

# **Axions and lattice QCD**

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# Axion

A hypothetical elementary particle introduced to solve a puzzle with the parity transformation in particle physics and a leading candidate for the dark matter particle .

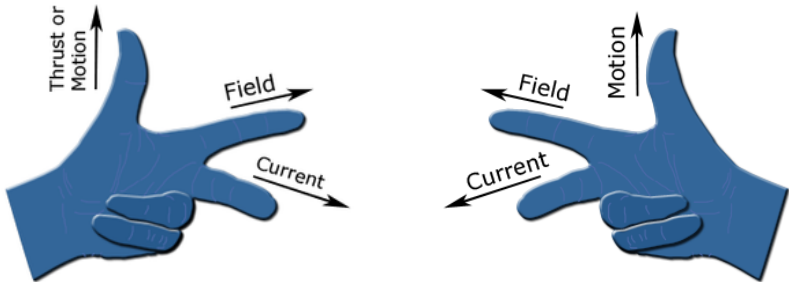
# Parity transformation (P)

Flip the sign of all spatial coordinates.

$$x \rightarrow -x, \quad y \rightarrow -y, \quad z \rightarrow -z$$

Parity transforms an object to its mirror image

$$L \rightarrow R, \quad R \rightarrow L$$



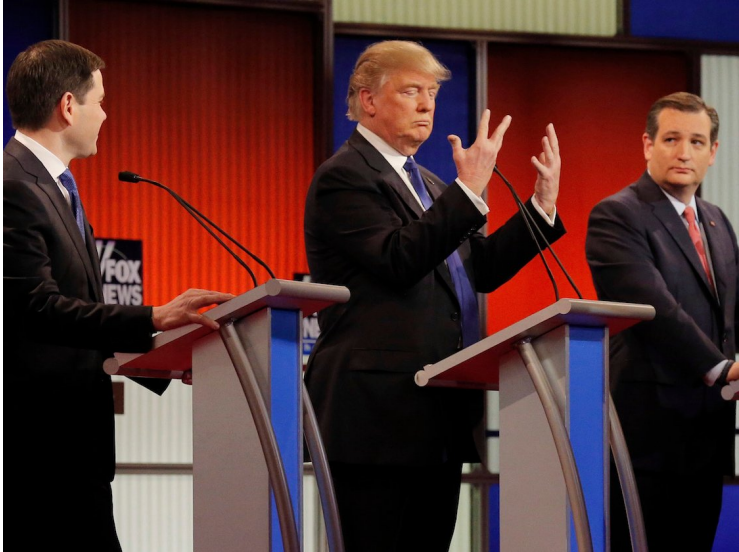
# Is P a symmetry?



Are the laws of physics the same in the mirror world?

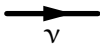
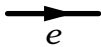
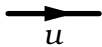
Or can I tell the difference between the original and the mirror image?

# Is P a symmetry?

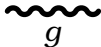


# Particle physics 101

**Particles:** up quark, down quark, electron and neutrino.

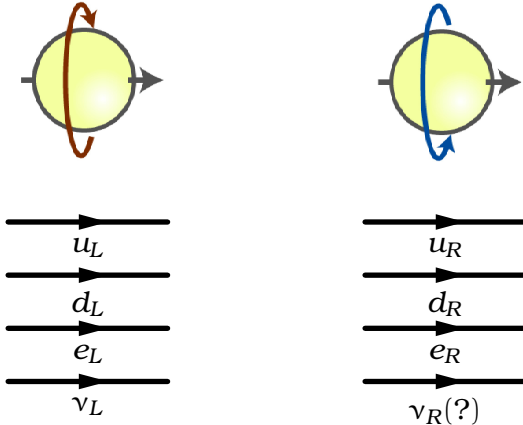


**Interactions:** strong, electromagnetic, weak and Higgs.



# Parity and particles

Particles have spin and are massless  $\rightarrow$



P exchanges **left-handed** particles with **right-handed**

# Electromagnetic interaction

$$\nabla \mathbf{E} = 4\pi\rho, \quad \nabla \mathbf{B} = 0, \quad \nabla \times \mathbf{E} + \partial_t \mathbf{B} = 0, \quad \nabla \times \mathbf{B} - \partial_t \mathbf{E} = 4\pi\mathbf{j}$$

$$P: \mathbf{E} \rightarrow -\mathbf{E}, \mathbf{B} \rightarrow \mathbf{B}, \rho \rightarrow \rho, \mathbf{j} \rightarrow -\mathbf{j}, \nabla \rightarrow -\nabla$$

Classical electrodynamics

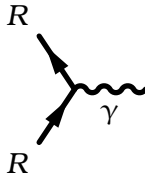
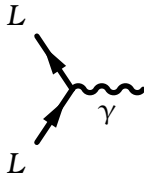


P symmetric



Quantum electrodynamics

$$L_{QED} = -\frac{1}{4}F_{\mu\nu}^2 + \psi_L^\dagger \sigma_\mu D_\mu \psi_L + \psi_R^\dagger \bar{\sigma}_\mu D_\mu \psi_R$$



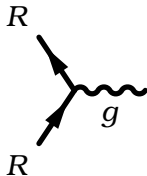
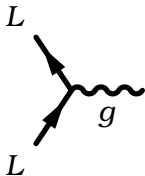


# Strong interaction

Generalized QED: 3x3-matrices and 3-vectors instead of numbers

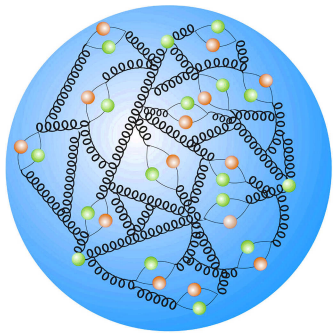
$$F_{\mu\nu} \rightarrow \begin{pmatrix} G_{\mu\nu}^{00} & G_{\mu\nu}^{01} & G_{\mu\nu}^{02} \\ G_{\mu\nu}^{10} & G_{\mu\nu}^{11} & G_{\mu\nu}^{12} \\ G_{\mu\nu}^{20} & G_{\mu\nu}^{21} & G_{\mu\nu}^{22} \end{pmatrix} \quad \psi \rightarrow \begin{pmatrix} \psi^0 \\ \psi^1 \\ \psi^2 \end{pmatrix}$$

$$L_{QCD} = -\frac{1}{4}\text{Tr}(G_{\mu\nu}^2) + (\psi_L^\dagger, \sigma_\mu D_\mu \psi_L) + (\psi_R^\dagger, \bar{\sigma}_\mu D_\mu \psi_R)$$



P symmetric

# Neutron electric dipole moment



P is symmetry  $[H, P] = 0$



neutron is P-eigenstate  $P|n\rangle \propto |n\rangle$



$$\langle n|\vec{d}|n\rangle = 0$$



EDM is P-odd  $P \cdot \vec{d} \cdot P = -\vec{d}$

Expt:  $\langle n|\vec{d}|n\rangle = -0.2(1.9) \times 10^{-26}$  ecm [Pendlebury '15]

# Higgs interaction

Higgs mechanism

# Higgs interaction

figs/bush.png

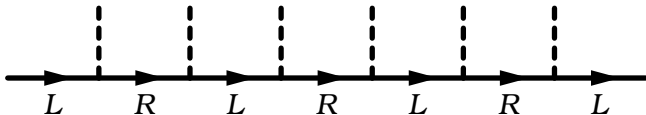
# Higgs interaction

figs/bush.png

“Left hand knows what the right hand is doing.”

# Higgs interaction

“Left hand knows what the right hand is doing.”



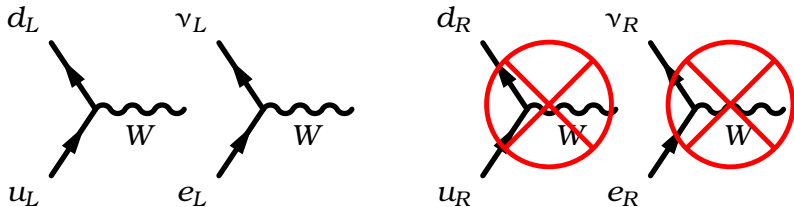
$$L_{Higgs} = m\psi_L^\dagger\psi_R + m\psi_R^\dagger\psi_L$$

Generate mass by combining massless L,R particles

P symmetric

# Weak interaction

Maximally violates parity, only interacts with L-particles.



P violating

Left:= the handedness of particles to which  $W$  couples

# The P puzzle

electromagnetic	✓
strong	✓
Higgs	✓
weak	✗

P is violated by the weak interaction  $\rightarrow$  it is not symmetry of Nature.

Why P is not violated by the others?

**Try:** P-invariance is consequence of remaining symmetries (Lorentz invariance, internal).



# The P puzzle in QED

$$L_{QED} = -\frac{1}{4}F_{\mu\nu}^2 + \psi_L^\dagger \sigma_\mu D_\mu \psi_L + \psi_R^\dagger \bar{\sigma}_\mu D_\mu \psi_R$$

What is the most general Lagrangian with Lorentz invariance and gauge symmetry?

$$L = L_{QED} + \theta \cdot F\tilde{F} \quad \text{with} \quad F\tilde{F} \equiv F_{\mu\nu}F_{\rho\sigma}\epsilon_{\mu\nu\rho\sigma}$$

Violates parity

$$F\tilde{F} \rightarrow -F\tilde{F}$$

Total derivative

$$F\tilde{F} = \partial_\mu K_\mu \quad \text{with} \quad K_\mu = \epsilon_{\mu\nu\rho\sigma} A_\nu F_{\rho\sigma}$$

then by Gauss-theorem

$$\int d^4x F\tilde{F} = \oint dn_\mu K_\mu = 0$$

P-invariance follows from Lorentz+gauge

# The P puzzle in QCD

Most general SU(3) symmetric Lagrangian

$$L = L_{QCD} + \theta \cdot G\tilde{G} \quad \text{with} \quad G\tilde{G} = \frac{1}{8\pi^2} \text{Tr} (G_{\mu\nu} G_{\rho\sigma} \epsilon_{\mu\nu\rho\sigma})$$

Violates parity

$$G\tilde{G} \rightarrow -G\tilde{G}$$

Total derivative

$$G\tilde{G} = \partial_\mu K_\mu \quad \text{with} \quad K_\mu = \text{Tr} \epsilon_{\mu\nu\rho\sigma} (A_\nu G_{\rho\sigma} + \frac{2}{3} A_\nu A_\rho A_\sigma)$$

then by Gauss-theorem

$$\int d^4x G\tilde{G} = \oint dn_\mu K_\mu \neq 0$$

$K_\mu$  can be non-zero at  $\infty$

P could be violated by  $\theta \cdot G\tilde{G}$ . Why not?

nEDM experiments  $\rightarrow \theta \lesssim 0.0000000001$

## P-violation in Higgs

$$L_{\text{Higgs}} = m \left( \psi_L^\dagger \psi_R + \psi_R^\dagger \psi_L \right) \quad \text{with real } m$$

Most general Lagrangian has **complex mass**:

$$L = m \psi_L^\dagger \psi_R + m^* \psi_R^\dagger \psi_L$$

It violates parity, but can be transformed away by an **axial transformation**:

$$\psi_L \rightarrow \psi_L, \quad \psi_R \rightarrow e^{-i \arg m} \psi_R$$

but P-violation does not go away:

$$L \rightarrow L + \arg m \cdot F\tilde{F} + \arg m \cdot G\tilde{G}$$

**Higgs P violation can be transformed to strong  $G\tilde{G}$**

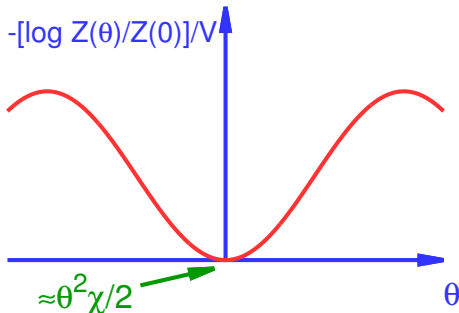
# The P puzzle

electromagnetic	✓	← by Lorentz+gauge invariance
strong	✓	} why $G\tilde{G}$ not appears in Nature?
Higgs	✓	
weak	✗	

## $\theta$ dependence of QCD

Calculate the Feynman path integral!

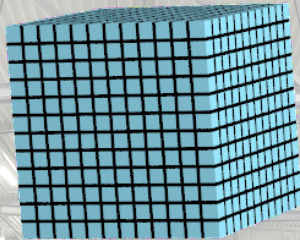
$$Z(\theta) = \int [dG][d\psi^\dagger][d\psi] \exp \left( i \int d^4x (L + \theta G\tilde{G}) \right)$$



Has a minimum at  $\theta = 0$ !

# Lattice QCD computation

Discretize!



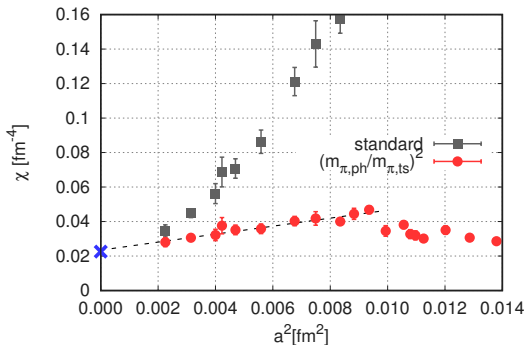
lattice spacing  $a \lesssim \frac{1}{10}$  proton size

Calculate!

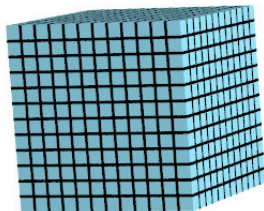
$10^9$  dimensional integrals using supercomputers

# $\theta$ dependence from lattice QCD

$$-\frac{1}{V} \log Z(\theta)/Z(0) = \frac{1}{2} \theta^2 \chi + \dots$$



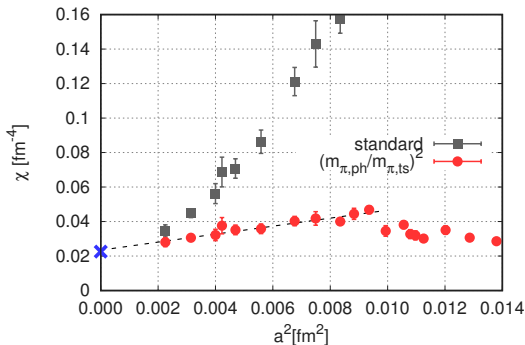
$a=0.134$  fm  
2.000 PCyears



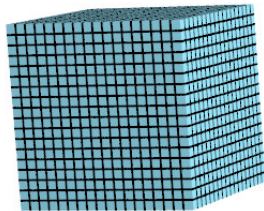
100.000 years on a PC, 1 year on a supercomputer.

# $\theta$ dependence from lattice QCD

$$-\frac{1}{V} \log Z(\theta)/Z(0) = \frac{1}{2} \theta^2 \chi + \dots$$



$a=0.095$  fm  
12.000 PCyears

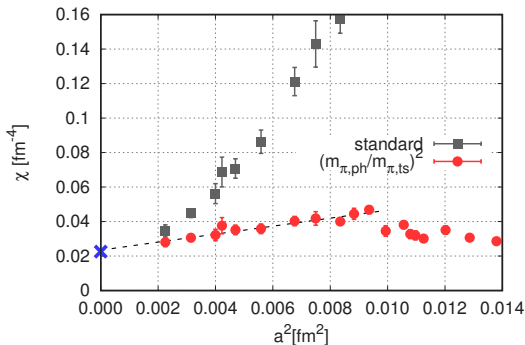


100.000 years on a PC, 1 year on a supercomputer.

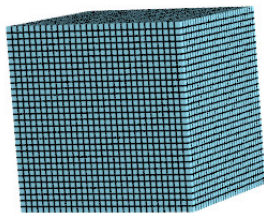


# $\theta$ dependence from lattice QCD

$$-\frac{1}{V} \log Z(\theta)/Z(0) = \frac{1}{2} \theta^2 \chi + \dots$$



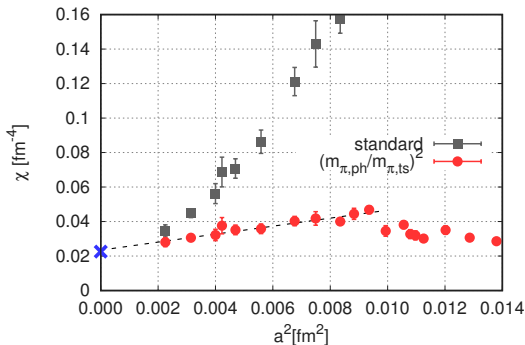
$a=0.064$  fm  
86.000 PCyears



100.000 years on a PC, 1 year on a supercomputer.

# $\theta$ dependence from lattice QCD

$$-\frac{1}{V} \log Z(\theta)/Z(0) = \frac{1}{2} \theta^2 \chi + \dots$$



Continuum limit

$$\chi^{1/4} = 76(2)(1) \text{MeV}$$

100.000 years on a PC, 1 year on a supercomputer.

# A solution by Peccei-Quinn

Make a dynamical field from  
the parameter!

figs/thetapot/plot.gif

$$L + \theta \cdot G\tilde{G} + \frac{1}{2}f^2 \cdot (\partial_\mu\theta)^2 + V(\theta, \partial_\mu\theta)$$

with  $V(\theta, \partial_\mu\theta)$  such, that minimum stays at  $\theta = 0$ .

Dynamical field  $\rightarrow$  **new particle**

# The axion

“Cleaning up the problem with the axial transformation”

[Weinberg, Wilczek]

$$L_a = \theta \cdot G\tilde{G} + \frac{1}{2}f_a^2 \cdot (\partial_\mu\theta)^2 + V(\theta, \partial_\mu\theta)$$

Mass  $\leftrightarrow$  Scale  $m_a^2 = \chi/f_a^2$  with  $\chi = 76(2)(1)\text{MeV}$

Interactions are model dependent

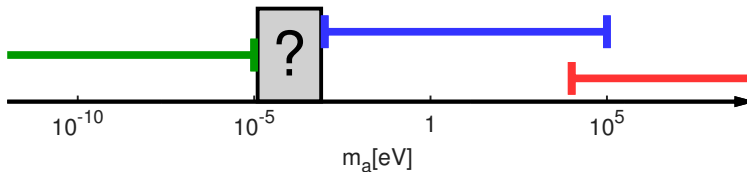
$$V(\theta, \partial\theta) = c \cdot \theta \cdot F\tilde{F} + V_1(\partial\theta)$$



Smaller mass more elusive.

# The axion window

Searching for axions is hard, since mass is unknown.



Exclusions on  $m_a$  from

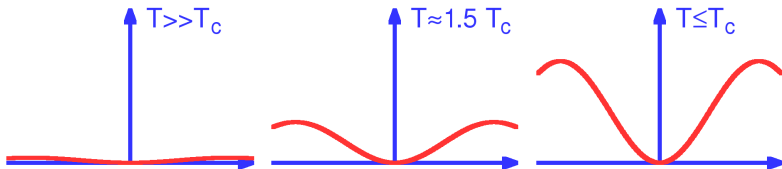
Early laboratory searches

Astrophysics (supernovae, red giants)

Axion is dark matter

# Axion production in the early Universe

Potential becomes flat at QCD transition ( $T_c \approx 150\text{MeV}$ )



Calculate the number of axions produced!

Rolling down the potential ( $\rightarrow \chi(T)$ ) + damped by expansion ( $\rightarrow \epsilon(T), p(T)$  equation of state).

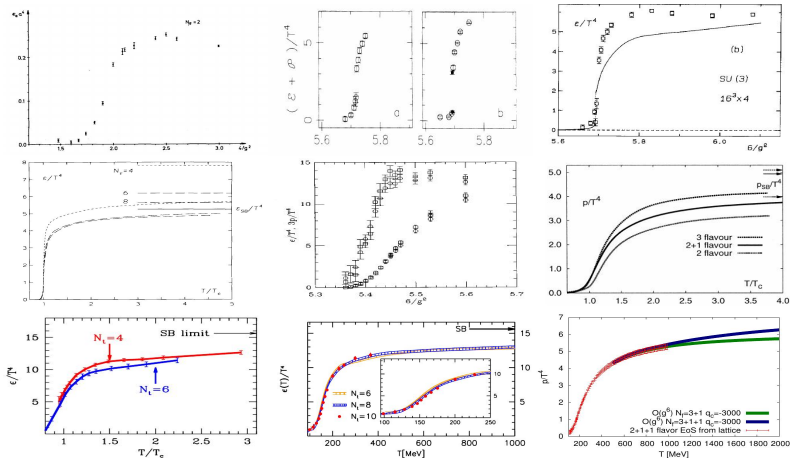
Need lattice QCD!

# Calculation of the axion mass based on high-temperature lattice quantum chromodynamics

S. Borsanyi<sup>1</sup>, Z. Fodor<sup>1,2,3</sup>, J. Guenther<sup>1</sup>, K. -H. Kampert<sup>1</sup>, S. D. Katz<sup>3,4</sup>, T. Kawanai<sup>2</sup>, T. G. Kovacs<sup>5</sup>, S. W. Mages<sup>2</sup>, A. Pasztor<sup>1</sup>, F. Pittler<sup>3,4</sup>, J. Redondo<sup>6,7</sup>, A. Ringwald<sup>8</sup> & K. K. Szabo<sup>1,2</sup>

[doi.org/10.1038/nature20115](https://doi.org/10.1038/nature20115)

# Equation of state from lattice QCD



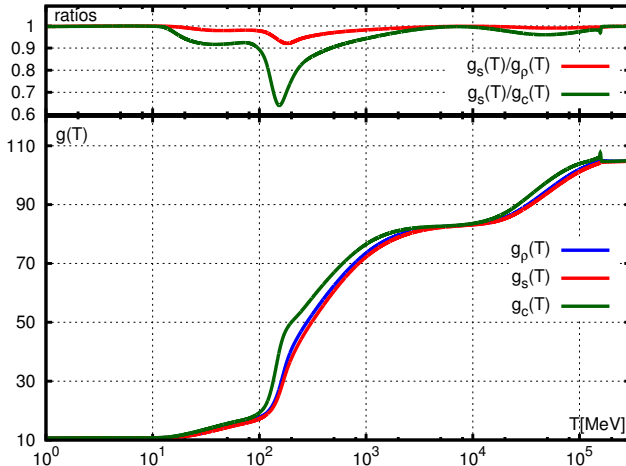
1981: pure SU(2),  $N_t = 2$

2016: SU(3) + u,d,s,c,b; cont. extrap. from  $N_t = 6 \dots 16$

First time without “left for future work”!



# Equation of state

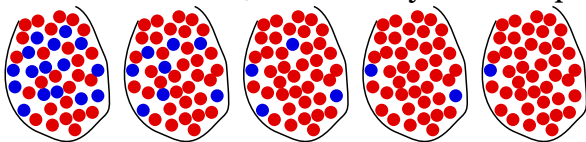


Full result= lattice QCD + weak [Laine,Meyer,Schroeder] +  
photon + neutrinos + leptons

# Determination of axion potential

## Challenge

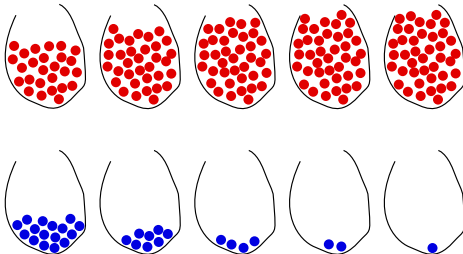
Determine the blue/red ratio by random pick!



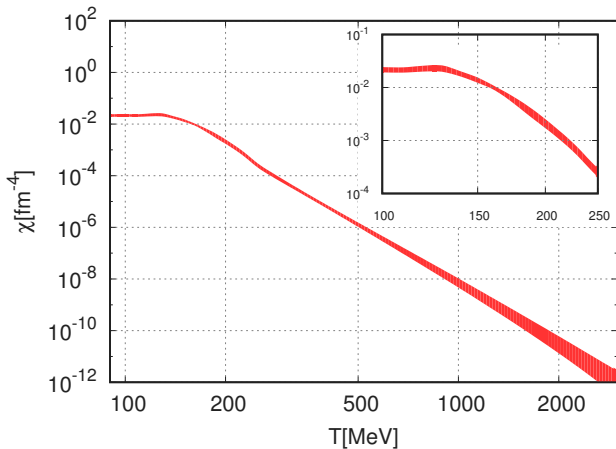
→ getting very difficult with  $T$  →

## Solution

Separate colors and determine the rate of change with  $T$ !



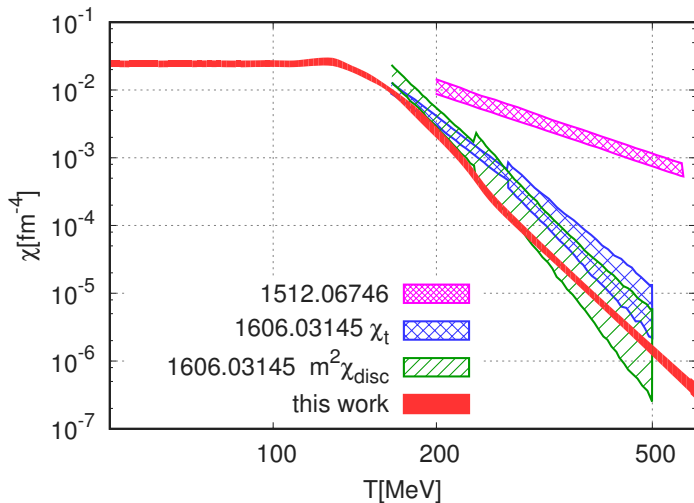
# Axion potential $\chi(T)$



Two challenges to solve:

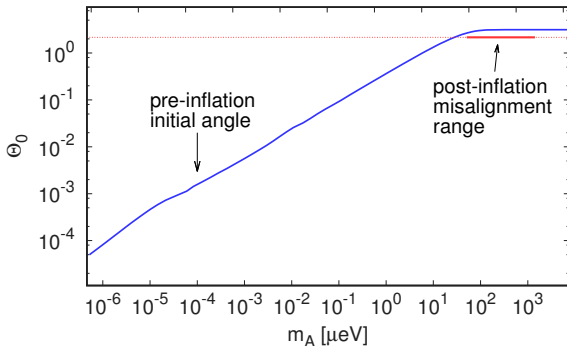
1. signal is small
2. lattice artefacts are large

# Comparison to others



# Results

All dark matter is axion:  $\Omega_{DM} \equiv \Omega_a(m_a, \theta_0) \rightarrow m_a(\theta_0)$



post-inflation: average all possible  $\theta_0$  values  $\rightarrow$

$$\langle m_a(\theta_0) \rangle = 28(1) \mu\text{eV}$$

(If not all DM is axion, then this is a lower bound.)

pre-inflation: single  $\theta_0$  in Universe,  $m_a$  can be anything