrightarrow$ completion

The advantages of the method are that

- the result does not depend on the choice of the parameter expansion (it is based on "PARTIAL SUMS)" ✓
- it takes into account the nature of the coefficients, i.e. that the **known terms** of the perturbative expansion in gg-fusion are **positive** ✓
- The corresponding **pdf** could be derived by following the work of Cacciari and Houdeau ✓

BU1

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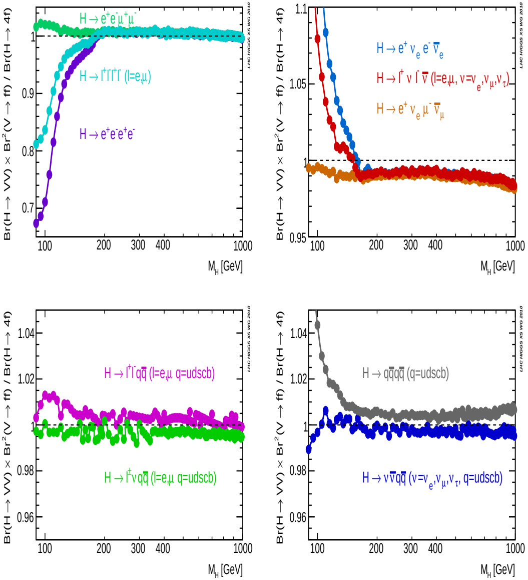
BU1

Pseudo Observables

What does the term "Higgs decay" mean? A mathematical expression? But what does it mean for such an expression to exist in the physical world? Trying to answer that question immediately raises other questions about the correspondence between mathematical objects and the physical world

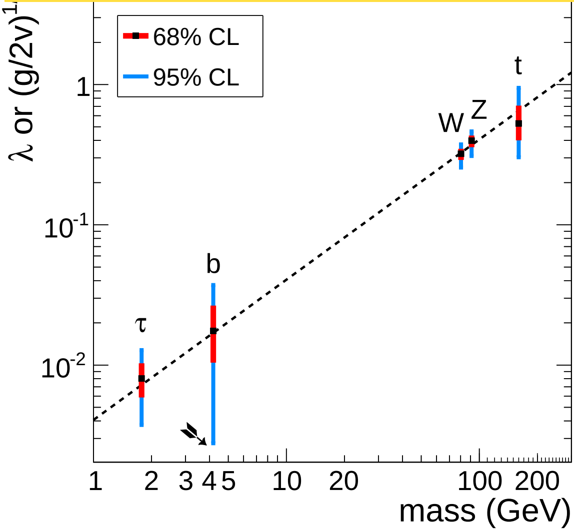
from PROPHECY4F



These plots are one of the best examples that

$$\mathrm{BR}\,(\mathrm{H} \to \mathrm{VV}) \quad \otimes \quad \mathrm{BR}^2\,(V \to \bar{\mathrm{f}}\mathrm{f})$$
$$\neq$$
$$\mathrm{BR}\,(\mathrm{H} \to 4\,\mathrm{f})$$

Trivial but true, $\looparrowright$ $\mathbf{H} \to \mathbf{VV}$ is not a physical OBSERVABLE, eventually it can be *defined* as " PSEUDO-OBSERVABLE" censored


λ or (g/2v)^1/2
68% CL
95% CL
1
10^-1
10^-2
t
W
Z
b
τ
1
2
3
4
5
10
20
100
200
mass (GeV)

The previous plot (**couplings** $\rightleftarrows$ **masses**) shows that

POs can be defined (couplings) **Iff** the rules of the game are respected ☛

- MODEL-INDEPENDENT couplings are *extracted* in some *effective* way that includes **QCD** but not **NLO EW**
- If one wants to obtain the **SM** (the straight line) $\looparrowright$ use RUNNING MASSES $m_{\mathrm{f}}(M_{\mathrm{H}})$

Theorem

$\nexists \qquad \mathrm{H} \to \mathrm{Z} + \gamma,\ \mathrm{H} \to \mathrm{VV}$ *etc.* *do not exist since*

✂$\mathbf{V} \quad \notin \qquad |\, \mathbf{in/out} >$ *bases of the Hilbert space*

High Precision Road

Dalitz Decay?

$M_{\mathrm{H}} = 125.5\ GeV$ $\qquad$ $\mathrm{BR}\left(\mathrm{H} \to e^+e^-\right) = 5.1 \times 10^{-9}$

while a *naive* estimate gives

$$\mathrm{BR}\left(\mathrm{H} \to \mathrm{Z}\gamma\right) \mathrm{BR}\left(\mathrm{Z} \to e^+e^-\right) = 5.31 \times 10^{-5}$$

4 orders of MAGNITUDE larger

How much is the corresponding PO extracted from full Dalitz Decay?

We could expect $\boldsymbol{\Gamma\left(\mathrm{H} \to e^+e^-\gamma\right) = 5.7\%\,\Gamma\left(\mathrm{H} \to \gamma\gamma\right)}$ but photon isolation must be discussed.

Categories

Terminology:

The name **Dalitz Decay** *must be reserved for the full process*

$$\mathbf{H \to \bar{f}f\gamma}$$

Subcategories:

$$\begin{cases} H \to Z^* \left(\to \bar{f}f\right) + \gamma & \text{✂ unphysical}^1 \\ H \to \gamma^* \left(\to \bar{f}f\right) + \gamma & \text{✂ unphysical} \\ H \to Z_c \left(\to \bar{f}f\right) + \gamma & \text{PO}^2 \end{cases}$$

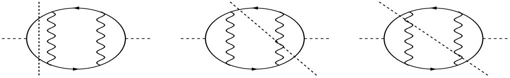
[1] Z^* is the off-shell Z

[2] Z_c is the Z at its complex pole

Understanding the problem

$$\boxed{\mathbf{H} \to \bar{\mathbf{f}}\mathbf{f} \quad \text{or} \quad \mathbf{H} \to \bar{\mathbf{f}}\mathbf{f} + n\gamma\,?}$$

Go to two-loop, the process is considerably more complex than, say, $\mathbf{H} \to \gamma\gamma$ *(***QED** *and* **QCD** *corrections)*. Think in terms of cuts of the three-loop $\mathbf{H}$ self-energy ☛



Moral: *Unless you* **Isolate photons**

you don't know which process you are talking about

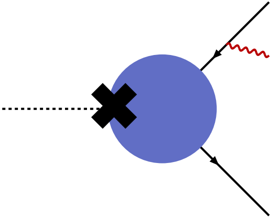
$\mathbf{H} \to \bar{\mathbf{f}}\mathbf{f}$ NNLO or $\mathbf{H} \to \bar{\mathbf{f}}\mathbf{f}\gamma$ NLO

The complete ***S***-matrix element will read as follows:

$$\begin{aligned} S &= \left|A^{(0)}\left(\mathrm{H}\to\bar{\mathrm{f}}\mathrm{f}\right)\right|^2 \\ &+ 2\,\mathrm{Re}\left[\mathrm{A}^{(0)}\left(\mathrm{H}\to\bar{\mathrm{f}}\mathrm{f}\right)\right]^\dagger \mathrm{A}^{(1)}\left(\mathrm{H}\to\bar{\mathrm{f}}\mathrm{f}\right) \\ &+ \left|A^{(0)}\left(\mathrm{H}\to\bar{\mathrm{f}}\mathrm{f}\gamma\right)\right|^2 ✗ \\ &+ 2\,\mathrm{Re}\left[\mathrm{A}^{(0)}\left(\mathrm{H}\to\bar{\mathrm{f}}\mathrm{f}\right)\right]^\dagger \mathrm{A}^{(2)}\left(\mathrm{H}\to\bar{\mathrm{f}}\mathrm{f}\right) \\ &+ 2\,\mathrm{Re}\left[\mathrm{A}^{(0)}\left(\mathrm{H}\to\bar{\mathrm{f}}\mathrm{f}\gamma\right)\right]^\dagger \mathrm{A}^{(1)}\left(\mathrm{H}\to\bar{\mathrm{f}}\mathrm{f}\gamma\right) ✗ \\ &+ \left|A^{(0)}\left(\mathrm{H}\to\bar{\mathrm{f}}\mathrm{f}\gamma\gamma\right)\right|^2. \end{aligned}$$

Don't get trapped by your intuition, the **IR/collinear** stuff will not survive in the limit $m_f \to 0$

There are **genuinely non-QED(QCD)** terms surviving the **zero-Yukawa** limit (a result known since the '80s)



$2f\,H$-BRs below $\mathbf{10^{-3}}$–$\mathbf{10^{-4}}$ pose additional TH problems Δ**BR** $\gg$ **BR**

Only the total *Dalitz Decay* has a meaning and *can be differentiated through cuts*

- The most important is the definition of *visible photon* to distinguish between $\overline{f}f$ and $\overline{f}f\gamma$
- Next cuts are on $M(\overline{f}f)$ to *isolate* pseudo-observables
    - with a small window around the Z-peak the pseudo-observable $H \to Z_c\gamma$ can be enhanced, but there is a contamination due to many non-resonant backgrounds ✓
    - at small di-lepton invariant masses γ^* dominates ✓

Distributions

Results: leptons

$$m\left(\bar{\mathrm{f}}\mathrm{f}\right) > 0.1\, M_{\mathrm{H}} \qquad m\left(\mathrm{f}\gamma\right) > 0.1\, M_{\mathrm{H}} \qquad m\left(\bar{\mathrm{f}}\gamma\right) > 0.1\, M_{\mathrm{H}}$$

$$\Gamma_{\mathrm{NLO}} = 0.233\ keV \quad \oplus \quad \begin{cases} \Gamma_{\mathrm{LO}} = 0.29 \times 10^{-6}\ keV & \mathrm{e} \\ \Gamma_{\mathrm{LO}} = 0.012\ keV & \mu \\ \Gamma_{\mathrm{LO}} = 3.504\ keV & \tau \end{cases}$$

☛ LO and NLO **do not interfere** (as long as masses are neglected in NLO), they belong to different helicity sets.

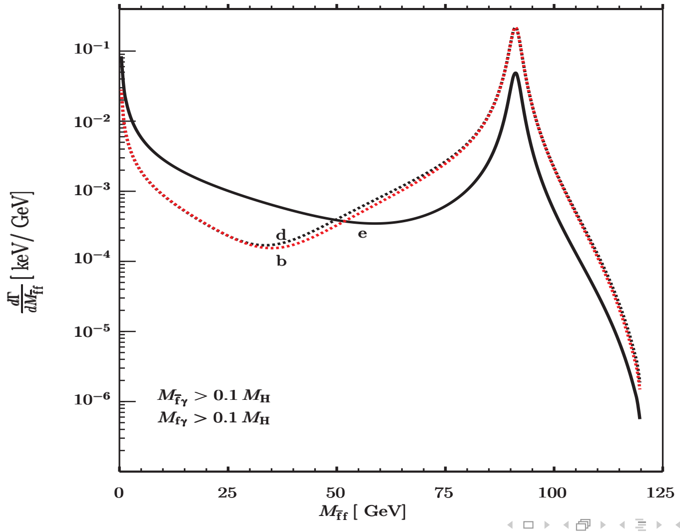
☛ Cuts à la **Dicus and Repko**

Results: quarks

$$m\left(\bar{\mathrm{f}}\mathrm{f}\right) > 0.1\,M_{\mathrm{H}} \qquad m\left(\mathrm{f}\gamma\right) > 0.1\,M_{\mathrm{H}} \qquad m\left(\bar{\mathrm{f}}\gamma\right) > 0.1\,M_{\mathrm{H}}$$

$$\begin{cases} \Gamma_{\mathrm{LO}} = 0.013\ keV & \Gamma_{\mathrm{NLO}} = 0.874\ keV \quad \mathrm{d} \\ \Gamma_{\mathrm{LO}} = 8.139\ keV & \Gamma_{\mathrm{NLO}} = 0.866\ keV \quad \mathrm{b} \end{cases}$$

☛ Note the effect of m_t


$\frac{d\Gamma}{dM_{\bar{f}f}}$ [keV/ GeV]
10^{-1}
10^{-2}
10^{-3}
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d
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$M_{\bar{f}\gamma} > 0.1\, M_H$
$M_{f\gamma} > 0.1\, M_H$
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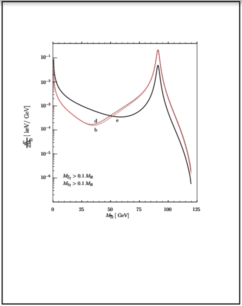
$M_{\bar{f}f}$ [GeV]



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dΓ/dM_f̄f [keV/ GeV]
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M_fγ > 0.1 M_H
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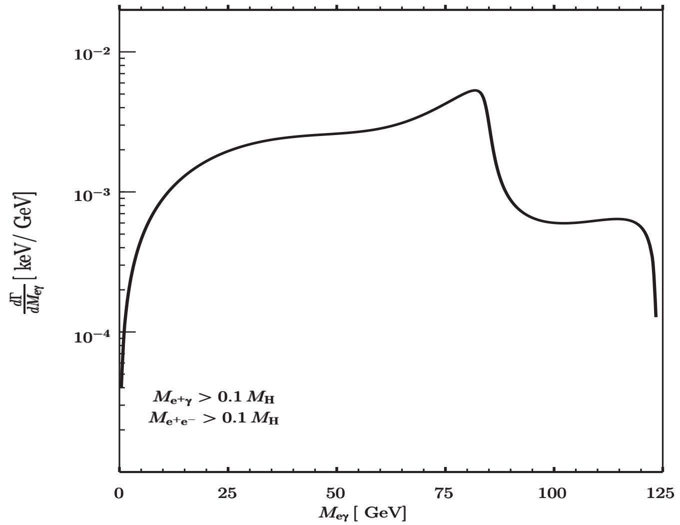
M_f̄f [GeV]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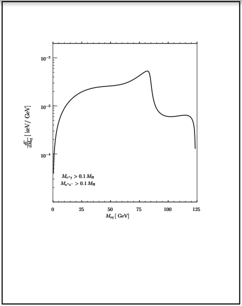
$M_{e\gamma}$ [GeV]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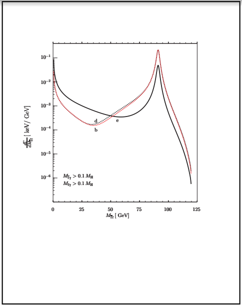
$M_{e\gamma}$ [GeV]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$\frac{d\Gamma}{dM_{\bar{f}f}}$ [keV/ GeV]
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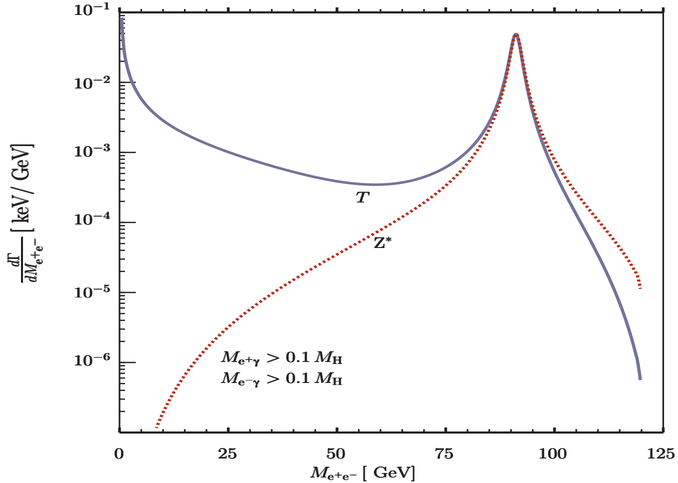
$M_{\bar{f}f}$ [GeV]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$\frac{d\Gamma}{dM_{e^+e^-}}$ [keV/ GeV]
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Z^*
$M_{e^+\gamma} > 0.1\, M_H$
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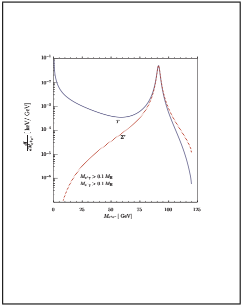
$M_{e^+e^-}$ [GeV]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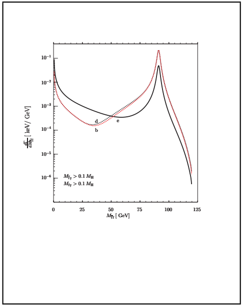
$\frac{d\Gamma}{dM_{e^+e^-}}$ [keV/ GeV]
$M_{e^+e^-}$ [GeV]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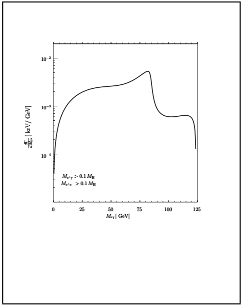
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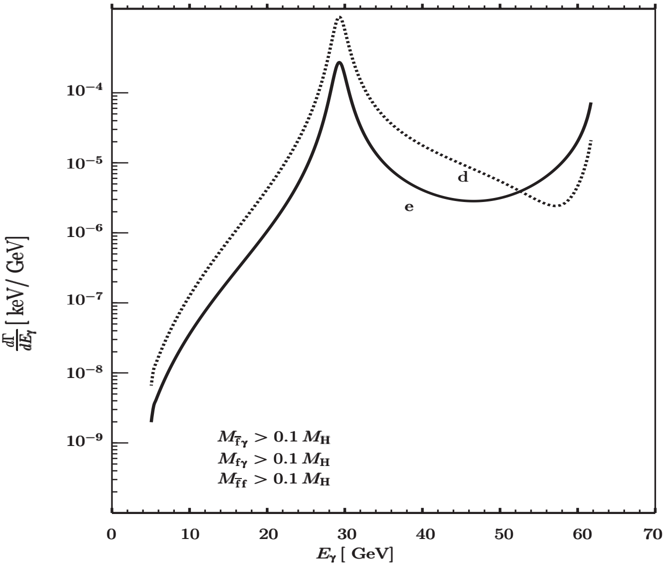
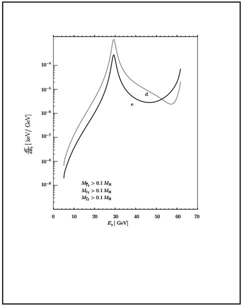
$\frac{d\Gamma}{dM_{e\gamma}}$ [keV/ GeV]
$M_{e\gamma}$ [GeV]



10^{-4}
10^{-5}
10^{-6}
10^{-7}
10^{-8}
10^{-9}
$\frac{d\Gamma}{dE_\gamma}$ [keV/ GeV]
d
e
$M_{\bar{f}\gamma} > 0.1\, M_H$
$M_{f\gamma} > 0.1\, M_H$
$M_{\bar{f}f} > 0.1\, M_H$
0 10 20 30 40 50 60 70
E_γ [GeV]


$\frac{d\Gamma}{dE_\gamma}$ [keV/ GeV]
E_γ [GeV]
d
e
$M_{\bar{f}\gamma} > 0.1\, M_H$
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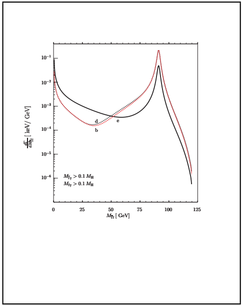
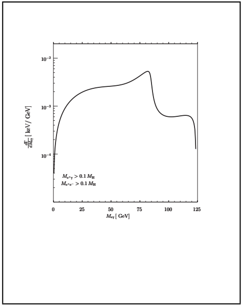
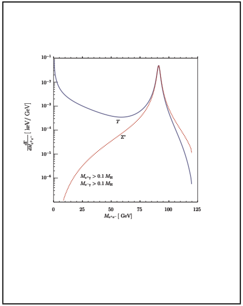
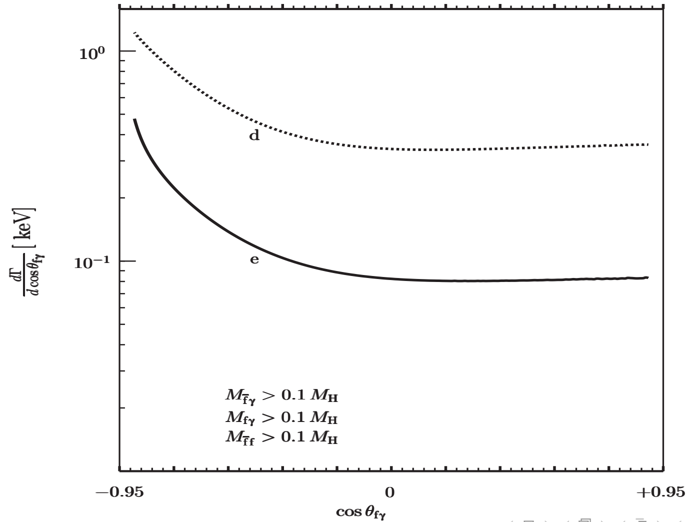
$M_{\bar{f}f} > 0.1\, M_H$
0 10 20 30 40 50 60 70

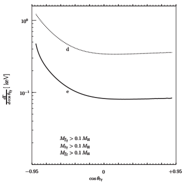

$\frac{d\Gamma}{dM_{\bar{f}f}}$ [keV/ GeV]
$M_{\bar{f}f}$ [GeV]
d
b
e
$M_{\bar{f}\gamma} > 0.1\, M_H$
$M_{f\gamma} > 0.1\, M_H$
0 25 50 75 100 125


$\frac{d\Gamma}{dM_{e\gamma}}$ [keV/ GeV]
$M_{e\gamma}$ [GeV]
$M_{e^+\gamma} > 0.1\, M_H$
$M_{e^+e^-} > 0.1\, M_H$
0 25 50 75 100 125

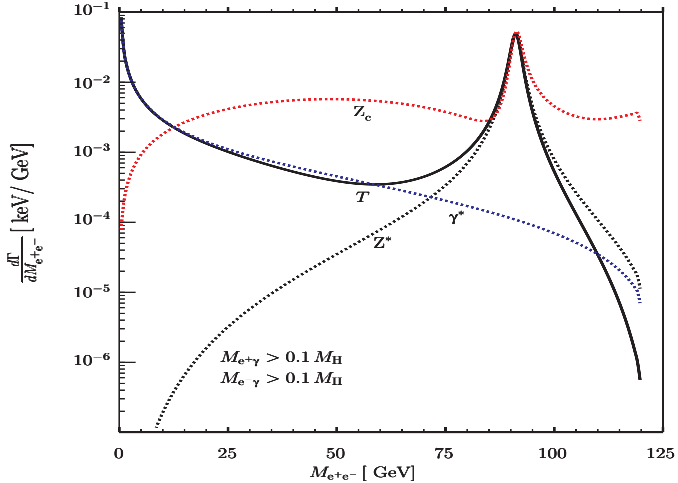

$\frac{d\Gamma}{dM_{e^+e^-}}$ [keV/ GeV]
$M_{e^+e^-}$ [GeV]
T
Z^*
$M_{e^+\gamma} > 0.1\, M_H$
$M_{e^-\gamma} > 0.1\, M_H$
0 25 50 75 100 125

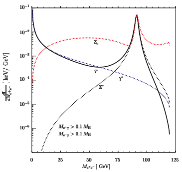


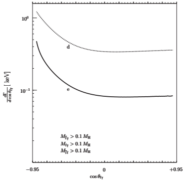
10^0
10^{-1}
$\frac{d\Gamma}{d\cos\theta_{f\gamma}}$ [keV]
d
e
$M_{\bar{f}\gamma} > 0.1\, M_H$
$M_{f\gamma} > 0.1\, M_H$
$M_{\bar{f}f} > 0.1\, M_H$
-0.95
0
$+0.95$
$\cos\theta_{f\gamma}$

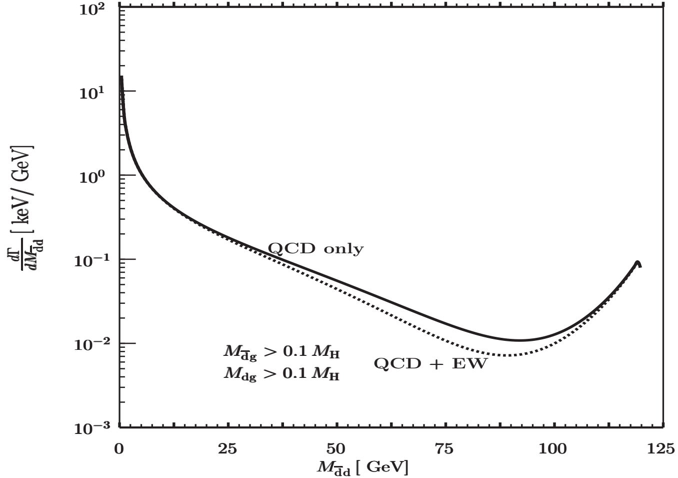


10^{0}
10^{-1}
d
e
$\frac{d\Gamma}{d\cos\theta_{f\gamma}}$ [keV]
$M_{\bar{f}\gamma} > 0.1\, M_H$
$M_{f\gamma} > 0.1\, M_H$
$M_{\bar{f}f} > 0.1\, M_H$
-0.95
0
$+0.95$
$\cos\theta_{f\gamma}$

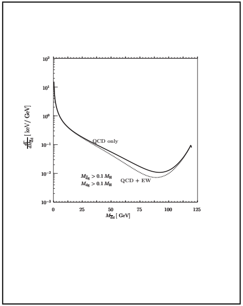

10^{-1}
10^{-2}
10^{-3}
10^{-4}
10^{-5}
10^{-6}
$\frac{d\Gamma}{dM_{e^+e^-}}$ [keV / GeV]
Z_c
T
γ^*
Z^*
$M_{e^+\gamma} > 0.1\, M_H$
$M_{e^-\gamma} > 0.1\, M_H$
0
25
50
75
100
125
$M_{e^+e^-}$ [GeV]

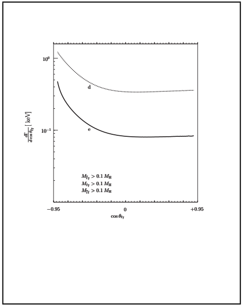

10^{-1}
10^{-2}
10^{-3}
10^{-4}
10^{-5}
10^{-6}
$\frac{d\Gamma}{dM_{e^+e^-}}$ [keV/ GeV]
Z_c
T
γ^*
Z^*
$M_{e^+\gamma} > 0.1\, M_H$
$M_{e^-\gamma} > 0.1\, M_H$
0
25
50
75
100
125
$M_{e^+e^-}$ [GeV]

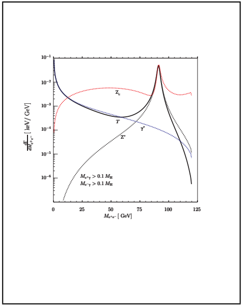

10^{0}
10^{-1}
$\frac{d\Gamma}{d\cos\theta_{f\gamma}}$ [keV]
d
e
$M_{\bar{f}\gamma} > 0.1\, M_H$
$M_{f\gamma} > 0.1\, M_H$
$M_{\bar{f}f} > 0.1\, M_H$
-0.95
0
$+0.95$
$\cos\theta_{f\gamma}$

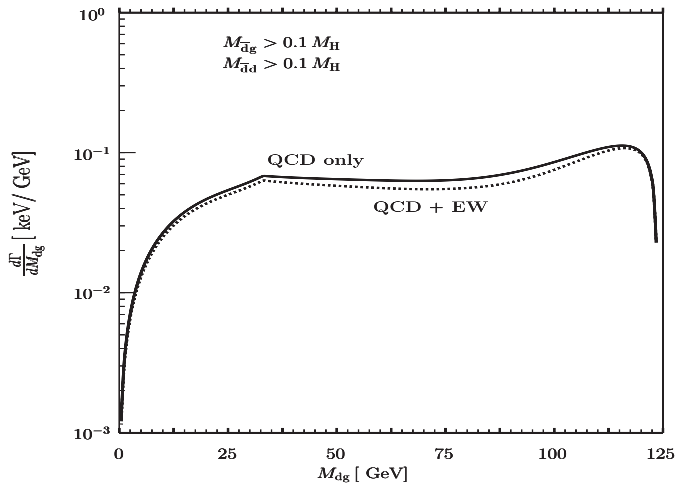


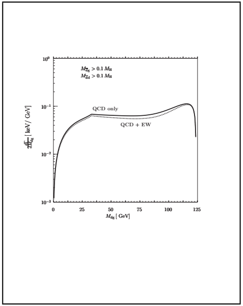
$\frac{d\Gamma}{dM_{\bar{d}d}}$ [keV/ GeV]
10^{2}
10^{1}
10^{0}
10^{-1}
10^{-2}
10^{-3}
0
25
50
75
100
125
$M_{\bar{d}d}$ [GeV]
QCD only
QCD + EW
$M_{\bar{d}g} > 0.1\, M_H$
$M_{dg} > 0.1\, M_H$

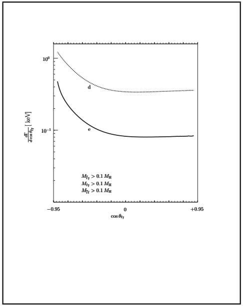

QCD only
QCD + EW
0
25
50
75
100
125

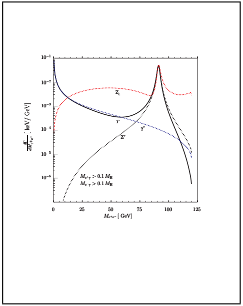

d
e
−0.95
0
+0.95

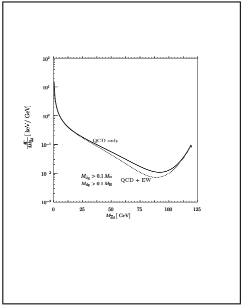

T
0
25
50
75
100
125

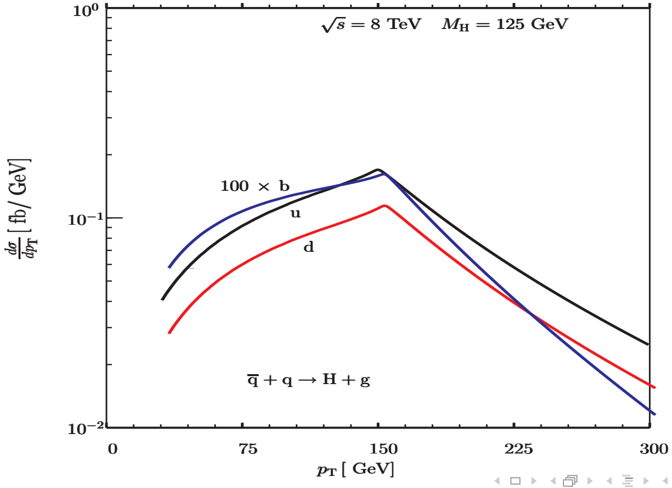

$M_{\overline{d}g} > 0.1\, M_H$
$M_{\overline{d}d} > 0.1\, M_H$
QCD only
QCD + EW
$\frac{d\Gamma}{dM_{dg}}$ [keV/ GeV]
M_{dg} [GeV]
10^0
10^{-1}
10^{-2}
10^{-3}
0
25
50
75
100
125

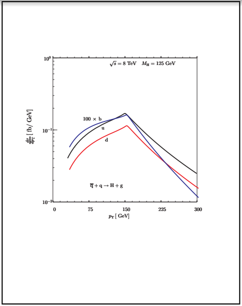

QCD only
QCD + EW
0
25
50
75
100
125

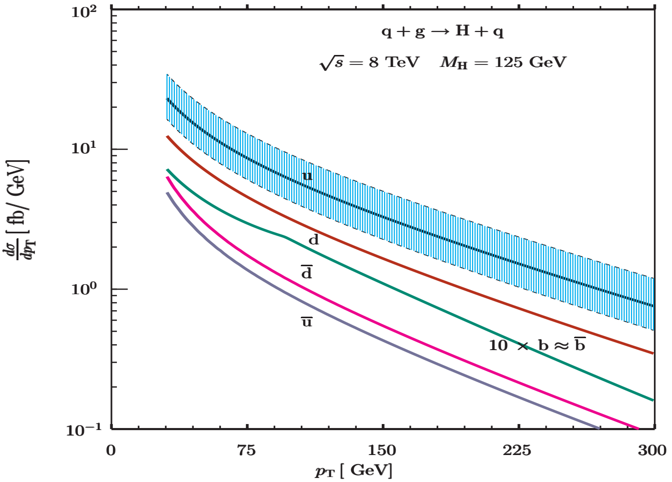

d
e
−0.95
0
+0.95

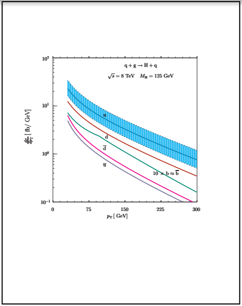

T
0
25
50
75
100
125

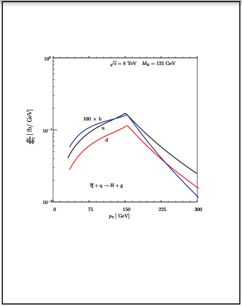

QCD only
QCD + EW
0
25
50
75
100
125

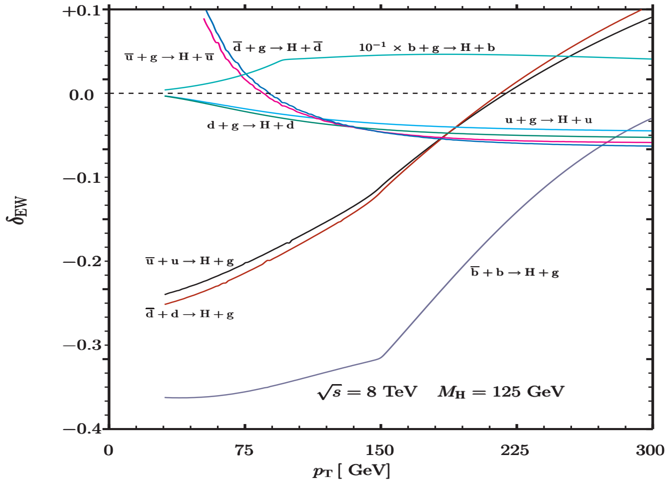

$\sqrt{s} = 8$ TeV $M_H = 125$ GeV
10^0
10^{-1}
10^{-2}
100 × b
u
d
$\overline{q} + q \to H + g$
0
75
150
225
300
p_T [GeV]
$\frac{d\sigma}{dp_T}$ [fb/ GeV]

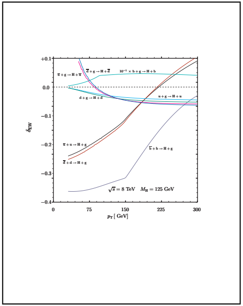

$\sqrt{s} = 8$ TeV $M_H = 125$ GeV
$\frac{d\sigma}{dp_T}$ [fb/ GeV]
10^0
10^{-1}
10^{-2}
100 × b
u
d
$\overline{q} + q \to H + g$
0
75
150
225
300
p_T [GeV]

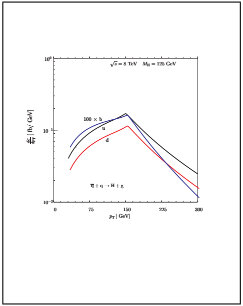

$q + g \to H + q$
$\sqrt{s} = 8$ TeV $M_H = 125$ GeV
u
d
$\overline{d}$
$\overline{u}$
$10 \times b \approx \overline{b}$
$\frac{d\sigma}{dp_T}$ [fb/ GeV]
p_T [GeV]
10^2
10^1
10^0
10^{-1}
0
75
150
225
300

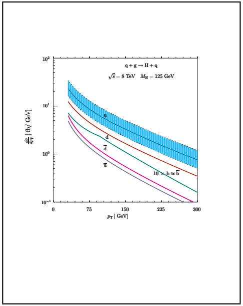

$q + g \to H + q$
$\sqrt{s} = 8$ TeV $M_H = 125$ GeV
$\frac{d\sigma}{dp_T}$ [fb/ GeV]
p_T [GeV]
u
d
$\overline{d}$
$\overline{u}$
$10 \times b \approx \overline{b}$
10^2
10^1
10^0
10^{-1}
0
75
150
225
300

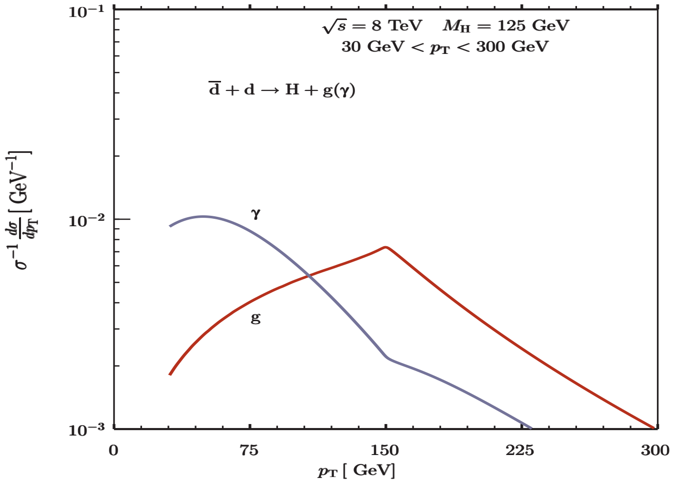

$\sqrt{s} = 8$ TeV $M_H = 125$ GeV
$\frac{d\sigma}{dp_T}$ [fb/ GeV]
p_T [GeV]
$100 \times b$
u
d
$\overline{q} + q \to H + g$
10^0
10^{-1}
10^{-2}
0
75
150
225
300

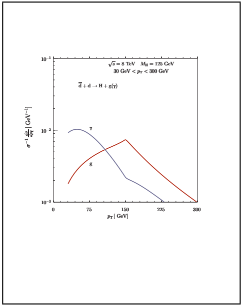

+0.1
0.0
−0.1
−0.2
−0.3
−0.4
δ_{EW}
0
75
150
225
300
p_T [GeV]
$\bar{u} + g \to H + \bar{u}$
$\bar{d} + g \to H + \bar{d}$
$10^{-1} \times b + g \to H + b$
$u + g \to H + u$
$d + g \to H + d$
$\bar{u} + u \to H + g$
$\bar{d} + d \to H + g$
$\bar{b} + b \to H + g$
$\sqrt{s} = 8$ TeV $M_H = 125$ GeV

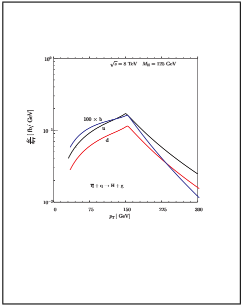

$\bar{u} + g \to H + \bar{u}$
$\bar{d} + g \to H + \bar{d}$
$10^{-1} \times b + g \to H + b$
$u + g \to H + u$
$d + g \to H + d$
$\bar{u} + u \to H + g$
$\bar{d} + d \to H + g$
$\bar{b} + b \to H + g$
$\sqrt{s} = 8$ TeV $M_H = 125$ GeV
δ_{EW}
p_T [GeV]

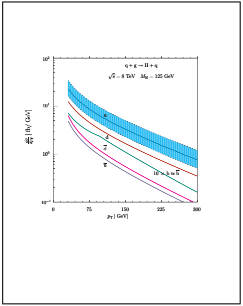

$\sqrt{s} = 8$ TeV $M_H = 125$ GeV
$100 \times b$
u
d
$\bar{q} + q \to H + g$
$\frac{d\sigma}{dp_T}$ [fb/ GeV]
p_T [GeV]

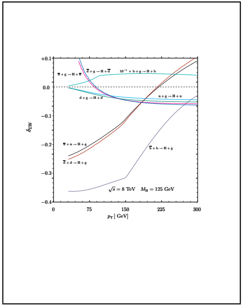

$q + g \to H + q$
$\sqrt{s} = 8$ TeV $M_H = 125$ GeV
u
d
$\bar{d}$
$\bar{u}$
$10 \times b \approx \bar{b}$
$\frac{d\sigma}{dp_T}$ [fb/ GeV]
p_T [GeV]

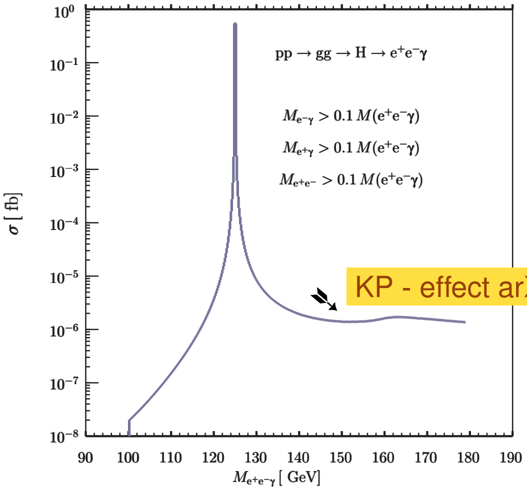

$\sqrt{s} = 8$ TeV $M_H = 125$ GeV
30 GeV $< p_T <$ 300 GeV
$\bar{d} + d \to H + g(\gamma)$
γ
g
10^{-1}
10^{-2}
10^{-3}
$\sigma^{-1} \frac{d\sigma}{dp_T}$ [GeV^{-1}]
0
75
150
225
300
p_T [GeV]


$\sqrt{s} = 8$ TeV $M_H = 125$ GeV
30 GeV < p_T < 300 GeV
$\overline{d} + d \to H + g(\gamma)$
γ
g
$\sigma^{-1} \frac{d\sigma}{dp_T}$ [GeV^{-1}]
p_T [GeV]


$\sqrt{s} = 8$ TeV $M_H = 125$ GeV
100 × b
u
d
$\overline{q} + q \to H + g$
$\frac{d\sigma}{dp_T}$ [fb/ GeV]
p_T [GeV]


$q + g \to H + q$
$\sqrt{s} = 8$ TeV $M_H = 125$ GeV
u
d
$\overline{d}$
$\overline{u}$
$10 \times b \approx \overline{b}$
$\frac{d\sigma}{dp_T}$ [fb/ GeV]
p_T [GeV]


δ_{EW}
$\sqrt{s} = 8$ TeV $M_H = 125$ GeV
p_T [GeV]



KP - effect arXiv:1206.4803

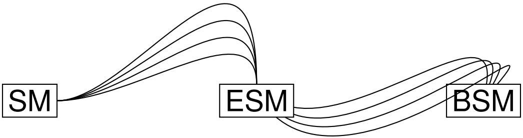
Observable	Pseudo-Observable
$\mathrm{H} \to \gamma\gamma$	
$\mathrm{H} \to \bar{\mathrm{f}}\mathrm{f}\gamma$	☞ $\mathrm{H} \to \mathrm{Z}\gamma$
$\mathrm{H} \to \bar{\mathrm{f}}\mathrm{f}$	
$\mathrm{H} \to \bar{\mathrm{f}}\mathrm{f}\bar{\mathrm{f}}'\mathrm{f}'$	☞ $\mathrm{H} \to \mathrm{VV}, \mathrm{Z}\gamma$

One needs to define when it is **4f** final state and when it is PAIR CORRECTION to **2f** final state (as it was done at LEP2)

Effective Field Theory

Renormalization - group view of the world

Let's consider the following path



The ontology of the SM on its scale should be understood as arising from the “emergent” effects of a more fundamental BSM at a finer scale

$$\mathscr{L}_{\rm ESM} = \mathscr{L}_{\rm SM} + \sum_{n>4} \sum_{i=1}^{N_n} \frac{a_i^n}{\Lambda^{n-4}} \mathscr{O}_i^{(d=n)}$$

$$\exists(\exists!) \quad \mathscr{L}_{\rm UCSM} \rightarrowtail \mathscr{O}_i\,?$$

UV completion of the SM (UCSM) or ESM?

Bottom-up or top-down approach to **ESM**?

- How many facts the theory explains: it is a draw
- Having the fewer auxiliary hypothesis: **SM** ↠ **UCSM** superior

The regulative ideal of an ultimate theory remains a powerful *aesthetic ingredient*

Strategy: How to interpret κ_X?

① *measure* $\boldsymbol{\kappa}$

$$\frac{\Gamma_{gg}}{\Gamma^{SM}_{gg}(m_H)} = \frac{\kappa_t^2 \cdot \Gamma^{tt}_{gg}(m_H) + \kappa_b^2 \cdot \Gamma^{bb}_{gg}(m_H) + \kappa_t \kappa_b \cdot \Gamma^{tb}_{gg}(m_H)}{\Gamma^{tt}_{gg}(m_H) + \Gamma^{bb}_{gg}(m_H) + \Gamma^{tb}_{gg}(m_H)}$$

② find $\mathcal{O}_i \Leftrightarrow \kappa_X$ (epistemological stop, true ESM believers stop here)

$$\mathscr{L}_{ESM} = \mathscr{L}_{SM} + \sum_{n>4} \sum_{i=1}^{N_n} \frac{a_i^n}{\Lambda^{n-4}} \mathcal{O}_i^{(d=n)}$$

③ find $\{\mathscr{L}_{BSM}\}$ that produces $\mathcal{O}_i$

κ_X **cannot** be arbitrary shifts of the **SM** diagrams $\looparrowright$ they require an **underlying** (at least **effective**) theory

We (**Yellow Report HXSWG vol. 3**) define an **EFT** based on

a **linear** representation of the EW gauge symmetry with a Higgs-doublet field, restricting ourselves to **dimension**-**6** operators relevant for Higgs physics **Buchmuller:1985jz, Grzadkowski:2010es.**

In a complete analysis all 59 **independent operators of Grzadkowski:2010es, including** 25 **four-fermion operators, have to be considered in addition to the selected** 34 **operators In weakly interacting theories the dimension-6 operators involving field strengths can only result from loops, while the others also result from tree diagrams (Arzt:1994gp)**

Lagrangian

$$\mathscr{L}_{\rm eff} = \mathscr{L}_{\rm SM}^{(4)} + \frac{1}{\Lambda^2}\sum_k \alpha_k \mathscr{O}_k,$$

$$\begin{aligned}\mathscr{L}_{\rm SM}^{(4)} = & -\frac{1}{4}G^A_{\mu\nu}G^{A\mu\nu} - \frac{1}{4}\mathrm{W}^I_{\mu\nu}\mathrm{W}^{I\mu\nu} - \frac{1}{4}\mathrm{B}_{\mu\nu}\mathrm{B}^{\mu\nu} \\ & + (D_\mu\Phi)^\dagger(D^\mu\Phi) + m^2\Phi^\dagger\Phi - \frac{1}{2}\lambda(\Phi^\dagger\Phi)^2 \\ & + i\bar{l}\not{D}l + i\bar{e}\not{D}e + i\bar{q}\not{D}q + i\bar{u}\not{D}u + i\bar{d}\not{D}d \\ & - (\bar{l}\Gamma_e e\Phi + \bar{q}\Gamma_u u\tilde{\Phi} + \bar{d}\Gamma_d d\Phi + \text{h.c.}),\end{aligned}$$

Operators

Φ^6 *and* $\Phi^4 D^2$	$\psi^2\Phi^3$	X^3
$\mathscr{O}_\Phi = (\Phi^\dagger\Phi)^3$	$\mathscr{O}_{e\Phi} = (\Phi^\dagger\Phi)(\bar{l}\Gamma_e e\Phi)$	$\mathscr{O}_G = f^{ABC}G^{A\nu}_\mu G^{B\rho}_\nu G^{C\mu}_\rho$
$\mathscr{O}_{\Phi\Box} = (\Phi^\dagger\Phi)\Box(\Phi^\dagger\Phi)$	$\mathscr{O}_{u\Phi} = (\Phi^\dagger\Phi)(\bar{q}\Gamma_u u\tilde{\Phi})$	$\mathscr{O}_{\tilde{G}} = f^{ABC}\tilde{G}^{A\nu}_\mu G^{B\rho}_\nu G^{C\mu}_\rho$
$\mathscr{O}_{\Phi D} = (\Phi^\dagger D^\mu\Phi)^*(\Phi^\dagger D_\mu\Phi)$	$\mathscr{O}_{d\Phi} = (\Phi^\dagger\Phi)(\bar{q}\Gamma_d d\Phi)$	$\mathscr{O}_W = \varepsilon^{IJK}W^{I\nu}_\mu W^{J\rho}_\nu W^{K\mu}_\rho$
		$\mathscr{O}_{\tilde{W}} = \varepsilon^{IJK}\tilde{W}^{I\nu}_\mu W^{J\rho}_\nu W^{K\mu}_\rho$
$X^2\Phi^2$	$\psi^2 X\Phi$	$\psi^2\Phi^2 D$
$\mathscr{O}_{\Phi G} = (\Phi^\dagger\Phi)G^A_{\mu\nu}G^{A\mu\nu}$	$\mathscr{O}_{uG} = (\bar{q}\sigma^{\mu\nu}\frac{\lambda^A}{2}\Gamma_u u\tilde{\Phi})G^A_{\mu\nu}$	$\mathscr{O}^{(1)}_{\Phi l} = (\Phi^\dagger i\overleftrightarrow{D}_\mu\Phi)(\bar{l}\gamma^\mu l)$
$\mathscr{O}_{\Phi\tilde{G}} = (\Phi^\dagger\Phi)\tilde{G}^A_{\mu\nu}G^{A\mu\nu}$	$\mathscr{O}_{dG} = (\bar{q}\sigma^{\mu\nu}\frac{\lambda^A}{2}\Gamma_d d\Phi)G^A_{\mu\nu}$	$\mathscr{O}^{(3)}_{\Phi l} = (\Phi^\dagger i\overleftrightarrow{D}^I_\mu\Phi)(\bar{l}\gamma^\mu\tau^I l)$
$\mathscr{O}_{\Phi W} = (\Phi^\dagger\Phi)W^I_{\mu\nu}W^{I\mu\nu}$	$\mathscr{O}_{eW} = (\bar{l}\sigma^{\mu\nu}\Gamma_e e\tau^I\Phi)W^I_{\mu\nu}$	$\mathscr{O}_{\Phi e} = (\Phi^\dagger i\overleftrightarrow{D}_\mu\Phi)(\bar{e}\gamma^\mu e)$
$\mathscr{O}_{\Phi\tilde{W}} = (\Phi^\dagger\Phi)\tilde{W}^I_{\mu\nu}W^{I\mu\nu}$	$\mathscr{O}_{uW} = (\bar{q}\sigma^{\mu\nu}\Gamma_u u\tau^I\tilde{\Phi})W^I_{\mu\nu}$	$\mathscr{O}^{(1)}_{\Phi q} = (\Phi^\dagger i\overleftrightarrow{D}_\mu\Phi)(\bar{q}\gamma^\mu q)$
$\mathscr{O}_{\Phi B} = (\Phi^\dagger\Phi)B_{\mu\nu}B^{\mu\nu}$	$\mathscr{O}_{dW} = (\bar{q}\sigma^{\mu\nu}\Gamma_d d\tau^I\Phi)W^I_{\mu\nu}$	$\mathscr{O}^{(3)}_{\Phi q} = (\Phi^\dagger i\overleftrightarrow{D}^I_\mu\Phi)(\bar{q}\gamma^\mu\tau^I q)$
$\mathscr{O}_{\Phi\tilde{B}} = (\Phi^\dagger\Phi)\tilde{B}_{\mu\nu}B^{\mu\nu}$	$\mathscr{O}_{eB} = (\bar{l}\sigma^{\mu\nu}\Gamma_e e\Phi)B_{\mu\nu}$	$\mathscr{O}_{\Phi u} = (\Phi^\dagger i\overleftrightarrow{D}_\mu\Phi)(\bar{u}\gamma^\mu u)$
$\mathscr{O}_{\Phi WB} = (\Phi^\dagger\tau^I\Phi)W^I_{\mu\nu}B^{\mu\nu}$	$\mathscr{O}_{uB} = (\bar{q}\sigma^{\mu\nu}\Gamma_u u\tilde{\Phi})B_{\mu\nu}$	$\mathscr{O}_{\Phi d} = (\Phi^\dagger i\overleftrightarrow{D}_\mu\Phi)(\bar{d}\gamma^\mu d)$
$\mathscr{O}_{\Phi\tilde{W}B} = (\Phi^\dagger\tau^I\Phi)\tilde{W}^I_{\mu\nu}B^{\mu\nu}$	$\mathscr{O}_{dB} = (\bar{q}\sigma^{\mu\nu}\Gamma_d d\Phi)B_{\mu\nu}$	$\mathscr{O}_{\Phi ud} = i(\tilde{\Phi}^\dagger D_\mu\Phi)(\bar{u}\gamma^\mu\Gamma_{ud}d)$

Examples & Details

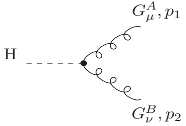
The parameters of the SM Lagrangian g, g', λ, m^2, and $\Gamma_{\rm f}$ keep their meaning in the presence of dimension-6 operators.

10.4.2 Higgs vertices

Here we list the most important Feynman rules for vertices involving exactly one physical Higgs boson. These are given in terms of the above-defined physical fields and parameters. In the coefficients of dimension-6 couplings we replaced v^2 by the Fermi constant via $v^2 = 1/(\sqrt{2}G_F)$.

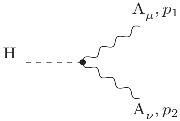
The triple vertices involving one Higgs boson read:

Hgg coupling:



$$= {\rm i}\frac{2g}{M_{\rm W}}\frac{1}{\sqrt{2}G_F\Lambda^2}\left[\alpha_{GG}(p_{2\mu}p_{1\nu} - p_1p_2g_{\mu\nu}) + \alpha_{G\widetilde{G}}\varepsilon_{\mu\nu\rho\sigma}p_1^\rho p_2^\sigma\right]\delta^{AB}, \tag{155}$$

HAA coupling:



$$= {\rm i}\frac{2g}{M_{\rm W}}\frac{1}{\sqrt{2}G_F\Lambda^2}\left[\alpha_{\rm AA}(p_{2\mu}p_{1\nu} - p_1p_2g_{\mu\nu}) + \alpha_{\rm A\widetilde{A}}\varepsilon_{\mu\nu\rho\sigma}p_1^\rho p_2^\sigma\right], \tag{156}$$

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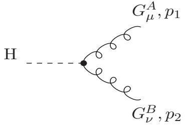

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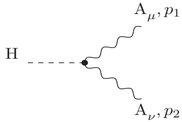

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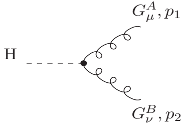

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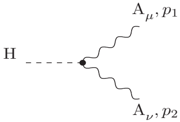

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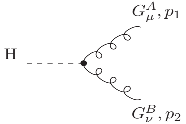

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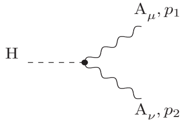

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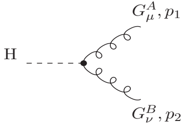

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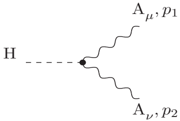

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Since κ_t, κ_b etc. can be made different **ONLY** by inserting $\mathcal{O}$ operators in **SM** vertices

Vademecum (**NLO + EFT**) trainee

- the **EFT** part has to be *implemented* into existing (**EW + QCD**) codes: formulation in arbitrary gauge (not U-gauge restricted) is needed
- Renormalization for the full **SM + EFT** Lagrangian is needed

Caveat

If one restricts the analysis to the calculation of **on-shell matrix** elements then additional operators are eliminated by the **Equations-Of-Motion** (EOM)

Only $\boldsymbol{S}$-matrix elements will be the same for equivalent operators but not the Green's functions $\therefore$

- since we are working with **unstable particles**,
- since we are inserting operators **inside loops** in a non-Abelian theory,
- since we want to use (off-shell) $\boldsymbol{S}, \boldsymbol{T}$ and $\boldsymbol{U}$ parameters to **constrain** the Wilson coefficients,

$\looparrowright$ the use of EOM should be taken with extreme caution

Caveat

$\boxed{\text{Note}}$ that for

$\Lambda \approx 5\ TeV$

we have

$$\boxed{1/(\sqrt{2}G_F\Lambda^2) \approx g^2/(4\pi)},$$

i.e. ➳ the contributions of $d = 6$ operators are $\approxeq$ loop effects.

➳ ➳ For higher scales, loop contributions tend to be more important ($\succcurlyeq$)

Example: SM loops dressed only with UV-admissible operators

$\mathrm{H} \to \gamma\gamma$

For $\mathrm{H} \to \gamma\gamma$ the SM amplitude reads

$$\mathscr{M}_{\mathrm{SM}} = F_{\mathrm{SM}} \left(\delta^{\mu\nu} + 2\,\frac{p_1^{\nu}\,p_2^{\mu}}{\overline{\mathrm{M}}_{\mathrm{H}}^{\,2}} \right) e_{\mu}\left(p_1\right) e_{\nu}\left(p_2\right)$$

$$F_{\mathrm{SM}} = -g\,\overline{\mathrm{M}}\,F_{\mathrm{SM}}^{\mathrm{W}} - \frac{1}{2}\,g\,\frac{M_{\mathrm{t}}^2}{\overline{\mathrm{M}}}\,F_{\mathrm{SM}}^{\mathrm{t}} - \frac{1}{2}\,g\,\frac{M_{\mathrm{b}}^2}{\overline{\mathrm{M}}}\,F_{\mathrm{SM}}^{\mathrm{b}}.$$

$$F_{\mathrm{SM}}^{\mathrm{W}} = 6 + \frac{\overline{\mathrm{M}}_{\mathrm{H}}^{\,2}}{\overline{\mathrm{M}}^{\,2}} + 6\left(\overline{\mathrm{M}}_{\mathrm{H}}^{\,2} - 2\,\overline{\mathrm{M}}^{\,2}\right) C_0\left(-\overline{\mathrm{M}}_{\mathrm{H}}^{\,2}, 0, 0; \overline{\mathrm{M}}, \overline{\mathrm{M}}, \overline{\mathrm{M}}\right),$$

$$F_{\mathrm{SM}}^{\mathrm{t}} = -8 - 4\left(\overline{\mathrm{M}}_{\mathrm{H}}^{\,2} - 4\,M_{\mathrm{t}}^2\right) C_0\left(-\overline{\mathrm{M}}_{\mathrm{H}}^{\,2}, 0, 0; M_{\mathrm{t}}, M_{\mathrm{t}}, M_{\mathrm{t}}\right),$$

We only need a subset of operators $\curvearrowright$

$$
\begin{aligned}
\widetilde{\mathscr{L}} \;=\;& A^1_{\mathrm{V}}\left(\Phi^\dagger\Phi - v^2\right) F^a_{\mu\nu}\,F^a_{\mu\nu} + A^2_{\mathrm{V}}\left(\Phi^\dagger\Phi - v^2\right) F^0_{\mu\nu}\,F^0_{\mu\nu}\\
+\;& A^3_{\mathrm{V}}\Phi^\dagger\tau_a\Phi\,F^a_{\mu\nu}\,F^0_{\mu\nu} + \frac{1}{2}A_{\partial\Phi}\,\partial_\mu\left(\Phi^\dagger\Phi\right)\partial_\mu\left(\Phi^\dagger\Phi\right)\\
+\;& A^1_{\Phi}\left(\Phi^\dagger\Phi\right)\left(D_\mu\Phi\right)^\dagger D_\mu\Phi + A^3_{\Phi}\left(\Phi^\dagger D_\mu\Phi\right)\left[\left(D_\mu\Phi\right)^\dagger\Phi\right]\\
+\;& \frac{1}{4\sqrt{2}}\,\frac{M_{\mathrm{t}}}{\overline{\overline{\mathrm{M}}}}\,A^1_{\mathrm{f}}\left(\Phi^\dagger\Phi - v^2\right)\overline{\psi}_{\mathrm{L}}\,\Phi\,\mathrm{t}_{\mathrm{R}}\\
+\;& \frac{1}{4\sqrt{2}}\,\frac{M_{\mathrm{b}}}{\overline{\overline{\mathrm{M}}}}\,A^2_{\mathrm{f}}\left(\Phi^\dagger\Phi - v^2\right)\overline{\psi}_{\mathrm{L}}\,\Phi^c\,\mathrm{b}_{\mathrm{R}} + \text{h. c.}
\end{aligned}
$$

$$
A^0_{\Phi} = A^1_{\Phi} + 2\,\frac{A^3_{\Phi}}{\hat{s}^2_\theta} + 4\,A_{\partial\Phi}.
$$

✌

$$\mathscr{M}_{\mathrm{H}\to\gamma\gamma} = \left(4\sqrt{2}\,G_{\mathrm{F}}\right)^{1/2}\left\{-\frac{\alpha}{\pi}\left[C^{\gamma\gamma}_{\mathrm{W}}\,F^{\mathrm{W}}_{\mathrm{SM}} + 3\sum_{\mathrm{q}} Q^2_{\mathrm{q}}\,C^{\gamma\gamma}_{\mathrm{q}}\,F^{\mathrm{q}}_{\mathrm{SM}}\right] + F_{\mathrm{AC}}\right\}$$

$$F_{\mathrm{AC}} = \frac{g_6}{\sqrt{2}}\,\overline{\mathrm{M}}^{\,2}_{\mathrm{H}}\left(\hat{s}^2_\theta\,A^1_{\mathrm{V}} + \hat{c}^2_\theta\,A^2_{\mathrm{V}} + \hat{c}_\theta\,\hat{s}_\theta\,A^3_{\mathrm{V}}\right).$$

$$g_6 = \frac{1}{G_{\mathrm{F}}\,\Lambda^2} = 0.085736\left(\frac{TeV}{\Lambda}\right)^2$$

✌ the scaling factors are given by

$$\mathcal{C}^{\gamma\gamma}_{\mathrm{W}} = \frac{1}{4}\,\overline{\mathrm{M}}^{\,2}\Big\{1 + \frac{g_6}{4\sqrt{2}}\Big[8\,A^3_{\mathrm{V}}\,\hat{c}_\theta\,\Big(\hat{s}_\theta + \frac{1}{\hat{s}_\theta}\Big) + A^0_{\Phi}\Big]\Big\}$$

$$\mathcal{C}^{\gamma\gamma}_{\mathrm{t}} = \frac{1}{8}\,M^2_{\mathrm{t}}\,\Big\{1 + \frac{g_6}{4\sqrt{2}}\Big[8\,A^3_{\mathrm{V}}\,\hat{c}_\theta\,\Big(\hat{s}_\theta + \frac{1}{\hat{s}_\theta}\Big) + A^0_{\Phi} - A^1_{\mathrm{f}}\Big]\Big\}$$

$$\mathcal{C}^{\gamma\gamma}_{\mathrm{b}} = \frac{1}{8}\,M^2_{\mathrm{b}}\,\Big\{1 + \frac{g_6}{4\sqrt{2}}\Big[8\,A^3_{\mathrm{V}}\,\hat{c}_\theta\,\Big(\hat{s}_\theta + \frac{1}{\hat{s}_\theta}\Big) + A^0_{\Phi} - A^2_{\mathrm{f}}\Big]\Big\}$$

Glimpsing at the headlines of the **complete** calculation for
$\mathbf{H \to \gamma\gamma}$

☛

- **SM** loops, dressed with admissible operators
- **New 33** loop-diagrams
- **Counter-terms**

Amplitude in *internal* notations ☛

```
g HAA= -int(q)*Qs(-1,[q]^2+mt^2)*Qs(-1,[q+p1]^2+mt^2)*Qs(-1,[q+p1+p2]^2+mt^2)*3*trace*(
 ( -1/2*g*mt/M + L^-2 *( 4*r2^-1*M^2*af1 - 2*M*aV1*mt - 1/2*a3K*M*g*mt + 2*adK*M*g*mt))*
 (-i_*(gd(s,q)+gd(s,p1)+gd(s,p2))+mt)*
 VAtt(nu,p2)*(-i_*(gd(s,q)+gd(s,p1))+mt)*
 VAtt(mu,p1)*(-i_*gd(s,q)+mt)+
 ( -1/2*g*mt/M + L^-2 *( 4*r2^-1*M^2*af1 - 2*M*aV1*mt - 1/2*a3K*M*g*mt + 2*adK*M*g*mt))*
 ( i_*gd(s,q)+mt)*
 VAtt(mu,p1)*( i_*(gd(s,q)+gd(s,p1))+mt)*
 VAtt(nu,p2)*( i_*(gd(s,q)+gd(s,p1)+gd(s,p2))+mt))-
 int(q)*Qs(-1,[q]^2+mb^2)*Qs(-1,[q+p1]^2+mb^2)*Qs(-1,[q+p1+p2]^2+mb^2)*trace*(
 ( -1/2*g*mb/M + L^-2 *( - 4*r2^-1*M^2*af2 - 2*M*aV1*mb - 1/2*a3K*M*g*mb + 2*adK*M*g*mb))*
 (-i_*(gd(s,q)+gd(s,p1)+gd(s,p2))+mb)*
 VAbb(nu,p2)*(-i_*(gd(s,q)+gd(s,p1))+mb)*
 VAbb(mu,p1)*(-i_*gd(s,q)+mb)+
 ( -1/2*g*mb/M + L^-2 *( - 4*r2^-1*M^2*af2 - 2*M*aV1*mb - 1/2*a3K*M*g*mb + 2*adK*M*g*mb))*
 ( i_*gd(s,q)+mb)*
 VAbb(mu,p1)*( i_*(gd(s,q)+gd(s,p1))+mb)*
 VAbb(nu,p2)*( i_*(gd(s,q)+gd(s,p1)+gd(s,p2))+mb))+
 + i_*L^-2 *(
    + 8*M*(sth^2*aV1 + cth^2*aV2 + sth*cth*aV3)*(p1(nu)*p2(mu) - d_(mu,nu)*p1.p2))+
 int(q)*Qs(-1,[q]^2+M^2)*Qs(-1,[q+p1]^2+M^2)*Qs(-1,[q+p1+p2]^2+M^2)*(
 dia1*VHWW(al,be,-q,q+p1+p2)*VAWmWp(nu,be,si,p2,-q-p1-p2,q+p1)*VAWmWp(mu,si,al,p1,-q-p1,q)+
 dia2*VHWW(be,al,q+p1+p2,-q)*VAWmWp(mu,al,si,p1,q,-q-p1)*VAWmWp(nu,si,be,p2,q+p1,-q-p1-p2)+
 dia3*VHPmWp(al,-p1-p2,-q)*VAWmWp(nu,al,be,p2,-q-p1-p2,q+p1)*VAPpWm(mu,be,p1,-q-p1)+

 .................

 dia30*VAAWP(mu,nu,al,p1,p2)*VHPpWm(al,-p1-p2,-q))+
 int(q)*Qs(-1,[q]^2+M0^2)*Qs(-1,[q+p1+p2]^2+M0^2)*(
 dia31*VHP0P0(-p1-p2,-q,q+p1+p2)*VAAP0P0(mu,nu,p1,p2))+
 int(q)*Qs(-1,[q]^2+mh^2)*Qs(-1,[q+p1+p2]^2+mh^2)*(
 dia32*VHHH(-p1-p2,q+p1,-q)*VAAHH(mu,nu,p1,p2))+
 int(q)*Qs(-1,[q]^2+M^2)*(dia33*VHAAWW(mu,nu,si,si));
```

```
id VHPmPp(p1?,p2?,p3?)=
  - 1/2*M^-1*mh^2*g
  + L^-2 * (
  - 2*M*mh^2*aV1 - 2*p2.p3*a1K*M*g + 1/2*(mh^2 + 2*p1.p1)*a3K*M*g - 2*(mh^2 + 2*p1.p1)*adK*M*g
*
id VHPmWp(be?,p1?,p2?)=
  - 1/2*(p1(be) - p2(be))*i_*g
  + L^-2 * (
  - 2*p2(be)*i_*a1K*M^2*g - 2*(p1(be) - p2(be))*i_*M^2*aV1
  - 1/2*(p1(be) - p2(be))*i_*a3K*M^2*g + 2*(p1(be) - p2(be))*i_*adK*M^2*g);
*
id VHPpWm(be?,p1?,p2?)=
  - 1/2*(p1(be) - p2(be))*i_*g
  + L^-2 * (
  - 2*p2(be)*i_*a1K*M^2*g - 2*(p1(be) - p2(be))*i_*M^2*aV1
  - 1/2*(p1(be) - p2(be))*i_*a3K*M^2*g + 2*(p1(be) - p2(be))*i_*adK*M^2*g);
*
id VHWW(al?,be?,p2?,p3?)=
  - d_(al,be)*M*g
  + L^-2 * (
  - 4*d_(al,be)*M^3*aV1 - d_(al,be)*a3K*M^3*g + 2*d_(al,be)*a1K*M^3*g
  + 4*d_(al,be)*adK*M^3*g + 8*(p2(be)*p3(al) - d_(al,be)*p2.p3)*M*aV1);
*
id VHZZ(al?,be?,p2?,p3?)=
  - d_(al,be)*M*cth^-2*g
  + L^-2 * (
  - 4*d_(al,be)*M^3*aV1*cth^-2 + d_(al,be)*a3K*M^3*cth^-2*g
  + 2*d_(al,be)*a1K*M^3*cth^-2*g + 4*d_(al,be)*adK*M^3*cth^-2*g
  - 8*(p2(be)*p3(al) - d_(al,be)*p2.p3)*M*aV3*cth*sth
  + 8*(p2(be)*p3(al) - d_(al,be)*p2.p3)*M*aV2*sth^2
  + 8*(p2(be)*p3(al) - d_(al,be)*p2.p3)*M*aV1*cth^2);
*
```

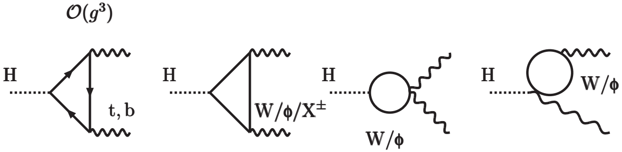
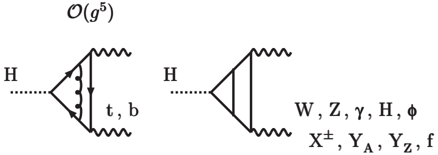


Figure 1: The three families of diagrams contributing to the amplitude for H $\to$ $\gamma\gamma$; W/ϕ denotes a W-line or a ϕ-line. $X^{\pm}$ denotes a FP-ghost line



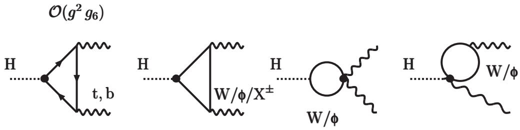


Figure 3: Example of one-loop SM diagrams with $\mathcal{O}$-insertions, contributing to the amplitude for $\mathrm{H} \to \gamma\gamma$

- • denotes operator insertion

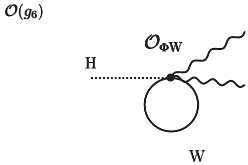


Figure 4: Example of one-loop $\mathcal{O}$-diagrams, contributing to the amplitude for $\mathrm{H} \to \gamma\gamma$

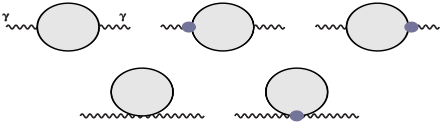


Figure 5: The photon self-energy with inclusion of $\mathcal{O}$-operators into SM one-loop diagrams. The last diagram contains vertices, like $\mathrm{AAHH}, \mathrm{AA}\phi^0\phi^0$, that do not belong to the SM part.

Renormalization

$$
\begin{aligned}
g &= g_{\text{ren}} \left[1 + \frac{g_{\text{ren}}^2}{16\,\pi^2} \left(dZ_g + g_6\, dZ_g^{(6)} \right) \frac{1}{\varepsilon} \right] \\
M_{\mathrm{W}} &= M_{\mathrm{W}}^{\text{ren}} \left[1 + \frac{1}{2}\, \frac{g_{\text{ren}}^2}{16\,\pi^2} \left(dZ_{M_{\mathrm{W}}} + g_6\, dZ_{M_{\mathrm{W}}}^{(6)} \right) \frac{1}{\varepsilon} \right]
\end{aligned}
$$

etc.

Wilson coefficients $\quad\rightarrow\quad W_i$

$$
\begin{aligned}
W_i &= \sum_j Z_{ij}^{\mathrm{wc}}\, W_j^{\text{ren}} \\
Z_{ij}^{\mathrm{wc}} &= \delta_{ij} + \left(g_{\text{ren}}^2\, dZ_{ij}^{\mathrm{wc}} + dZ_{ij}^{\mathrm{wc},6} \right) \frac{1}{\varepsilon}
\end{aligned}
$$

Conclusions?

Data-driven Theory? Or

If you're looking for your lost keys, failing to find them in the kitchen is not evidence against their being somewhere else in the house

Conclusions?

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Conclusions?

- Higgs-landscape: asking the right questions takes as much skill as giving the right answers
- *We are afraid of ideas, of experimenting, of change. We shrink from thinking a problem through to a logical conclusion.* (Anne Sullivan Macy)
- *I am turned into a sort of machine for observing facts and grinding out conclusions* (Charles Darwin)
- *El sueño de la razõn produce monstruos* (Francisco Goya)

Thanks for your attention

Backup

assumptions/inferences

- Given the (few) known coefficients in the perturbative expansion we estimate the next (few) coefficients and the corresponding partial sums by means of sequence transformations. This is the first step towards “ reconstructing ” the physical observable.
- The sequence transformations have been tested on a number of test sequences.
- A function can be uniquely determined by its asymptotic expansion if certain conditions are satisfied (Sokal).
- Borel procedure is a summation method which, under the above conditions, determines uniquely the sum of the series. It should be taken into account that there is a large class of series that have Borel sums (analytic in the cut-plane) and there is evidence that Levin-Weniger transforms produce approximations to these Borel sums. This is one of the arguments of plausibility supporting our results.
- The QCD scale variation uncertainty decreases when we include new (estimated) partial sums.
- All known and predicted coefficients are positive and all transforms predict convergence within a narrow interval.
- Missing a formal proof of uniqueness, we assume uninformative prior between the last known partial sum and the (largest) predicted partial sum.

F. Wilczek hep-ph/9311302

typical strong interaction scale we'll be getting higher and higher powers of the strong interactions scale over Q^2. Keeping the first few terms should be a good approximation even at 1.8 GeV. It is very helpful that the mass dimensions of the gauge invariant operators start at 4.

The Wilson coefficients, the operator product coefficients $\mathcal{C}$ above, obey renormalization group equations. They can be calculated in perturbation theory in the effective coupling at large Q^2, of course. However, at Q^2 of approximately m_τ^2 we cannot simply ignore plausible non-perturbative corrections and still guarantee worthwhile accuracy. A term of the form $\Lambda^2_{\text{QCD}}/Q^2$ would show up, through the mechanism of dimensional transmutation, as a contribution proportional to $\exp\left(-c/\alpha_s\right)$ in this coefficient, where c is a calculable numerical constant. It is an important question whether there is such a contribution, because if there were, and they were not under tight control, it is formally of such a magnitude as to ruin the useful precision of the predictions. Such a correction would be bigger than the ones coming from higher operators because these operators have dimension 4, so their coefficients have Q^2 over Λ^2 squared, which is *a priori* smaller.

Mueller [7] has given an important, although not entirely rigorous, argument that no Λ^2/Q^2 term can appear. The argument is a little technical, so I won't be able to do it full justice here but I will attempt to convey the main idea. The argument is based on the idea that at each successive power of 1 over Q^2 one can make the perturbation series in QCD, which is a badly divergent series in general, at least almost convergent, that is Borel summable, by removing a finite number of obstructions. Furthermore the obstructions are captured and parameterized by the low dimension operators mentioned before. Once these obstructions are removed, the remaining (processed) perturbation expression converges on the correct result for the full theory. Neither in the obstructions nor in the residual perturbative expression do the potentially dangerous terms occur – which means that they don't occur at all.

Maybe I should draw a picture of this [Figure 4]. One has the current product, and one is doing an analysis of its behavior when large virtual momentum is flowing through the current lines. The principle of the operator product expansion is to exhibit the powers of Q^2 by breaking the propagators in the graph into hard and soft parts. Any soft part costs you a power of one over Q^2 so you want the minimal number. If you just take out a couple of lines you have one of those low dimension operators, so those are interpreted as the operators, with the causality and unitarity. The usual demonstrations that these properties hold order by order in perturbation theory can be adapted to the re-processed version, which is more complicated but has the virtue of actually defining an answer. In fact we can agree that it gives *the* answer, since after all the whole point of quantum field theory is to give non-trivial realizations of the axioms, and that is what we have found.

QCD is not quite so favorable as this ideal, which occurs only for massive super-renormalizable theories in low dimensions. There are several known obstructions to Borel summability in QCD, which go by frightening names: ultraviolet and infrared renormalons, instantons, and threshold-induced oscillations. What Mueller did was to analyze these known sources of possible dangerous terms. He argued that the infrared renormalons are essentially just the higher-order terms in the operator product expansion, the ultraviolet renormalons generate singularities in $g(\alpha)$ away from the real axis whose influence on the truncated form of $g(\alpha)$ one actually computes can be minimized by judicious mappings in the α plane, that the threshold-induced oscillations are negligible quantitatively, and that the instanton contribution is both small and in principle calculable.

So now I have fleshed out my earlier description of Mueller's argument a bit. The key underlying assumption is that the known obstacles to Borel summability are the only ones. In principle, one can test this circle of ideas by calculating the operator product coefficients directly in the full theory (*i.e.* numerically, using lattice gauge theory techniques). If they were to fail, it would signify that there is an important gap in our understanding of quantum field theory.

On the experimental side, the Aleph group has tested the framework leading to this operator product expansion by comparing the resulting specific predictions for decay into semi-inclusive final states with specific quantum numbers, including the Q^2 dependence (which you can look at by looking at final states of different mass) [9]. They got a good fit with no one over Q^2 term and with matrix elements of the lowest dimension relevant operators $m\bar{\psi}\psi$, $\mathrm{tr}G_{\mu\nu}G_{\mu\nu}$ fitted to other experiments. These quantities also appear in other similar applications, where observed hadron parameters are correlated using the so-called QCD or ITEP sum rules, which arise by saturated various operator products. By taking suitable moments one can define quantities that are insensitive to the higher dimension operators, and for these the predictions of perturbative QCD are especially stringent.

I went into some detail into the analysis of tau decay because I think it's not

Structure of the calculation

- Process: $\mathrm{H} \to \bar{\mathrm{f}}\mathrm{f}\gamma$, $\mathrm{f} = \mathrm{l}, \mathrm{q}$,
  including b with non-zero m_t
- Setup: $m_\mathrm{f} = 0$ at NLO. Calculation based on helicity amplitudes
  LO and NLO do not interfere (with $m_\mathrm{f} = 0$)

Cuts available in the H rest-frame

Please complain but it took years to interface *POWHEG* and *Prophecy4f* ......

$\mathrm{gg} \to \bar{\mathrm{f}}\mathrm{f}\gamma$? Can be done, *But* ......

HTO-DALITZ Features

- Internal cross-check, loops are evaluated both *analytically* and *numerically* (using BST-algorithm)
- The code makes extensive use of *In-House* abbreviation algorithms (if $a+b$ appears twice or more it receives an abbreviation and it is pre-computed only once).
- All functions are **collinear-free**
- High performances thanks to gcc-4.8.0
- *Open MPI* version under construction, *GPU* version in a preliminar phase
- Returns the full result and also the unphysical components

Man at work

- Extensions: as it was done during Lep times, there are diagrams where both the Z and the γ propagators should be Dyson-improved, i.e.

$$\alpha_{\mathrm{QED}}(0) \to \alpha_{\mathrm{QED}}(\mathrm{virtuality}) \qquad \rho_{\mathrm{f}} - \text{parameter included}$$

- However, the interested sub-sets are not gauge invariant, $\therefore$ appropriate subtractions must be performed (at virtuality $= 0$, s_{Z}, the latter being the Z complex-pole).

Misunderstandings

- use $M\left(\bar{\mathrm{f}}\mathrm{f}\gamma\right)$ and require $\mid M - M_{\mathrm{Z}} \mid \leq n\Gamma_{\mathrm{Z}}$. *This is not the photon we are discussing*
  Photons are collinear to leptons only if emitted by leptons but those are Yukawa-suppressed.
  In any case $M\left(\bar{\mathrm{f}}\mathrm{f}\gamma\right) = M_{\mathrm{H}}$ or it is *Not* Dalitz decay
- Requiring a cut on the opening angle between leptons and the photon to define *isolated photons* is highly recommended, *But* at the moment we are still in the Higgs rest-frame (*Miracles take a bit longer*)

Return

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Return

with the physical mass parameters

$$
\begin{aligned}
M_{\mathrm{W}}^2 &= \frac{1}{4}g^2v^2\left[1+2\frac{v^2}{\Lambda^2}\alpha_{\Phi\mathrm{W}}\right],\\
M_{\mathrm{Z}}^2 &= \frac{1}{4}(g^2+g'^2)v^2\left[1+\frac{v^2}{2\Lambda^2}\left(4\alpha_{\mathrm{ZZ}}+\alpha_{\Phi D}\right)\right],\\
M_{\mathrm{H}}^2 &= \lambda v^2\left[1+\frac{v^2}{2\Lambda^2}\left(4\alpha_{\Phi\Box}-\frac{6}{\lambda}\alpha_{\Phi}-\alpha_{\Phi D}\right)\right],\\
m_{\mathrm{f}} &= \frac{1}{\sqrt{2}}U^{\mathrm{f}}\Gamma_{\mathrm{f}}U^{\mathrm{f},\dagger}v\left[1-\frac{1}{2}\frac{v^2}{\Lambda^2}\alpha_{\mathrm{f}\phi}\right].
\end{aligned}
\qquad (151)
$$

In (150) we have used the usual 't Hooft–Feynman gauge-fixing term

$$
\mathcal{L}_{\mathrm{fix}} = -C_+C_- - \frac{1}{2}(C_{\mathrm{Z}})^2 - \frac{1}{2}(C_{\mathrm{A}})^2 - \frac{1}{2}C_G^AC_G^A \qquad (152)
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with

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C_G^A = \partial_\mu G^{A\mu}, \qquad C_{\mathrm{A}} = \partial_\mu \mathrm{A}^\mu, \qquad C_{\mathrm{Z}} = \partial_\mu \mathrm{Z}^\mu + M_{\mathrm{Z}}\phi^0, \qquad C_\pm = \partial_\mu \mathrm{W}^{\pm\mu} \pm \mathrm{i}M_{\mathrm{W}}\phi^\pm \qquad (153)
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in terms of the physical fields and parameters, which gives rise to the same propagators as in the SM.

In the following, the abbreviations c_{w} and s_{w} are defined via the physical masses

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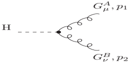
10.4.2 Higgs vertices

Here we list the most important Feynman rules for vertices involving exactly one physical Higgs boson. These are given in terms of the above-defined physical fields and parameters. In the coefficients of dimension-6 couplings we replaced v^2 by the Fermi constant via $v^2 = 1/(\sqrt{2}G_F)$.

The triple vertices involving one Higgs boson read:

Hgg coupling:



$$
= \mathrm{i}\frac{2g}{M_{\mathrm{W}}}\frac{1}{\sqrt{2}G_F\Lambda^2}\left[\alpha_{GG}(p_{2\mu}p_{1\nu}-p_1p_2g_{\mu\nu})+\alpha_{G\tilde{G}}\varepsilon_{\mu\nu\rho\sigma}p_1^\rho p_2^\sigma\right]\delta^{AB}, \qquad (155)
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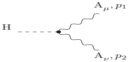

$$
= \mathrm{i}\frac{2g}{M_{\mathrm{W}}}\frac{1}{\sqrt{2}G_F\Lambda^2}\left[\alpha_{\mathrm{AA}}(p_{2\mu}p_{1\nu}-p_1p_2g_{\mu\nu})+\alpha_{\mathrm{A}\tilde{\mathrm{A}}}\varepsilon_{\mu\nu\rho\sigma}p_1^\rho p_2^\sigma\right], \qquad (156)
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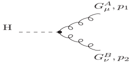

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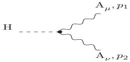

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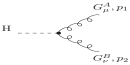

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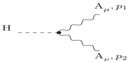

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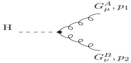

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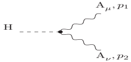

$$
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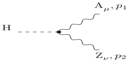
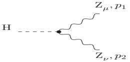
$$

HAZ coupling:



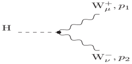
$$= \mathrm{i}\frac{g}{M_\mathrm{W}}\frac{1}{\sqrt{2}G_F\Lambda^2}\left[\alpha_{\mathrm{AZ}}(p_{2\mu}p_{1\nu} - p_1p_2g_{\mu\nu}) + \alpha_{\mathrm{A}\tilde{\mathrm{Z}}}\varepsilon_{\mu\nu\rho\sigma}p_1^\rho p_2^\sigma\right], \tag{157}$$

HZZ coupling:



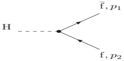
$$\begin{aligned} &= \mathrm{i}g\frac{M_\mathrm{Z}}{c_\mathrm{w}}g_{\mu\nu}\left[1 + \frac{1}{\sqrt{2}G_F\Lambda^2}\left(\alpha_{\Phi\mathrm{W}} + \alpha_{\Phi\Box} + \frac{1}{4}\alpha_{\Phi D}\right)\right] \\ &\quad + \mathrm{i}\frac{2g}{M_\mathrm{W}}\frac{1}{\sqrt{2}G_F\Lambda^2}\left[\alpha_{\mathrm{ZZ}}(p_{2\mu}p_{1\nu} - p_1p_2g_{\mu\nu}) + \alpha_{\mathrm{Z}\tilde{\mathrm{Z}}}\varepsilon_{\mu\nu\rho\sigma}p_1^\rho p_2^\sigma\right], \end{aligned} \tag{158}$$

HWW coupling:



$$\begin{aligned} &= \mathrm{i}gM_\mathrm{W}g_{\mu\nu}\left[1 + \frac{1}{\sqrt{2}G_F\Lambda^2}\left(\alpha_{\Phi\mathrm{W}} + \alpha_{\Phi\Box} - \frac{1}{4}\alpha_{\Phi D}\right)\right] \\ &\quad + \mathrm{i}\frac{2g}{M_\mathrm{W}}\frac{1}{\sqrt{2}G_F\Lambda^2}\left[\alpha_{\Phi\mathrm{W}}(p_{2\mu}p_{1\nu} - p_1p_2g_{\mu\nu}) + \alpha_{\Phi\widetilde{\mathrm{W}}}\varepsilon_{\mu\nu\rho\sigma}p_1^\rho p_2^\sigma\right], \end{aligned} \tag{159}$$

Hff coupling:

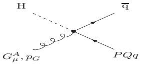


$$= -\mathrm{i}\frac{g}{2}\frac{m_\mathrm{f}}{M_\mathrm{W}}\left[1 + \frac{1}{\sqrt{2}G_F\Lambda^2}\left(\alpha_{\Phi\mathrm{W}} + \alpha_{\Phi\Box} - \frac{1}{4}\alpha_{\Phi D} - \alpha_{\mathrm{f}\phi}\right)\right], \tag{160}$$

where $\mathrm{f} = \mathrm{e}, \mathrm{u}, \mathrm{d}$.

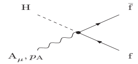
The quadruple vertices involving one Higgs field, one gauge boson and a fermion–antifermion pair are given by ($\mathrm{q} = \mathrm{u}, \mathrm{d}$, $\mathrm{f} = \mathrm{u}, \mathrm{d}, \nu_\mathrm{l}, \mathrm{e}$, and $\hat{\mathrm{f}} = \mathrm{q}$ for $\mathrm{f} = \mathrm{u}, \mathrm{d}$ and $\hat{\mathrm{f}} = \mathrm{l}$ for $\mathrm{f} = \mathrm{e}$):

$\mathrm{Hgq}\overline{\mathrm{q}}$ coupling:



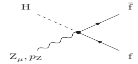
$$= \mathrm{i}g\frac{m_q}{M_\mathrm{W}}\mathrm{i}p_{G\nu}\sigma^{\mu\nu}\frac{\lambda^A}{2}\left[\frac{1+\gamma_5}{2}\alpha_{\mathrm{q}G} + \frac{1-\gamma_5}{2}\alpha^*_{\mathrm{q}G}\right], \tag{161}$$

HAf$\bar{\text{f}}$ coupling:



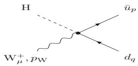
$$= \mathrm{i}g\frac{m_f}{M_\text{W}}\mathrm{i}p_{\text{A}\nu}\sigma^{\mu\nu}\left[\frac{1+\gamma_5}{2}(\alpha_{\text{fB}}c_\text{w} + 2I_3^f\alpha_{\text{fW}}s_\text{w}) + \frac{1-\gamma_5}{2}(\alpha^*_{\text{fB}}c_\text{w} + 2I_3^f\alpha^*_{\text{fW}}s_\text{w})\right], \quad (162)$$

HZf$\bar{\text{f}}$ coupling:



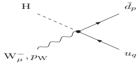
$$\begin{aligned} &= \mathrm{i}g\frac{m_f}{M_\text{W}}\mathrm{i}p_{\text{Z}\nu}\sigma^{\mu\nu}\left[\frac{1+\gamma_5}{2}(2I_3^\text{f}\alpha_{\text{fW}}c_\text{w} - \alpha_{\text{fB}}s_\text{w}) + \frac{1-\gamma_5}{2}(2I_3^\text{f}\alpha^*_{\text{fW}}c_\text{w} - \alpha^*_{\text{fB}}s_\text{w})\right] \\ &\quad + \mathrm{i}2M_\text{Z}\gamma^\mu\left[\frac{1-\gamma_5}{2}\left(\alpha^{(1)}_{\Phi\hat{\text{f}}} - 2I_3^\text{f}\alpha^{(3)}_{\Phi\hat{\text{f}}}\right) + \frac{1+\gamma_5}{2}\alpha_{\Phi\text{f}}\right], \end{aligned} \quad (163)$$

HW$^+$d$\overline{\text{u}}$ coupling:



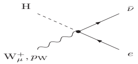
$$\begin{aligned} &= \mathrm{i}g\frac{\sqrt{2}}{M_\text{W}}\mathrm{i}p_{\text{W}\nu}\sigma^{\mu\nu}V_{pq}\left[\frac{1+\gamma_5}{2}m_\text{d}\alpha_{\text{dW}} + \frac{1-\gamma_5}{2}m_\text{u}\alpha^*_{\text{uW}}\right] \\ &\quad - \mathrm{i}\sqrt{2}M_\text{W}\gamma^\mu\left[\frac{1-\gamma_5}{2}\,2\alpha^{(3)}_{\Phi\text{q}}V_{pq} + \frac{1+\gamma_5}{2}(\Gamma_{\text{ud}})_{pq}\alpha_{\Phi\text{ud}}\right], \end{aligned} \quad (164)$$

HW$^-$u$\overline{\text{d}}$ coupling:



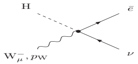
$$\begin{aligned} &= \mathrm{i}g\frac{\sqrt{2}}{M_\text{W}}\mathrm{i}p_{\text{W}\nu}\sigma^{\mu\nu}V^\dagger_{pq}\left[\frac{1+\gamma_5}{2}m_\text{u}\alpha_{\text{uW}} + \frac{1-\gamma_5}{2}m_\text{d}\alpha^*_{\text{dW}}\right] \\ &\quad - \mathrm{i}\sqrt{2}M_\text{W}\gamma^\mu\left[\frac{1-\gamma_5}{2}\,2\alpha^{(3)}_{\Phi\text{q}}V^\dagger_{pq} + \frac{1+\gamma_5}{2}(\Gamma^\dagger_{\text{ud}})_{pq}\alpha^*_{\Phi\text{ud}}\right], \end{aligned} \quad (165)$$

HW$^+$e$\overline{\nu}_\text{e}$ coupling:



$$= \mathrm{i}g\frac{\sqrt{2}}{M_\text{W}}\mathrm{i}p_{\text{W}\nu}\sigma^{\mu\nu}\frac{1+\gamma_5}{2}m_\text{e}\alpha_{\text{eW}} - \mathrm{i}\sqrt{2}M_\text{W}\gamma^\mu\frac{1-\gamma_5}{2}\,2\alpha^{(3)}_{\Phi\text{l}}, \quad (166)$$

HW$^-\nu_\text{e}$e$^+$ coupling:



$$= \mathrm{i}g\frac{\sqrt{2}}{M_\text{W}}\mathrm{i}p_{\text{W}\nu}\sigma^{\mu\nu}\frac{1-\gamma_5}{2}m_\text{e}\alpha^*_{\text{eW}} - \mathrm{i}\sqrt{2}M_\text{W}\gamma^\mu\frac{1-\gamma_5}{2}\,2\alpha^{(3)}_{\Phi l}. \quad (167)$$

Decoupling and $SU(2)_C$

- Heavy degrees of freedom $\hookrightarrow \mathbf{H} \to \gamma\gamma$: to be fully general one has to consider effects due to heavy fermions $\in \boldsymbol{R_f}$ and heavy scalars $\in \boldsymbol{R_s}$ of $\boldsymbol{SU(3)}$. Colored scalars disappear from the low energy physics as their mass increases . However, the same is not true for fermions.

Decoupling and $SU(2)_C$

- Heavy degrees of freedom $\hookrightarrow \mathbf{H} \to \gamma\gamma$: to be fully general one has to consider effects due to heavy fermions $\in \boldsymbol{R_f}$ and heavy scalars $\in \boldsymbol{R_s}$ of $\boldsymbol{SU(3)}$. Colored scalars disappear from the low energy physics as their mass increases . However, the same is not true for fermions.

- Renormalization: whenever $\rho_{\mathrm{LO}} \neq 1$, quadratic power-like contribution to $\Delta\rho$ are absorbed by renormalization of the new parameters of the model $\rightsquigarrow \rho$ is not a measure of the custodial symmetry breaking.
  Alternatively one could examine models containing $\boldsymbol{SU(2)_L \otimes SU(2)_R}$ multiplets.