

Multiquark Resonances

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CHARGED EXOTIC RESONANCES

BES III
BELLE (2013)

$$e^+e^- \rightarrow \gamma \Upsilon(4.26)$$

$$\hookrightarrow \pi^+ Z_c^- (3.9)$$

$$\hookrightarrow \pi^- J/\psi$$

$$J^{PC} = 1^+ 1^{+-}, \Gamma \sim 50 \text{ MeV}$$

another decay modes of Υ is

$$\Upsilon(4.26) \rightarrow \gamma X(3.872)$$

$$\hookrightarrow DD^*$$

$$J/\psi \pi^+ \pi^-$$

$$J/\psi \pi^+ \pi^- \pi^0$$

Z_c IS A 4-QUARK RESONANCE

$$c\bar{c} d\bar{u}$$

Charged $Z_c(3900)$

Found in $Y(4260) \rightarrow Z_c^\pm(3900) \pi^\mp \rightarrow J/\psi \pi^\pm \pi^\mp$

Exotic charged charmonium-like state!

$$G = G_\pi C_{J/\psi} = -1(-1) = +1$$

$$P = +1 \text{ (S-wave)}$$

$\Rightarrow Z_c^0$ has $J^{PC} = 1^{+-}$

$$I^G J^{PC} = 1^+ 1^{+-}$$

BESIII, arXiv:1303.5949

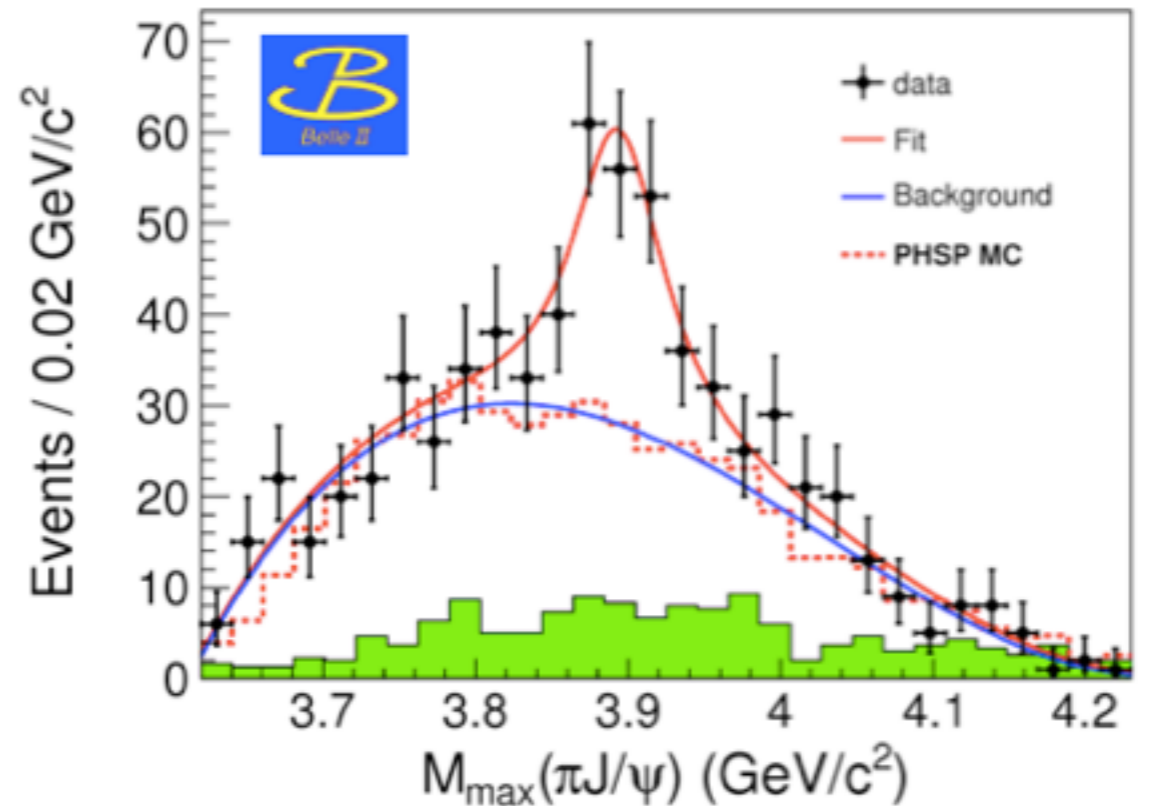
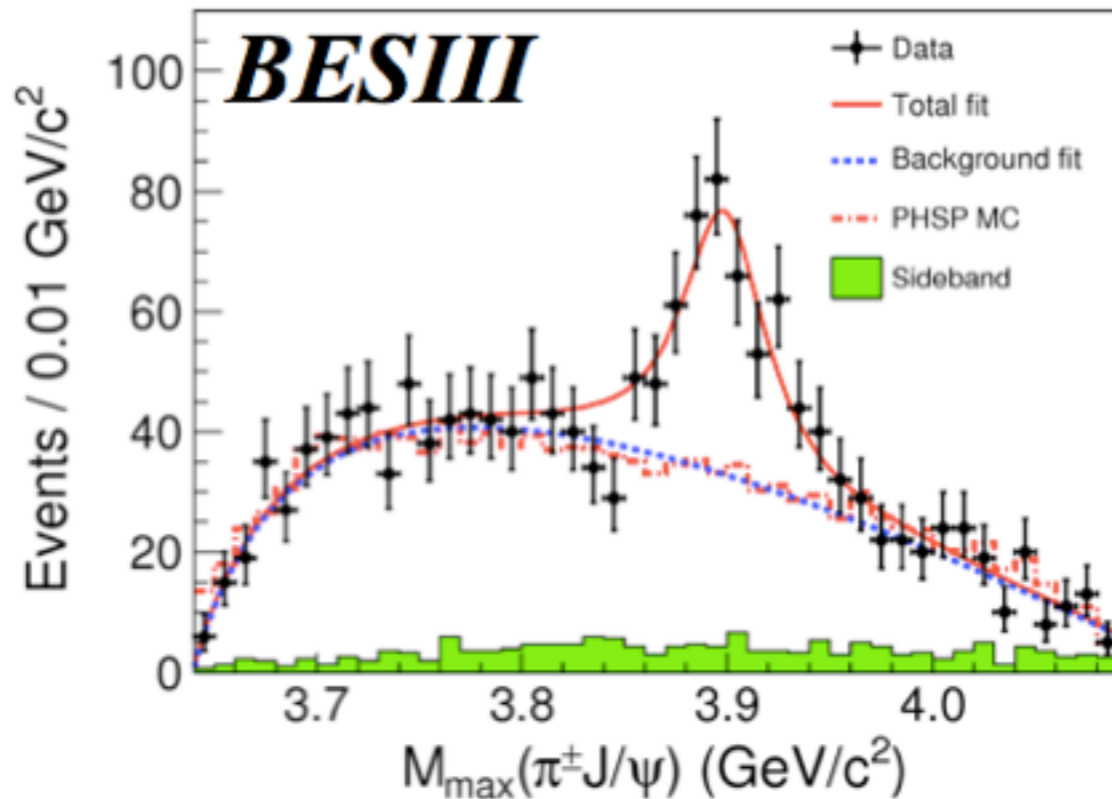
$$M = 3899.0 \pm 3.6 \pm 4.9 \text{ MeV}$$

$$\Gamma = 46 \pm 10 \pm 20 \text{ MeV}$$

Belle, arXiv:1304.0121

$$M = 3894.5 \pm 6.6 \pm 4.5 \text{ MeV}$$

$$\Gamma = 63 \pm 24 \pm 26 \text{ MeV}$$



X(3872)

DISCOVERED BY BELLE IN 2003

CONFIRMED BY BaBar, DØ, CDF, CMS, LHCb & ATLAS!

4 $pp \rightarrow X(3872) @ CMS$

4 Measurement of the cross section ratio

PROMPT PRODUCTION

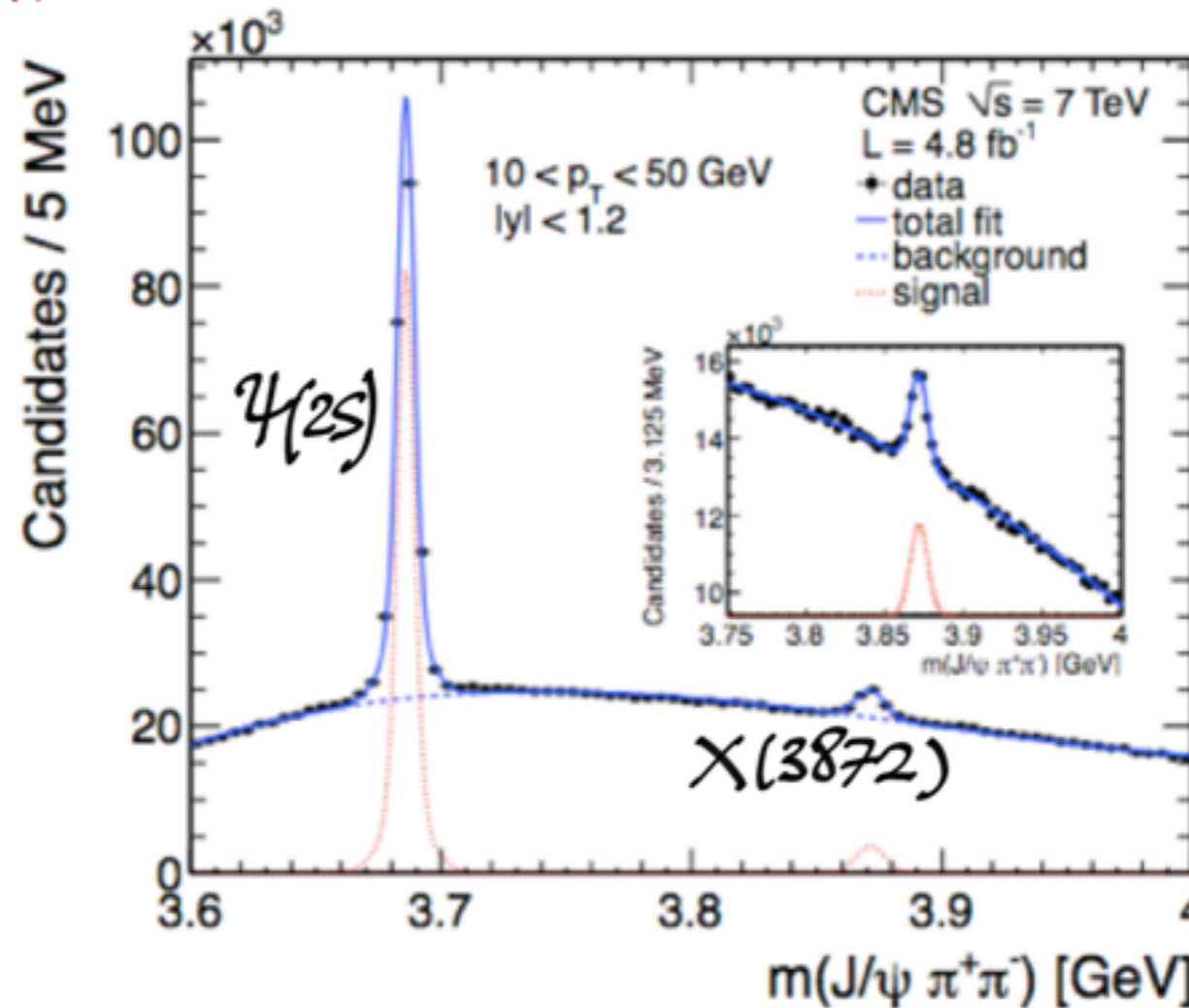


Figure 1: The $J/\psi \pi^+ \pi^-$ invariant-mass spectrum for $10 < p_T < 50$ GeV and $|y| < 1.2$. The lines represent the signal-plus-background fits (solid), the background-only (dashed), and the signal-only (dotted) components. The inset shows an enlargement of the $X(3872)$ mass region.

INTERESTING FACTS ABOUT 'X'

1. $J^{PC} = 1^{++}$

2. The largest BR is

$$X \rightarrow \bar{D}^0 D^{*0}$$

and $M_X \equiv M_{D^0} + M_{D^{*0}} !$

3. ... but it was discovered in another channel

$$X \rightarrow J/\psi \rho^0$$

and $M_X \equiv M_{J/\psi} + M_{\rho^0} !!$

What about $X^\pm \rightarrow J/\psi \rho^\pm$? Never observed.

4. $X \rightarrow J/\psi \omega$

$$\& \frac{BR(X \rightarrow J/\psi \rho^0)}{BR(X \rightarrow J/\psi \omega)} \approx 1$$

Is there any relation between $Z_c(3.9)$ and $X(3.872)$?

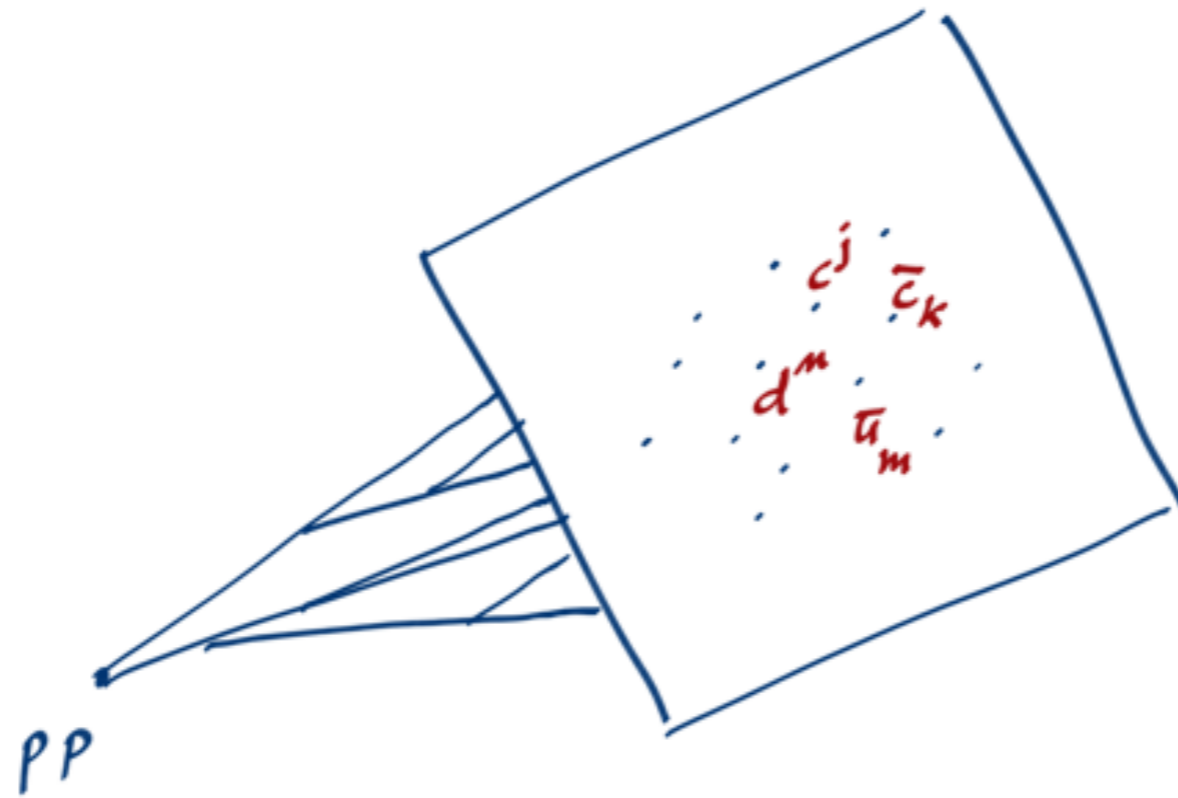
- X is neutral — no charged partner ever found
- Z_c appears in all three states of charge
- X and Z_c almost degenerate
- opposite G-parity

Both compact tetraquarks?... Where is X^+ ?

Both molecules?... One of the two has *positive* binding energy?!

CHARGED EXOTIC RESONANCES

It should be searched in prompt pp collisions at LHC



The relative motion must be compatible with the formation of a compact tetraquark

$$\epsilon_{ijn} c^j d^m \epsilon^{ikm} \bar{c}_k \bar{u}_m$$

CHARGED EXOTIC RESONANCES

According to the VIRIAL THEOREM

$$\bar{T} (= -\bar{E}) \simeq \frac{1}{2} m_c \alpha_c^2 (2m_c) \simeq 50 \text{ MeV}$$

However color might be neutralized also in two singlets

$$\begin{array}{l} c^j \bar{u}_j \quad \bar{c}_k d^k \\ \text{or} \\ c^j \bar{c}_j \quad \bar{u}_k d^k \end{array}$$

in all spin configurations preserving the spin of the heavy quark pair.

Bound states may not be formed (continuum spectrum)
yet we may expect

$$E < E_{\text{max}} < |\bar{E}|$$

'HADRONIZATION STATE'

Superposition with unknown coefficients

$$\Psi = [cd][\bar{c}\bar{u}] + \psi\pi^- + \psi'\pi^- + \eta_c\rho + \bar{D}D^* + \bar{D}^*D$$

$$= \psi_d + \sum_i \psi_{m_i}$$

$$\Psi = \psi_Q + \psi_P = Q\Psi + P\Psi$$

$$H\Psi = E\Psi$$

$$(E - H_{PP})\psi_P = H_{PQ}\psi_Q$$

$$(E - H_{QQ})\psi_Q = H_{QP}\psi_P$$

$$H_{PP} = H_0 + V_1$$

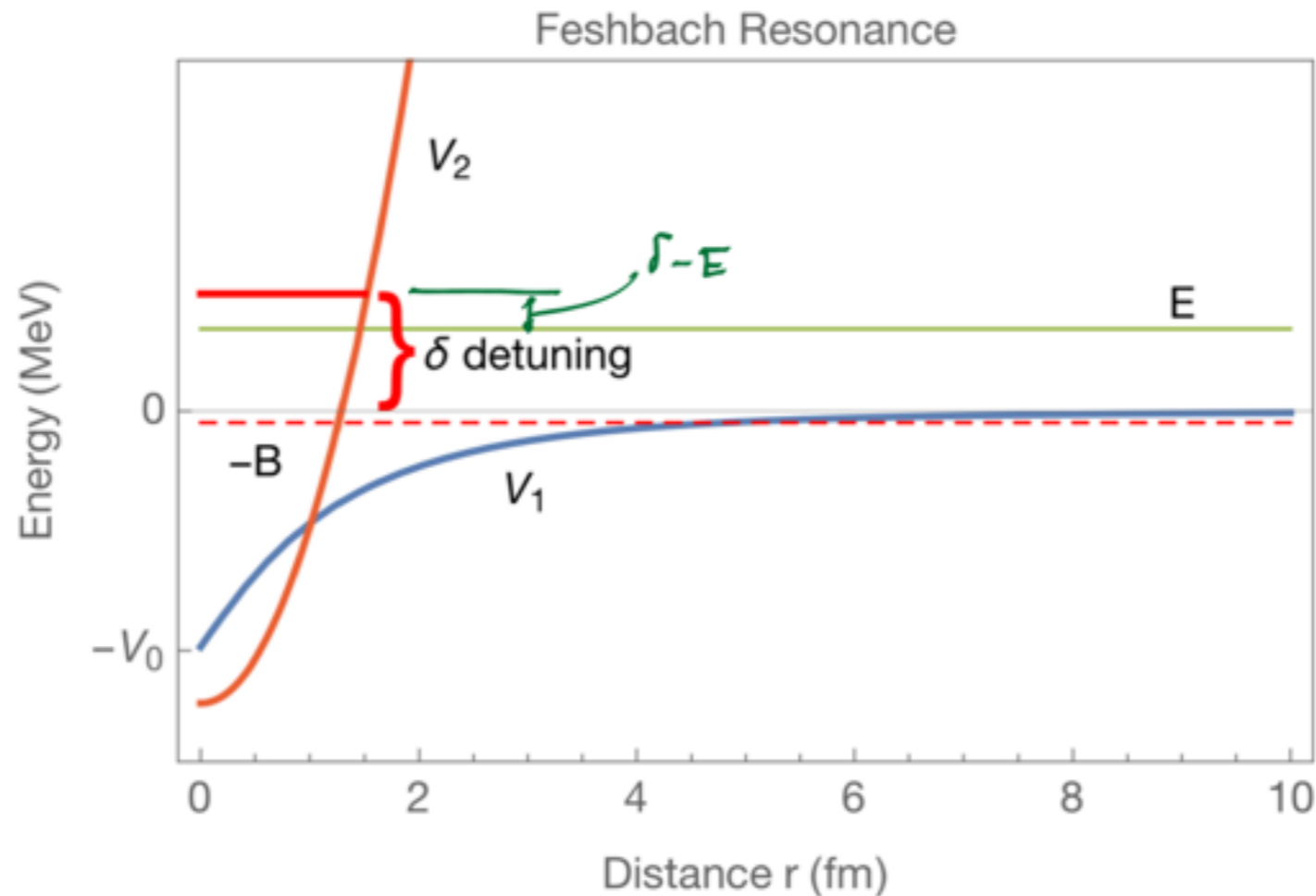
$$H_{QQ} = H_0 + V_2$$

$$\Leftrightarrow \boxed{(E - H_{PP} - V_I)\psi_P = 0}$$

EFFECTIVE INT.
IN THE P SPACE
($P \rightarrow Q \rightarrow P$)

$$V_I = H_{PQ} \frac{1}{E - H_{QQ} + i\epsilon} H_{QP}$$

OPEN & CLOSED CHANNELS



Because of V_I the scattering length in P is

$$a = a_p - c \frac{|\langle \psi_m | H_{QP} \psi_\alpha \rangle|^2}{E_m - E_\alpha + i\epsilon}$$

$(c > 0)$

$$\equiv \left(1 - \frac{2c}{\delta - E + i\epsilon} \right) a_p$$

SCATTERING IN THE OPEN CHANNEL

Would-be hadron molecules in P , may momentarily rearrange their internal structure into compact tetraquarks: this happens w/ probability

$$\sigma \sim \frac{|2ma|}{p} \sim \frac{1}{p} \frac{\epsilon}{(\delta - E)^2 + \epsilon^2}$$

$$d\Gamma \sim p v \sigma \sim \delta(E - \delta) |2ca_p| (2m)^{1/2} \sqrt{E} dE$$

$$E < E_{\max}$$

If δ falls within E_{\max}

$$\Gamma \sim (2m)^{1/2} |2ca_p| \sqrt{\delta}$$

THIS MIGHT HAPPEN ONLY FOR ONE OF THE THRESHOLDS IN ψ , THE CLOSER (in energy) FROM BELOW TO $E_{n=0}$ -

CHARGED EXOTIC RESONANCES

The $Z_c(3.99)$ is not the only one.

$$Z_c^{\prime\bar{F}}(4.025) \rightarrow \pi^{\bar{F}} h_c$$

$$Z_b^{\bar{F}}(10.61) \rightarrow \pi^{\bar{F}} \Upsilon(n\phi)$$

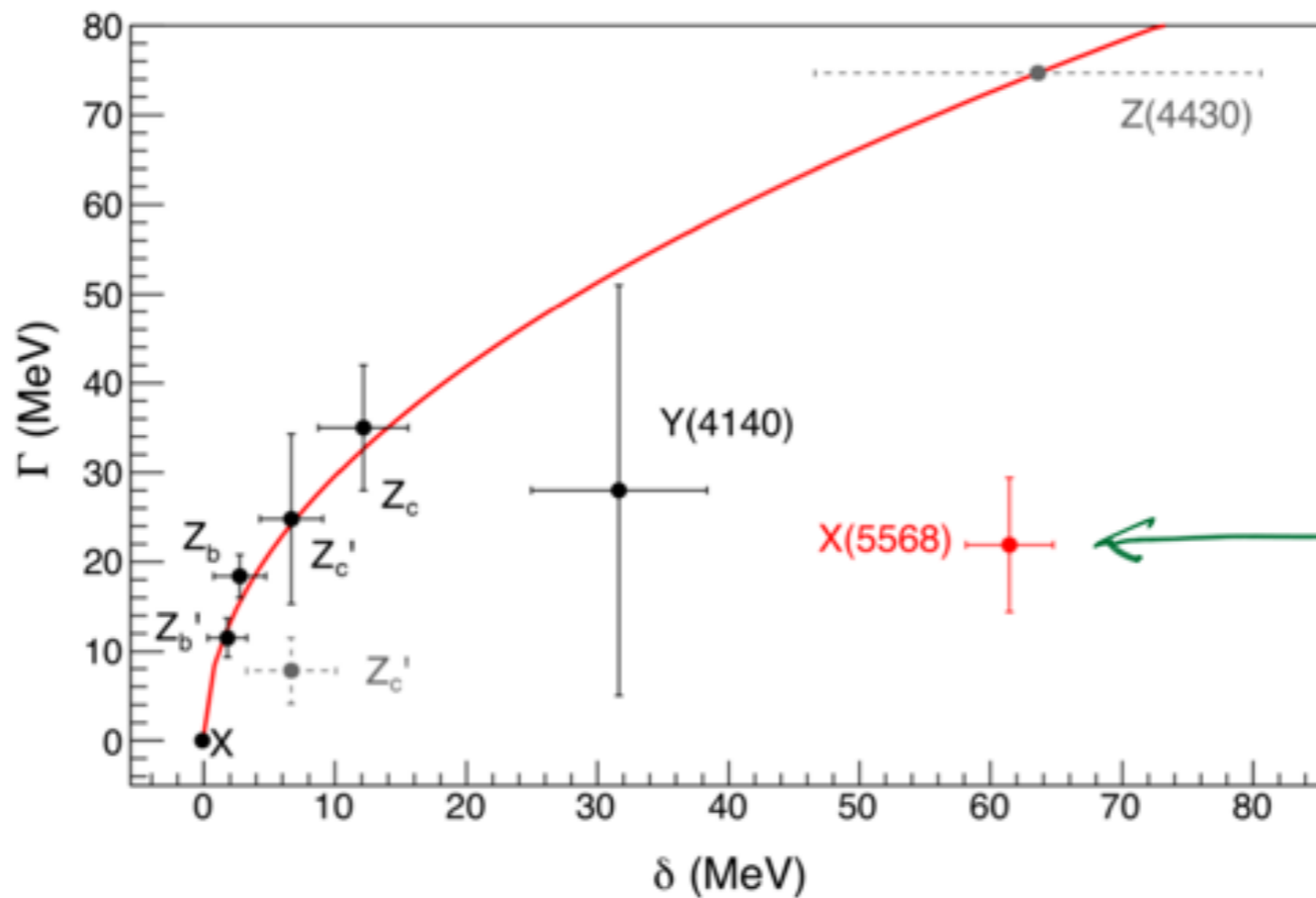
BELLE(2012)

$$Z_b^{\prime\bar{F}}(10.65) \rightarrow \pi^{\bar{F}} h_b(kP)$$

	MOL	δ
Z_c	DD^*	~ 24 MeV
Z_c'	D^*D^*	~ 8 MeV
Z_b	BB^*	~ 6 MeV
Z_b'	B^*B^*	≈ 1 MeV

' \sqrt{s} - LAW'

A. ESPOSITO, A. PILLONI, ADP 1603.07667 (PLB)



THE DIQUARK.
LEVEL E_0 CAN
BE COMPUTED

X $\bar{D}^0 D^{*0}$
 Z_c $\bar{D}^0 D^{*+}$
 Z_c' $\bar{D}^{*0} D^{*+}$
 Y $J/4 \phi$
 Z_b $\bar{B}^0 B^{*+}$
 Z_b' $\bar{B}^{*0} B^{*+}$

X(5568) $B_s^0 \pi^+$

FEB 2016

A NEW $B_s^0 \pi^\pm$ STATE CLAIMED BY DØ

[DØ: arXiv:1602.07588]

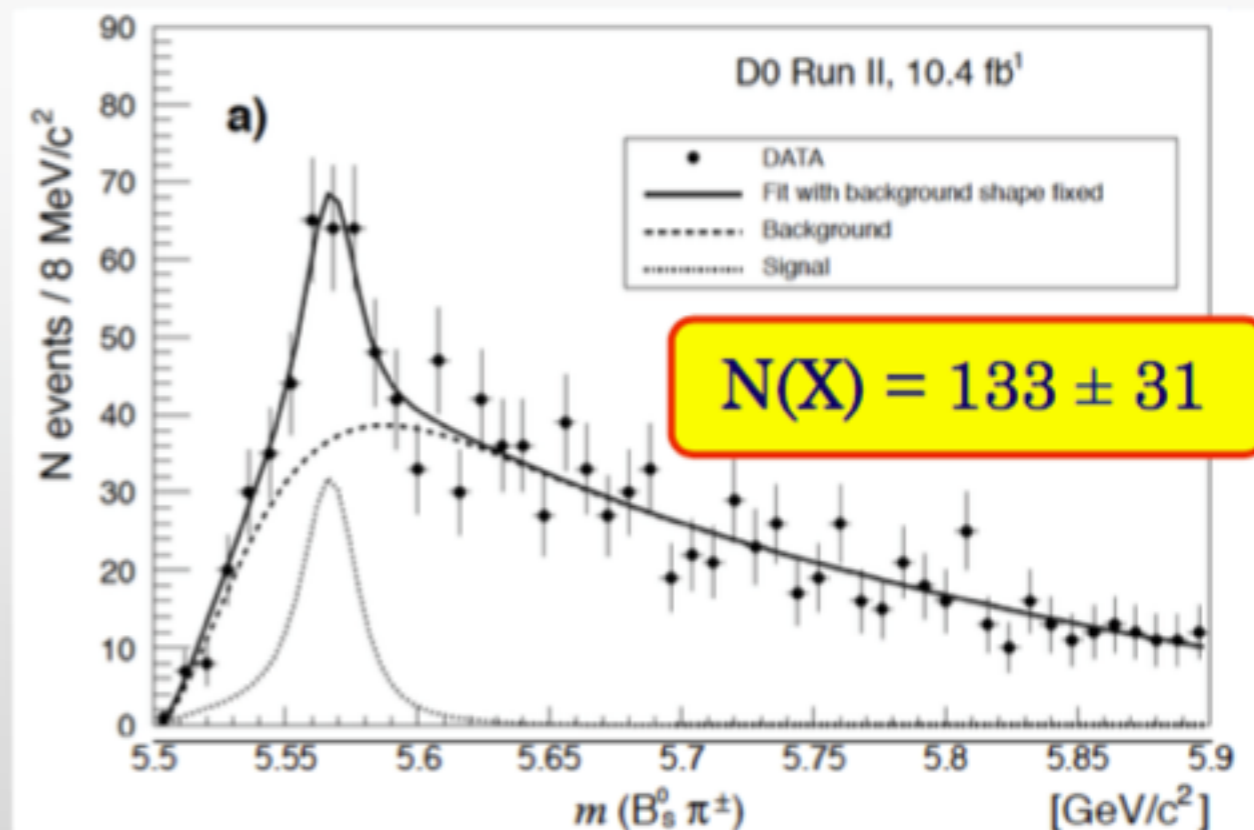
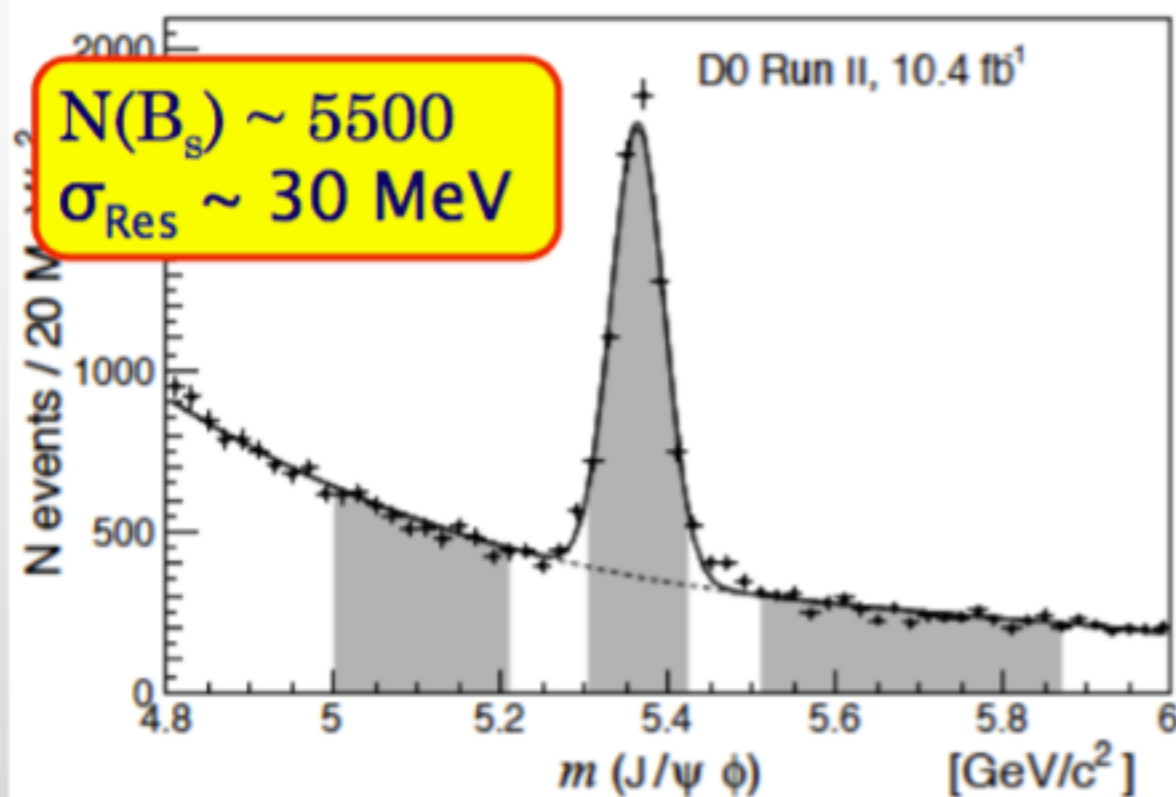
Claimed observation with 5.1σ significance of an exotic state

✓ $X(5568) \rightarrow B_s^0 \pi^\pm$, $B_s^0 \rightarrow J/\psi \phi$, $J/\psi \rightarrow \mu^+ \mu^-$, $\phi \rightarrow K^+ K^-$

$$M = 5567.8 \pm 2.9_{-1.9}^{+0.9} \text{ MeV}/c^2$$

$$\Gamma = 21.9 \pm 6.4_{-2.5}^{+5.0} \text{ MeV}/c^2$$

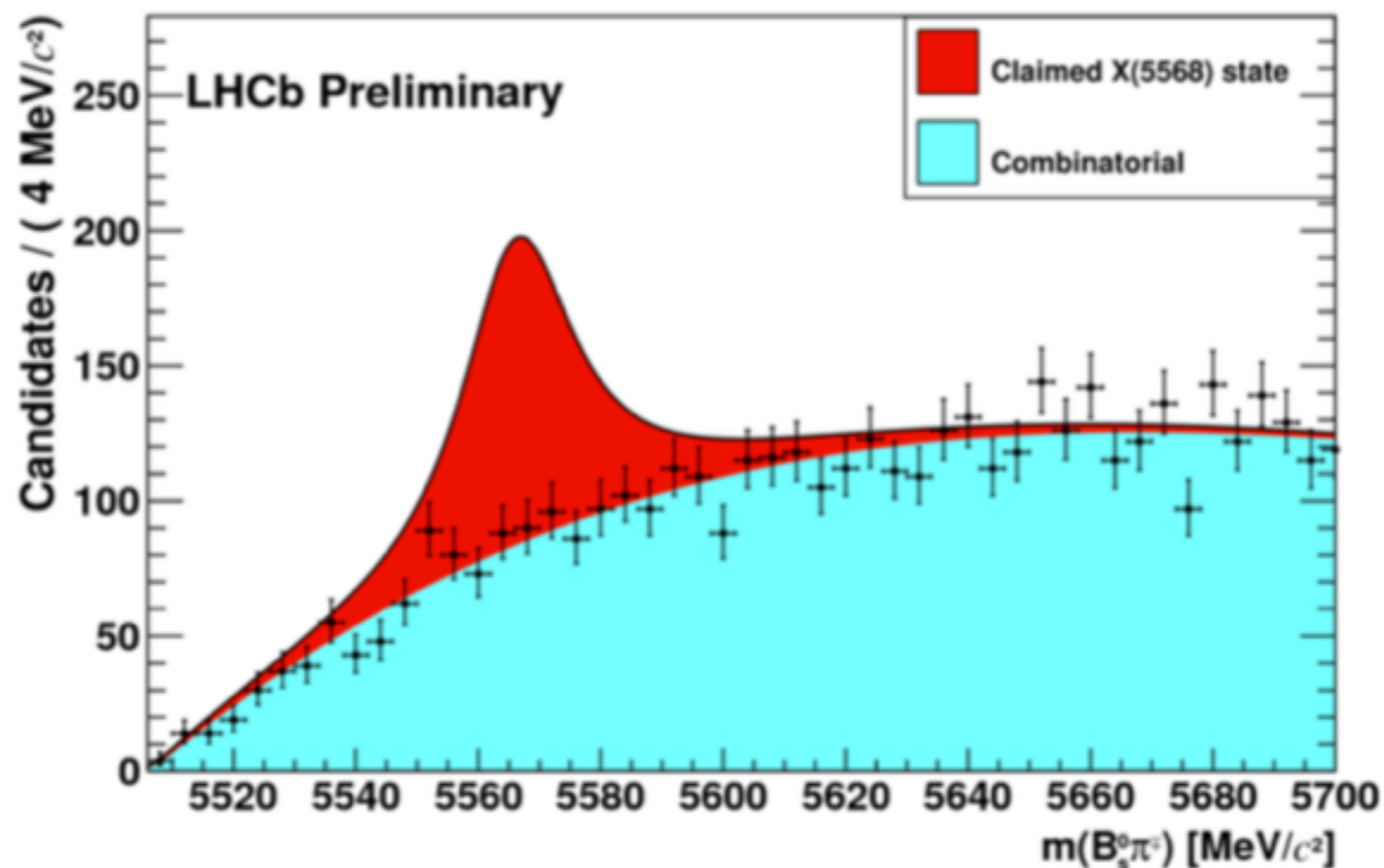
✓ Fraction of B_s^0 from X decay: $\rho_X^{\text{DØ}} = (8.6 \pm 1.9 \pm 1.4) \%$



JUST FOR CURIOSITY...

If $\rho_X^{\text{LHCb}} = \rho_X^{\text{D}\emptyset} = 8.6\%$, how would the X(5568) signal look like?

(Both modes combined: $p_T(B_s) > 10 \text{ GeV}/c$)



STRONG INTERACTION FESHBACH PHENOMENON?

1. Not all meson-meson thresholds correspond to a resonance (loosely bound molecules)
2. Not all 'diquarkonia' manifest in the spectrum.

A FAVOURABLE DETUNING MUST OCCUR

$$\delta \in [0, E_{\max}]$$

This might explain why some observed isospin multiplets are incomplete.

- A pure diquarkonium theory predicts

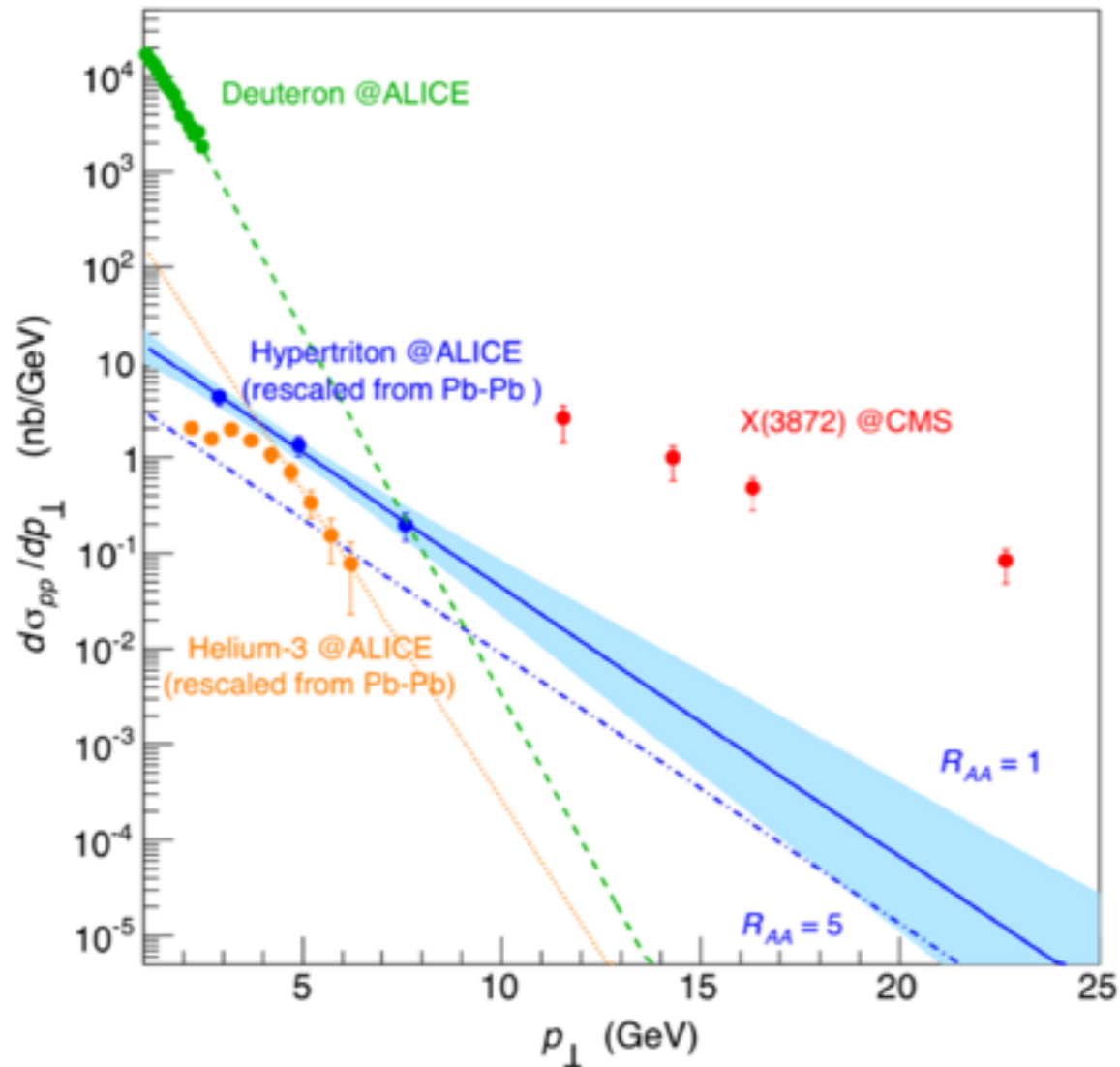
$$X^0, X^+, X^- @ 3872 \text{ MeV}$$

NOT (YET??) SEEN.

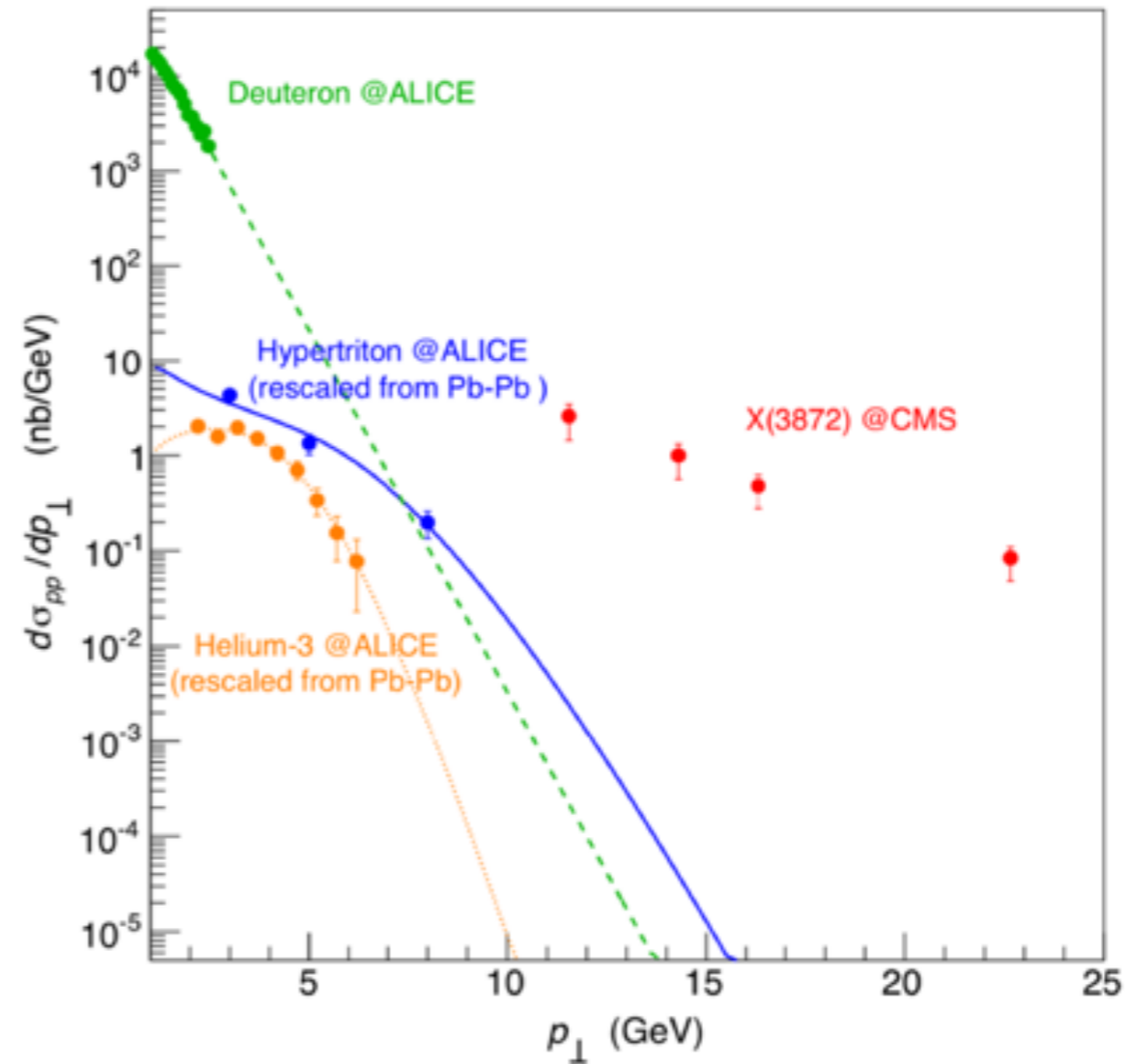
- What about X_b ?

DATA FROM ALICE

For those who think that we are observing real hadron molecules



Exponential fit -



'blast wave' fit -

A Esposito et al. PRD 92 (2015) 034028

Diquarks and Exotic Hadrons

Jaffe and Wilczek hep-ph/0307341 (PRL)

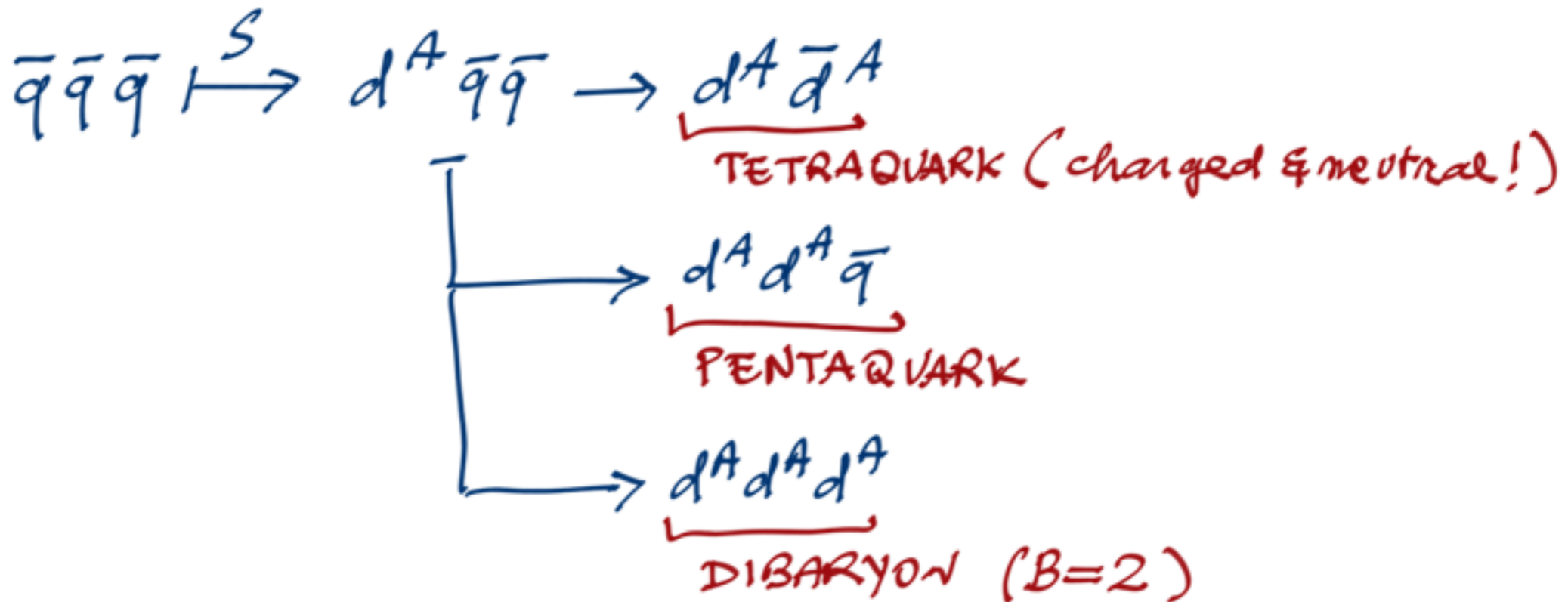
Maiani, Piccinini, ADP, Riquer, hep-ph/0412098 (PRD)

DIQUARKS

$$d^{S,A} = q^\alpha \Gamma q'^\beta \pm q^\beta \Gamma q'^\alpha$$

BUILD NEW HADRONS WITH

$$\begin{cases} q \mapsto \bar{d}^A \\ \bar{q} \mapsto d^A \end{cases}$$



DIQUARKONIA

$$\left\{ \begin{array}{l} H \approx 2 k_{Qq} \cdot (\vec{S}_q \cdot \vec{S}_Q + \vec{S}_{\bar{q}} \cdot \vec{S}_{\bar{Q}}) \\ \text{the spin of heavy-light diquarks can be } 0, 1 \\ \text{and is conserved in strong interactions.} \end{array} \right.$$

From phenomenology

$$|S_{c\bar{c}}, S_{q\bar{q}}\rangle \left\{ \begin{array}{l} Z = \frac{1}{\sqrt{2}} (|1, 0\rangle - |0, 1\rangle) \\ Z' = \frac{1}{\sqrt{2}} (|1, 0\rangle + |0, 1\rangle) \end{array} \right.$$

$C = (-)^{L+S_{q\bar{q}}+S_{c\bar{c}}}$

which in the diquarkonium basis

$$|S_{cq}, S_{\bar{c}\bar{q}}\rangle \left\{ \begin{array}{l} Z = \frac{1}{\sqrt{2}} (|1, 0\rangle - |0, 1\rangle) \\ Z' = |1, 1\rangle_{L=S} \quad C = (-)^J \\ X = \frac{1}{\sqrt{2}} (|1, 0\rangle + |0, 1\rangle) \quad 1^{++} \end{array} \right.$$

DIQUARKONIA

$$(H)^{1+-} = \begin{vmatrix} -k & 0 \\ 0 & k \end{vmatrix} + 2m_{[qq]} \mathbb{1} \rightarrow \begin{aligned} M_{2'} - M_2 &= 2k \\ M_2 + M_{2'} &= 4m_{[qq]} \end{aligned}$$

$$(H)^{1++} = -k + 2m_{[qq]} \mathbb{1} \rightarrow M_2 \simeq M_X$$

Therefore

$$m_{[bq]} \simeq 5315 \text{ MeV}$$

$$k_{bq} \simeq 22.5 \text{ MeV}$$

and for the $D\phi$ state $[\bar{b}\bar{q}][sq']$ in $B_S\pi$

$$M = m_{[bq]} + m_{[sq]} + 2k_{bq} \vec{S}_{\bar{b}} \cdot \vec{S}_{\bar{q}} + 2k_{sq} \vec{S}_s \cdot \vec{S}_{q'}$$

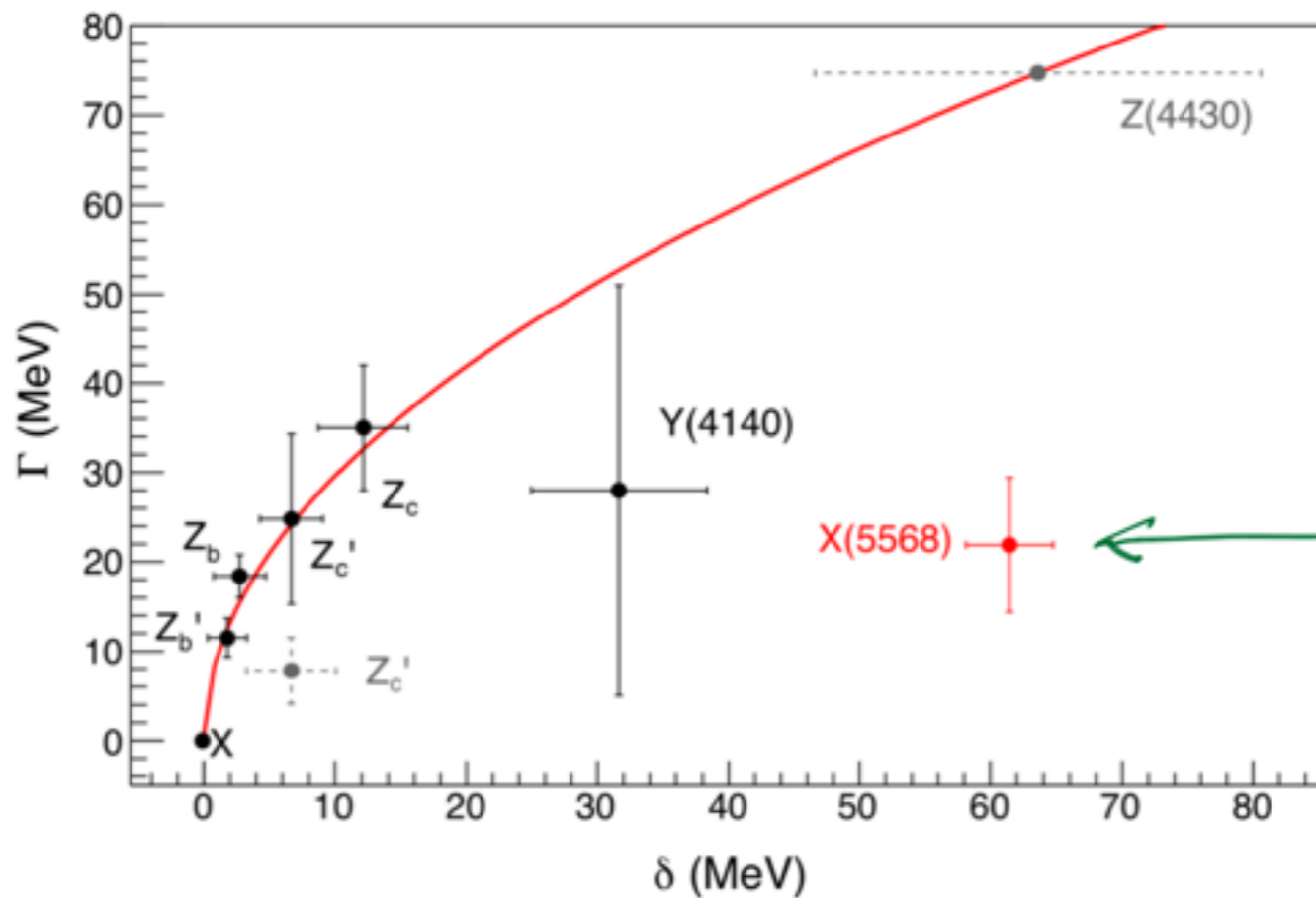
$$= m_{[bq]} + (m_{[sq]} - 3/2 k_{sq}) - 3/2 k_{bq}$$

$$\simeq 5771 \text{ MeV (vs 5568 observed)}$$

too large detuning δ w.r.t Γ_{exp}

' \sqrt{s} - LAW'

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X(5568) $B_s^0 \pi^+$

FEB 2016

DIQUARKONIA

$$(H)^{1+-} = \begin{vmatrix} -k & 0 \\ 0 & k \end{vmatrix} + 2m_{[qq]} \mathbb{1} \rightarrow \begin{aligned} M_{Z'} - M_Z &= 2k \\ M_Z + M_{Z'} &= 4m_{[qq]} \end{aligned}$$

$$(H)^{1++} = -k + 2m_{[qq]} \mathbb{1} \rightarrow M_Z \simeq M_X$$

$$(H)^{2++} = k$$

$$(H)^{0++} = -3k$$

$$(H)^{0++'} = k$$

} more diquarkonia ...

but no good matches
in the 'p-channel'

Pentaquarks

based on 1507.04980 with L. Maiani and V. Riquer (Sapienza U.)

THE PENTAQUARK

Highly undesirable option for molecules (before discovery)

Perfect molecule (after discovery)

LHCb 2015

$$\Lambda_b(bud) \rightarrow K^- P^+ \\ \hookrightarrow J/\psi p$$

$$P^+ = \bar{c} c u u d \Rightarrow \text{negative parity}$$

TWO STATES OBSERVED

$$J^P = 3/2^- \text{ @ } 4380 \text{ MeV}$$

$$J^P = 5/2^+ \text{ @ } 4550 \text{ MeV}$$

$L=0$ & $L=1$ Pentaquarks?

Note: Lower baryons have $P=+$ / pentaq. have $P=-$!
Lower mesons have $P=-$ / tetraq. have $P=+$!

MASS DIFFERENCE

ISN'T $\Delta M = 170 \text{ MeV}$ too **SMALL** for one unit of L ?

($\Delta M = 300 \text{ MeV}$ for $\Lambda(1405) - \Lambda(1116)$)

On the other hand, from $\Sigma_c - \Lambda_c$ we find

$$M_{[qq']_{S=1}} - M_{[qq']_{S=0}} \simeq 200 \text{ MeV}$$

So

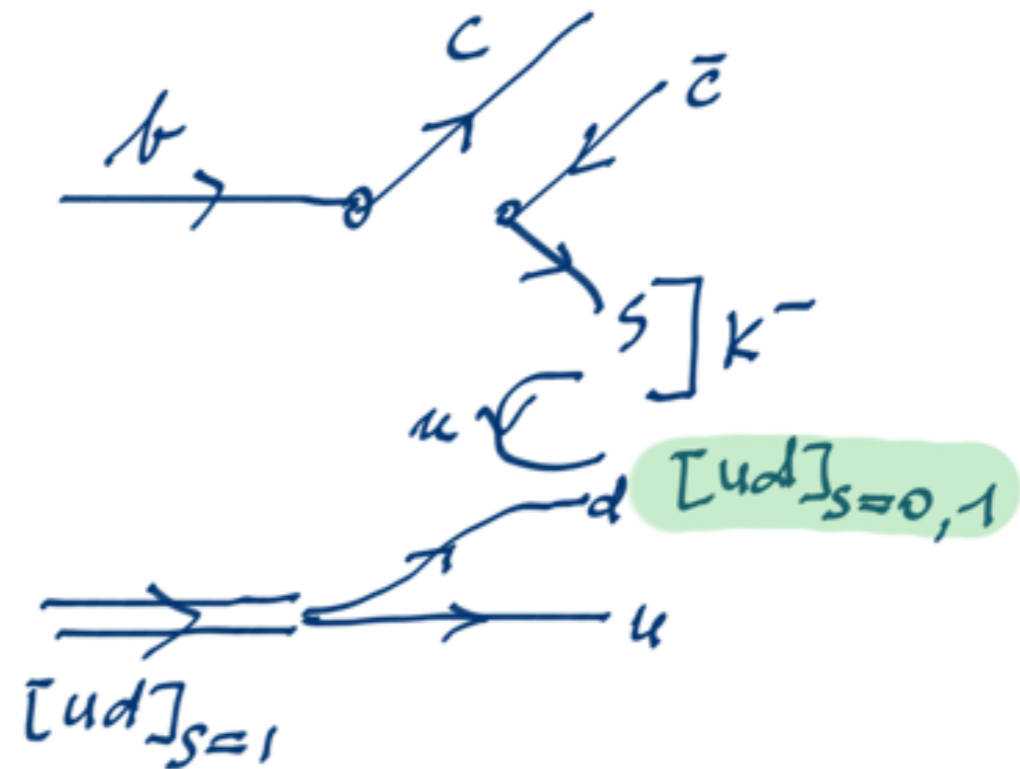
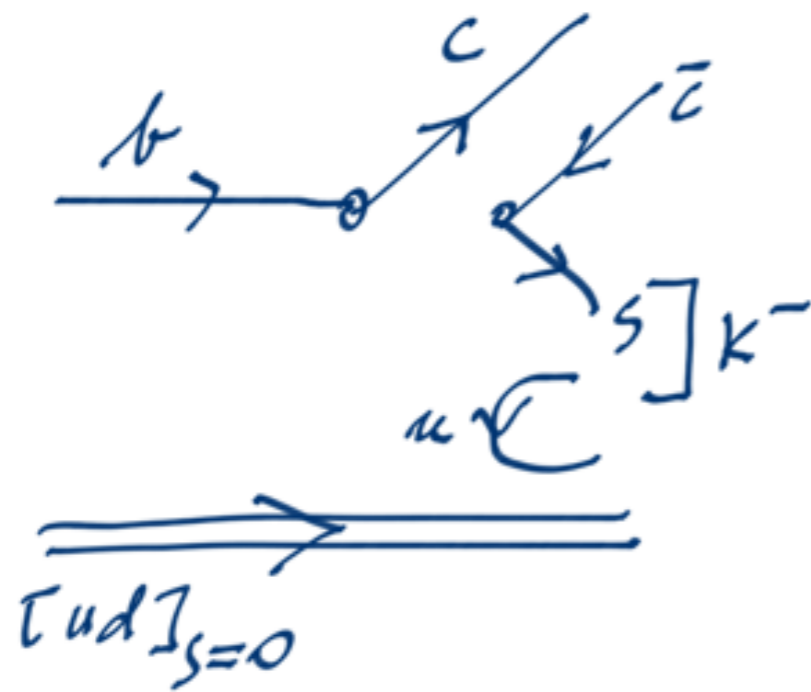
$$P(3/2^-) = \bar{c} [cq]_{S=1} [q'q'']_{S=1} @ L=0$$

$$P(5/2^+) = \bar{c} [cq]_{S=1} [q'q'']_{S=0} @ L=1$$

... combine diquark spin & orbital angular momentum -
Other states?

$\Lambda \rightarrow K^- P^+$

Λ_b baryon might contain a $[ud]_{S=0}$, "good" diquark.
 but $P(3/2^-)$ should contain $[ud]_{S=1}$, whereas $P(5/2^+)$
 has $[ud]_{S=0}$.



One can show that both pentaquarks have $S_{c\bar{c}} = 1$ so that
 HQ spincons. allows decay into J/ψ .

Flavor

$$\langle P, M | H_w (\Delta I=0, \Delta S=-1) | \Lambda_b \rangle$$

8_F

3_F

$\bar{3}_F$

(from s, d, u)

(from [ud])

therefore P is either 8 or $10_{-}^{(*)}$

We might expect

$$\Lambda_b \rightarrow \pi P_{10}^{S=-1} \rightarrow \pi J/\psi \Sigma(1385)$$

$$\Lambda_b \rightarrow K P_{10}^{S=-2} \rightarrow K J/\psi \Xi(1530)$$

or even

$$\Sigma_b \rightarrow \phi P_{10}^{S=-3} \rightarrow \phi J/\psi \Sigma^-(1672)$$

$$(*) \begin{cases} 8 \otimes 10 = 8 \oplus 10 \oplus 27 \oplus 35 \\ 8 \otimes 8 = 1 \oplus 8 \oplus 8 \oplus 10 \oplus \bar{10} \oplus 27 \end{cases}$$

Large N and tetraquarks

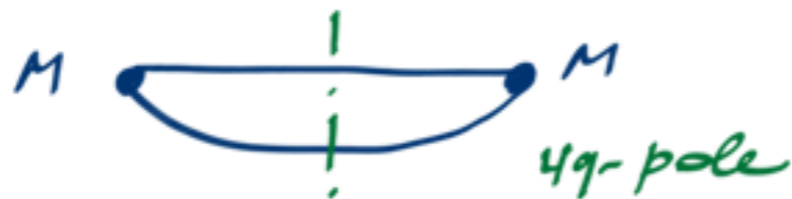
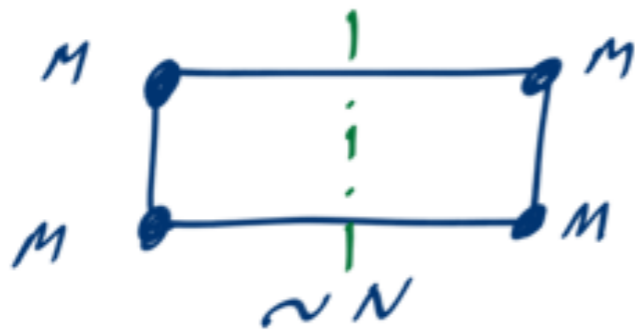
based on 1605.04839 with L. Maiani and V. Riquer

see G. Rossi and G. Veneziano 1605.04285 for an alternative approach

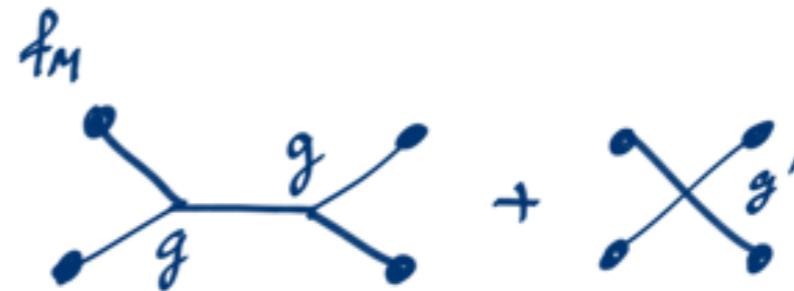
Large N-QCD & Tetraquarks

Reputation of tetraquarks obscured by some considerations by Witten and Coleman (Witten Nucl. Phys. B160 (1979))

Quark theory



Meson theory



$$g \sim \frac{1}{\sqrt{N}}$$



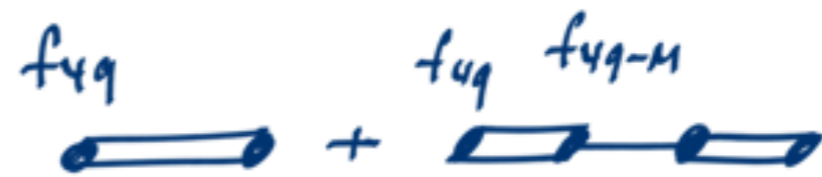
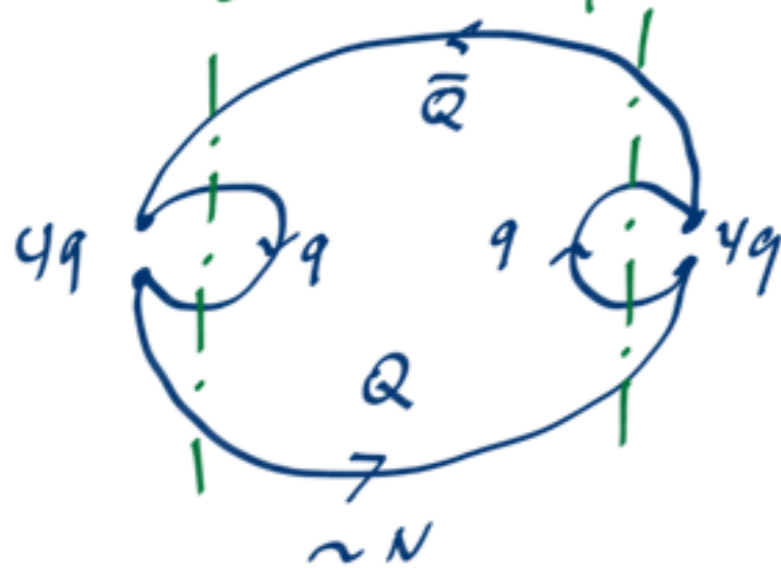
Large N-QCD & Tetraquarks

IF TETRAQUARKS DEVELOP A POLE, IT WILL BE IRRELEVANT IF THE RESIDUE IS OF ORDER $1/N$ WRT DISCONNECTED PARTS.

[S. Weinberg PRL 2013]

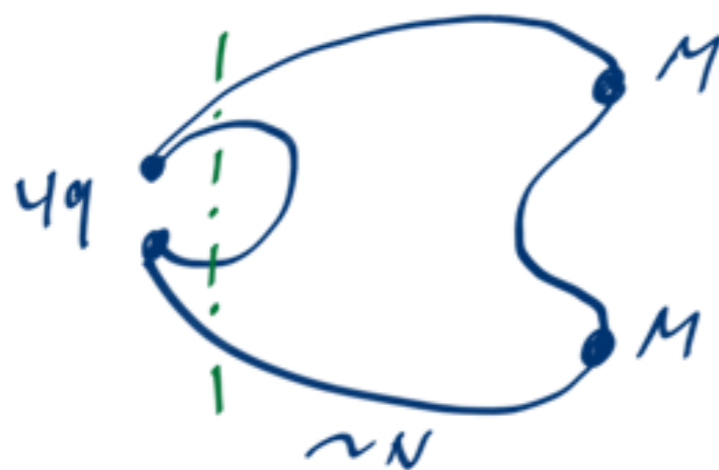
MOREOVER THE WIDTH $\Gamma_{4q} \sim 1/N$.

Consider the case of neutral tetraquarks



$$f_{4q} \sim \sqrt{N}$$

$$f_{4q-M} \sim N^0$$



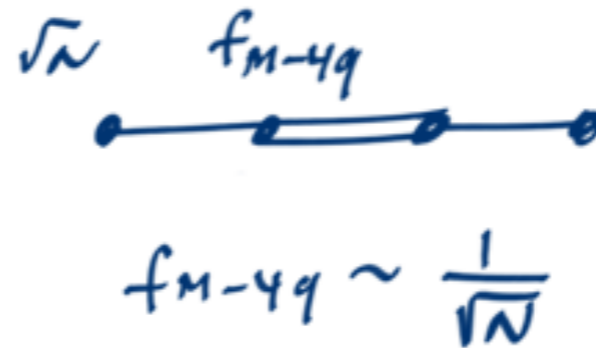
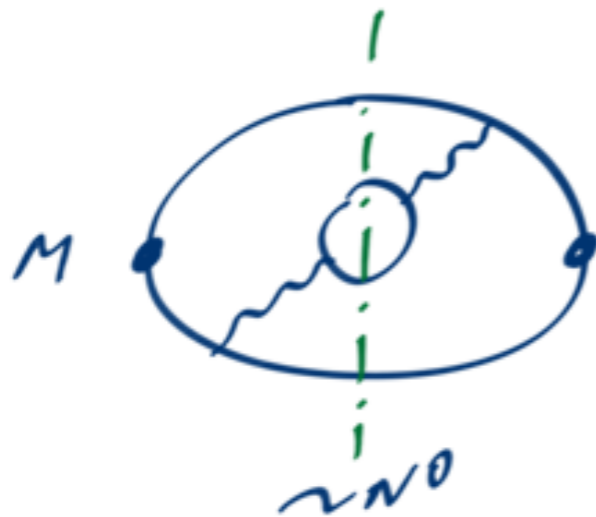
$$G \sim \frac{1}{\sqrt{N}}$$

$$G \sim N^0 \cdot \frac{1}{\sqrt{2}}$$

Large N - QCD & Tetraquarks

16/05/16

1. Aren't those cuts equivalently leading to Witten's argument?
2. Why the mixing should appear at a different N order in different diagrams?

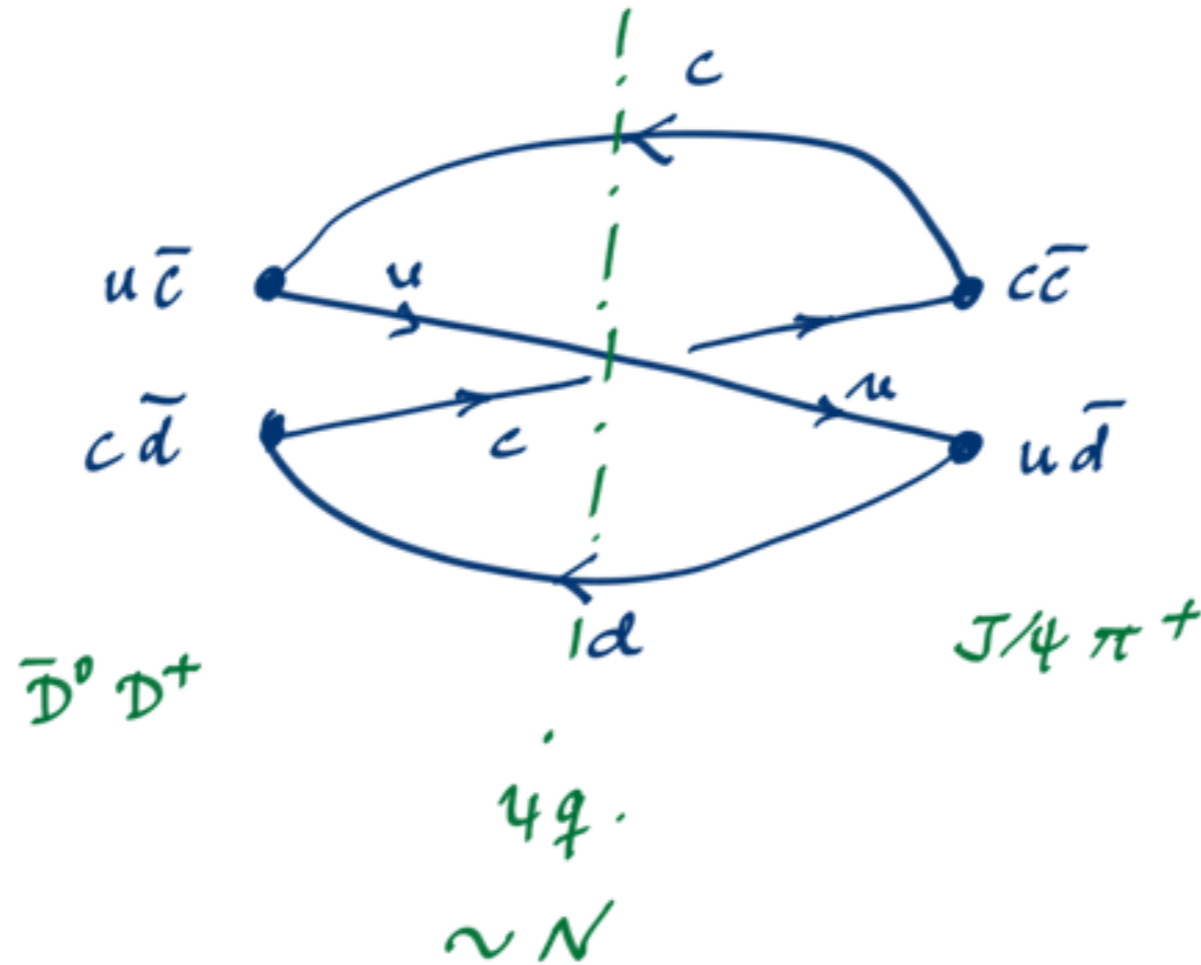


15 / 15

4qgj

Large N-QCD & Tetraquarks

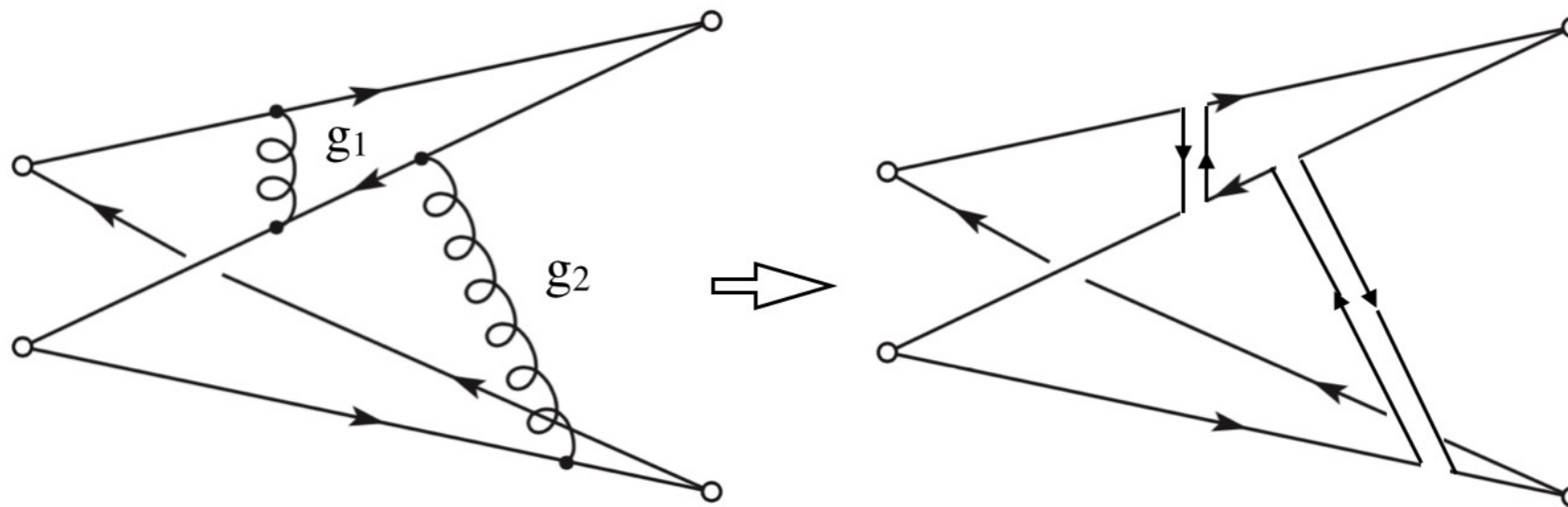
Consider a charged tetraquark compl. funct.



Can be 'untwisted': does it really contain a tetraquark pole?

ONE CAN SHOW THAT DIQUARKONIA GIVE THE RIGHT DESCRIPTION OF THE INTERMEDIATE $4q$ -STATE!

Non-planar diagrams

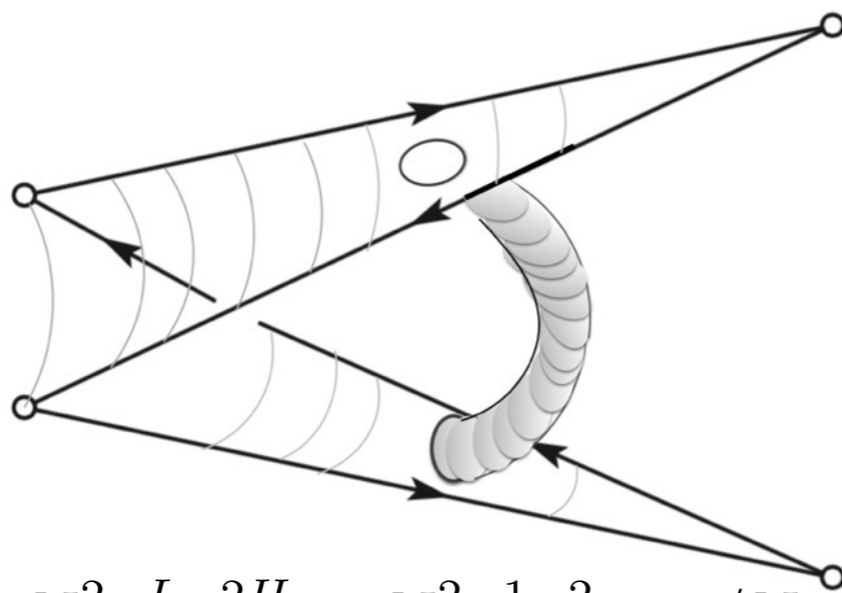


(a)

Non-planar gluons produce the interaction needed to make the color singlet bilinear to merge into a tetraquark.

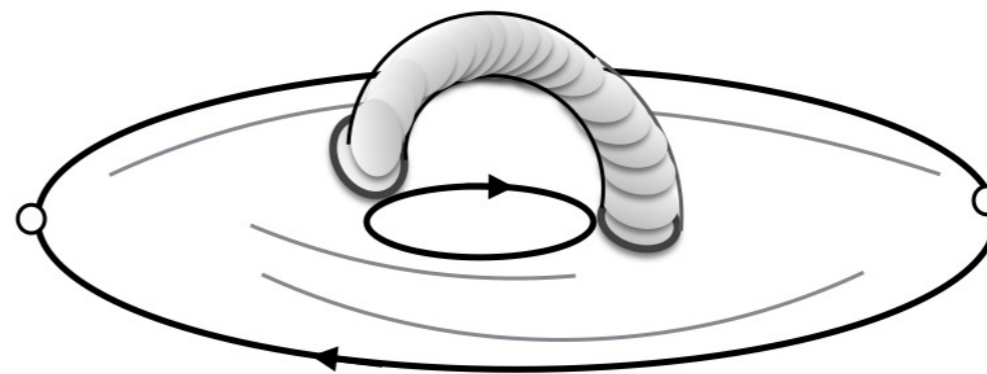
$$f_{4q\text{-charged}}^2 \sim \frac{1}{N}$$

$$f_{4q} \cdot g \cdot \sqrt{N}\sqrt{N} \sim \frac{1}{N} \Rightarrow g \sim \frac{1}{N\sqrt{N}}$$



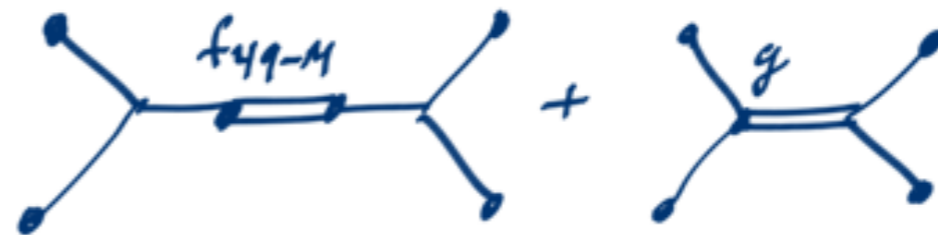
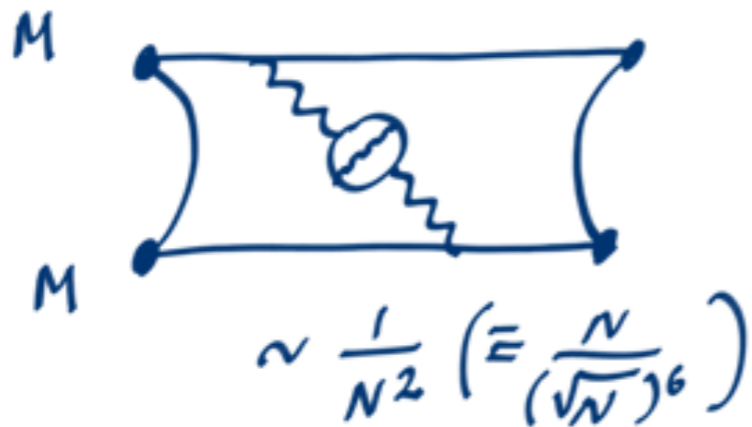
$$N^{2-L-2H} = N^{2-1-2} = 1/N$$

(b)



(c)

Large N-QCD & Tetraquarks



$$g \sim \frac{1}{N^2}$$

$$f_{4q-M} \sim \frac{1}{N\sqrt{N}}$$



SUMMARY

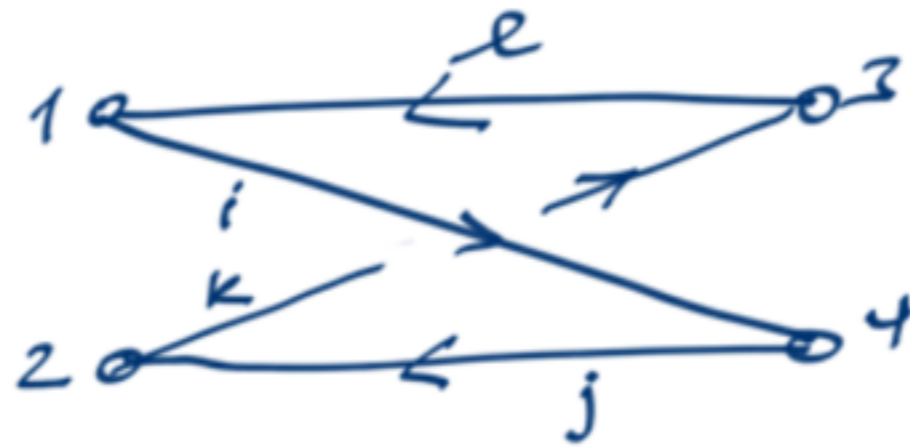
About 20 (exotic) resonances have been discovered

1. Some have complete I-multiplets, some do not
2. Most of them close to M-M thresholds (above!)
3. Diquarks predict charged states & pentaquarks -
But need complete I-multiplets -
Several quantum numbers predicted $0^{++}, 2^{++} \dots$
4. Very similar problems with loosely bound molecules -

— A FESHBACH PHENOMENON AT WORK IN
STRONG INTERACTIONS?

Combining diquarkonia predictions and threshold
positions, a new picture might emerge —

Backup



$$(\Gamma_1)_{il} (\Gamma_4^+)_{ji} (\Gamma_2)_{kj} (\Gamma_3^+)_{lk} = \langle \Gamma_1 \Gamma_3^+ \Gamma_2 \Gamma_4^+ \rangle$$

(spin matrices: $\sigma^2, \sigma^2 \sigma^a$)

	$\bar{D}^0 (D^{*+})^b$	$D^+ (\bar{D}^{*0})^b$	$(\bar{D}^{*0} \wedge D^{*+})^b$	
$\eta_c (\rho^+)^a$	1	1	1) $\cdot \frac{1}{2} \delta^{ab}$
$\psi^a \pi^+$	1	1	-1	
$(\psi \wedge \rho^+)^a$	1	-1	0	

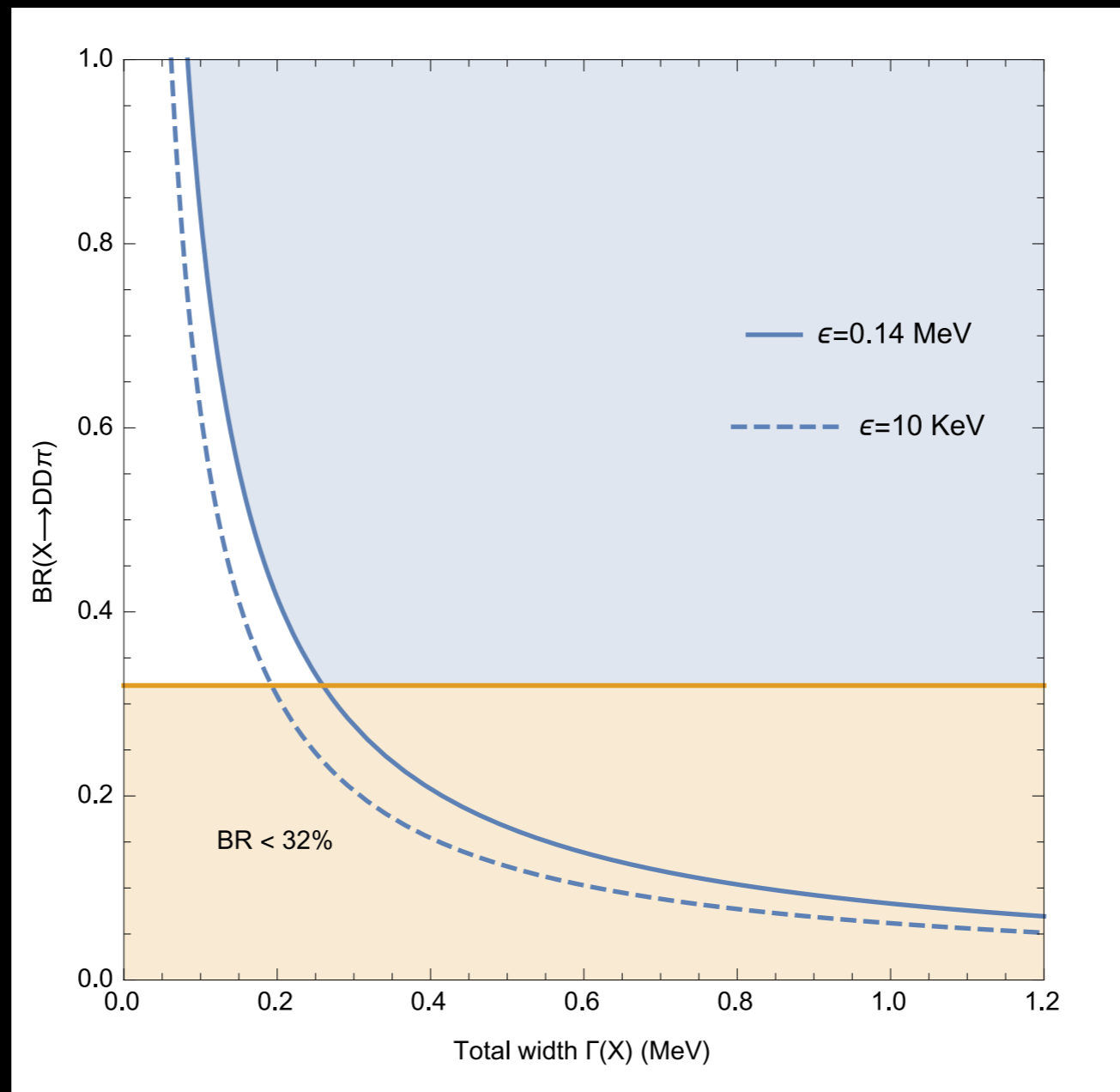
The $J^P = 1^+, G = -1$ state corresponds to

$$i \frac{(\psi \wedge \rho^+)^a}{\sqrt{2}} \pm \frac{\bar{D}^0 D^{*+} - D^+ \bar{D}^{*0}}{\sqrt{2}} \quad (*)$$

Binding energy and decay rates

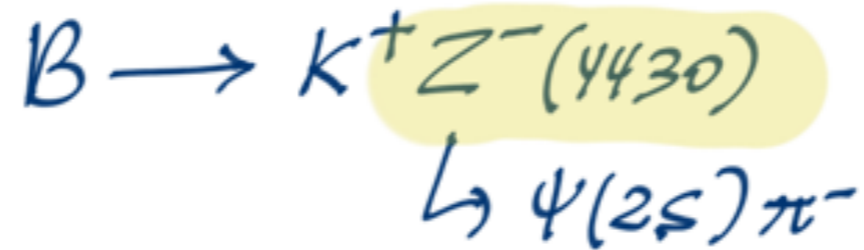
$$B \simeq \frac{G^4}{512 \pi^2} \frac{m^5}{(m_a m_b)^4}$$

$$\mathcal{B}(X \rightarrow DD\pi) \cdot \Gamma(X) \sim G^2 \sim \sqrt{B}$$



CHARGED RESONANCES

LHCb 2014 confirms BELLE 2007 (& disproves Babar 2007)



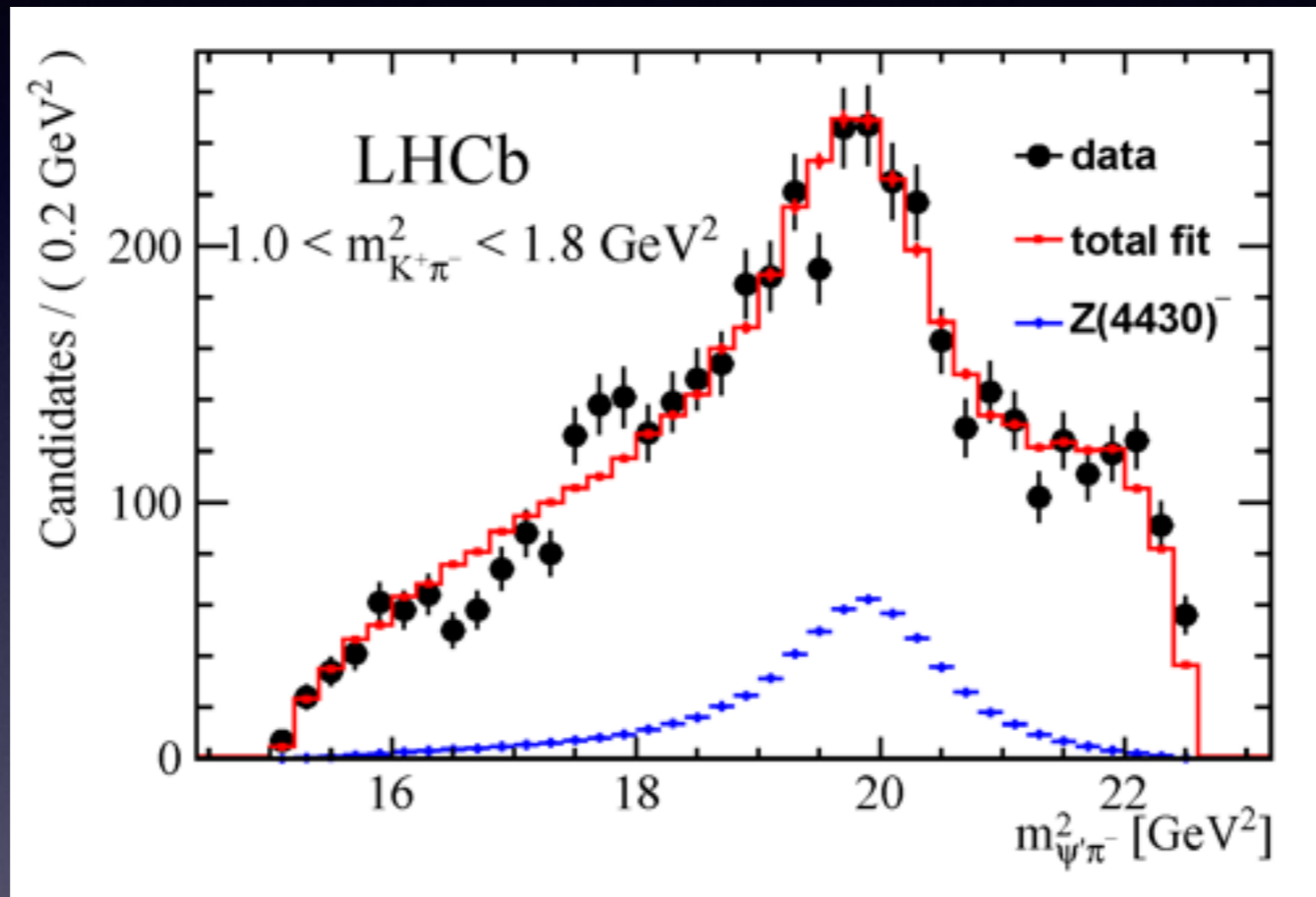
0708.3997 " A CRUCIAL CONSEQUENCE OF Z^- IS A CHARGED STATE IN $J/\psi \pi^\pm$ AT 3880 MeV " (Z^- is its radial excitation)

$$M(\psi(2S)) - M(\psi) \simeq M(Z(4430)) - M(Z(3880))$$

BES III FOUND $Z_c(3900)$ IN 2013.

CHARGED RESONANCES HAVE NOT (YET?) BEEN SEARCHED/OBSERVED IN PP PROMPT COLLISIONS.

Z(4430)⁻ at LHCb | April 2014



$$B \rightarrow K^+ (\psi(2S) \pi^-)_{J^PC = 1^{++}}$$

Signal: 13.9 σ

Other assignments ruled out at 9.7 σ

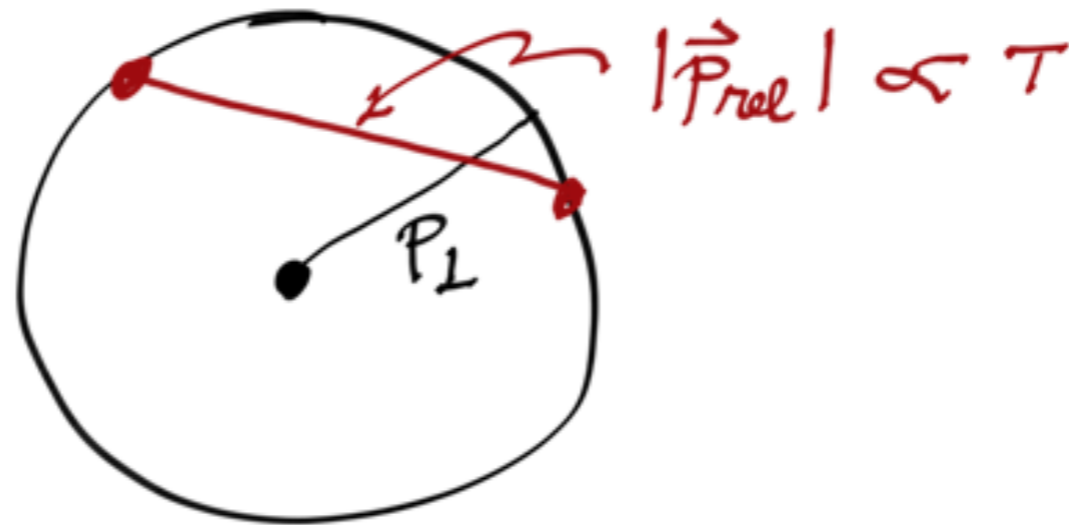
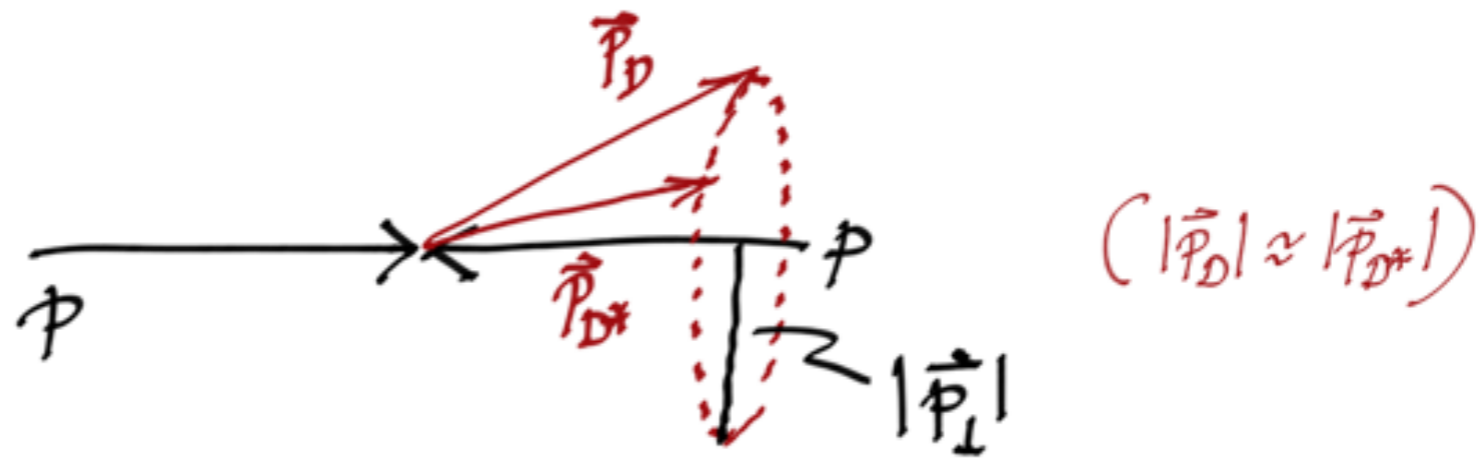
First observed by BELLE in 2007 and not confirmed by BaBar at that time

A $D^0 D^{*0}$ MOLECULE?

The previous arguments rely on $T \simeq 0$

What is T (barycentric energy of DD^* after subtraction of rest masses)

in pp collisions at LHC with HIGH P_{\perp} cuts?

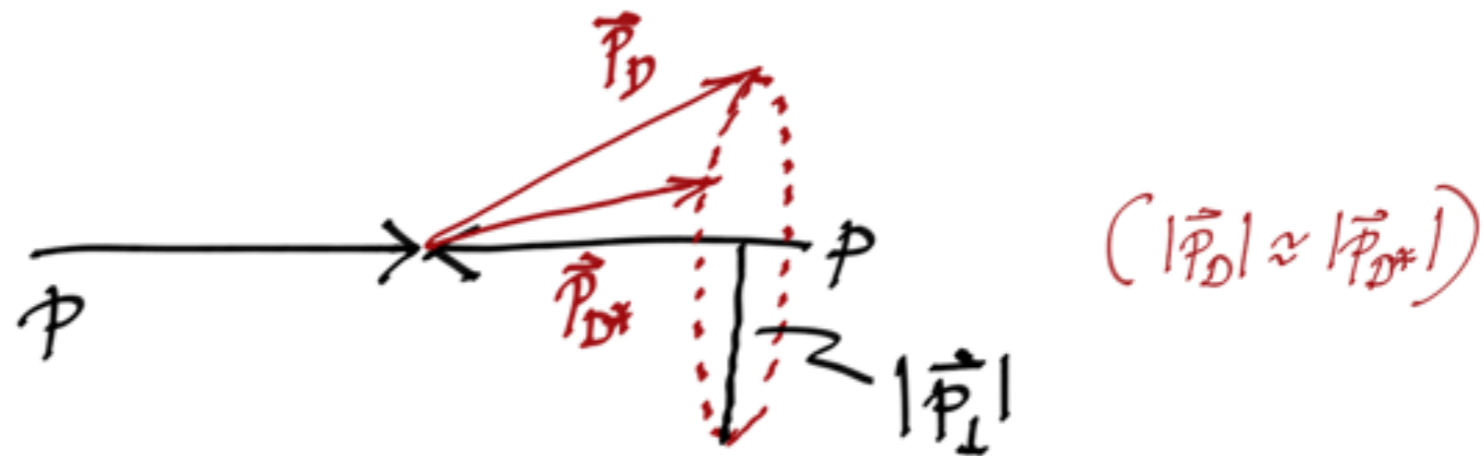


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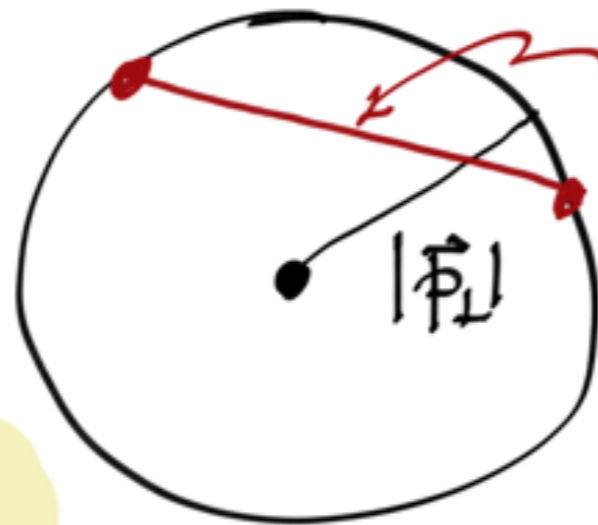
in pp collisions at LHC with HIGH P_{\perp} cuts?



$$|\vec{P}_{rel}| = 1.27 |\vec{P}_{\perp}|$$

$$P(|\vec{P}_{rel}|) =$$

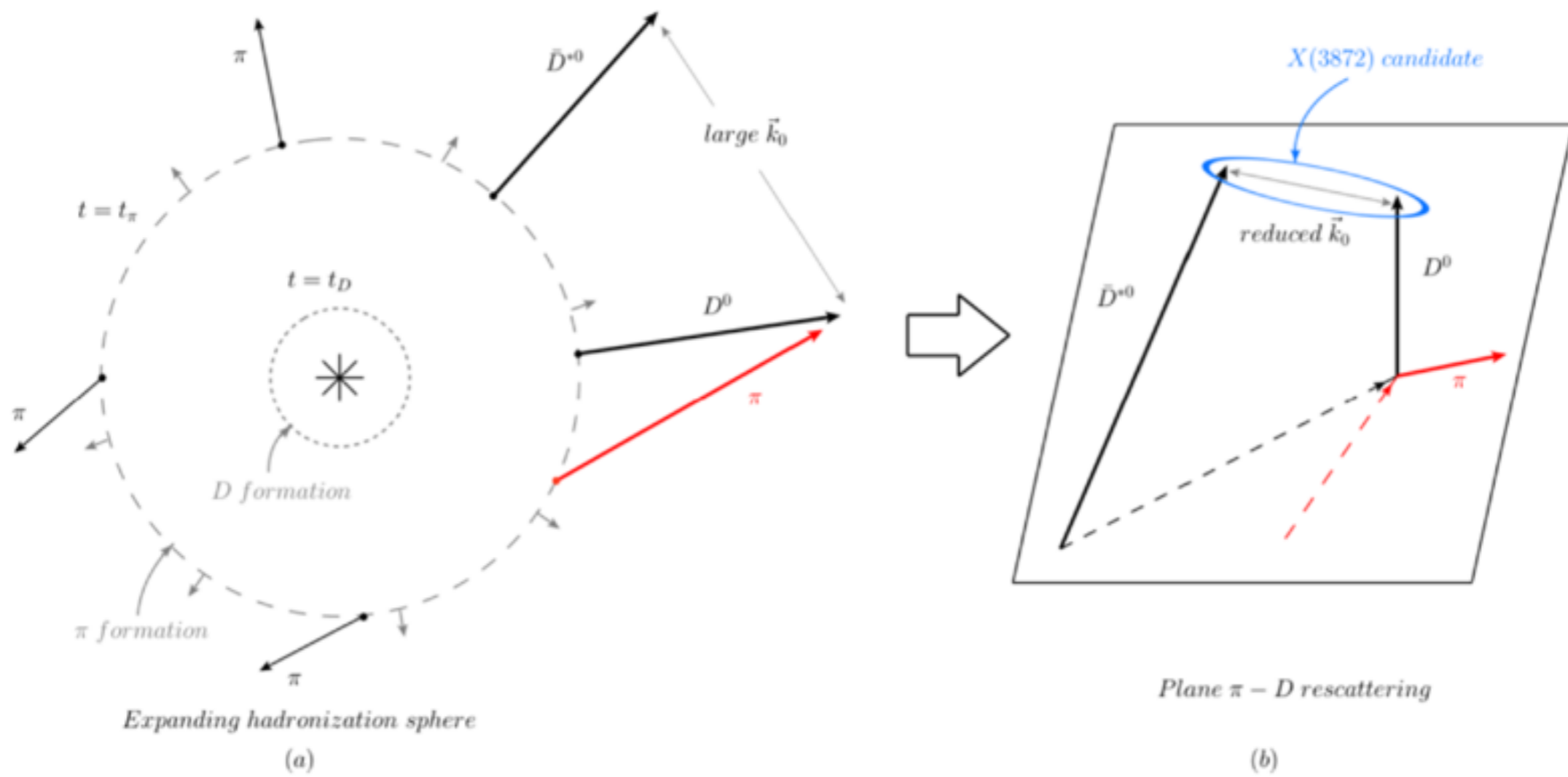
$$= \frac{1/\pi}{\sqrt{|\vec{P}_{\perp}|^2 - (|\vec{P}_{rel}|/2)^2}}$$



$$|\vec{P}_{rel}| \propto T$$

$$P(|\vec{P}_{rel}|) \sim 40\%$$

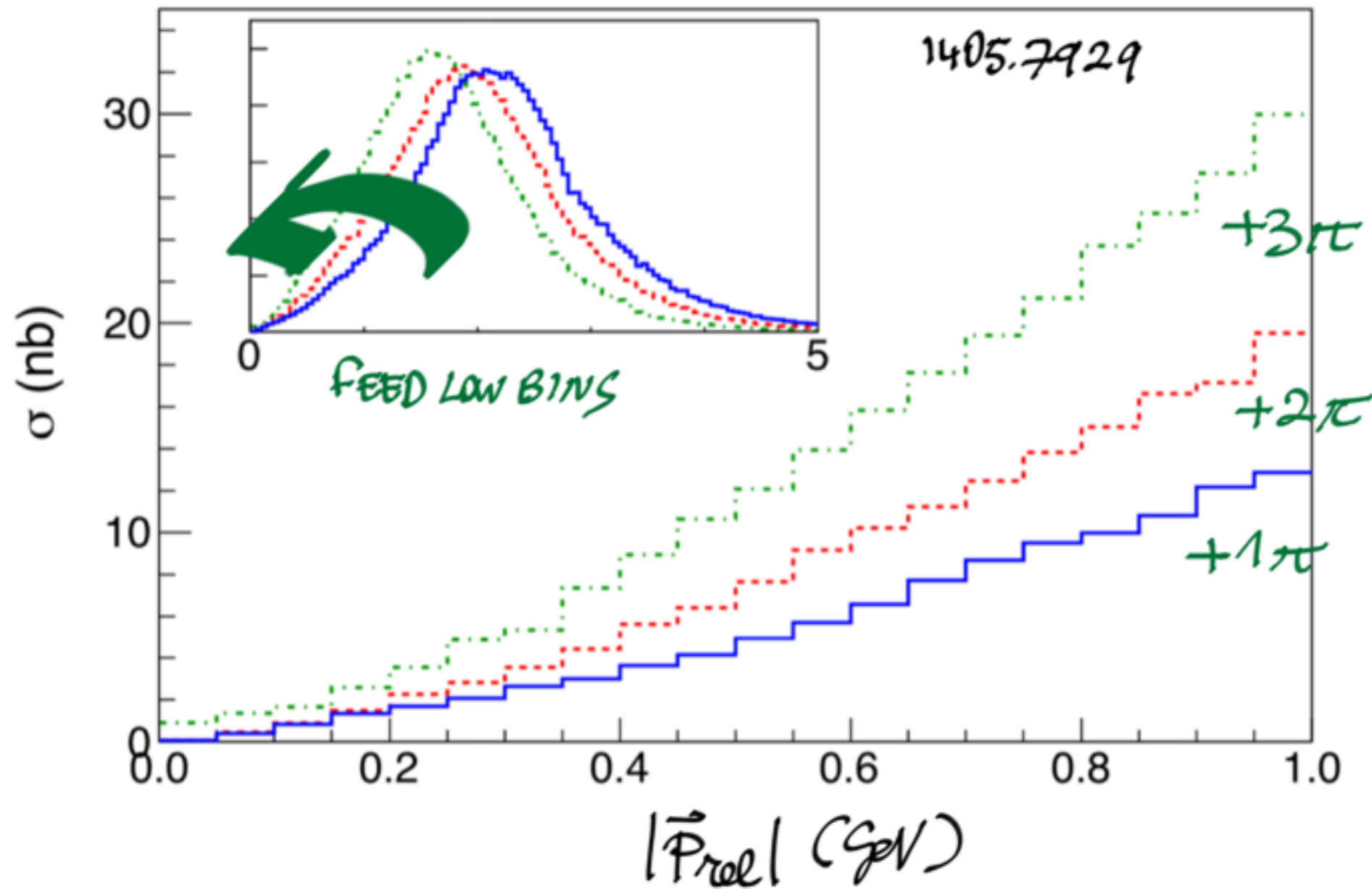
RESCATTERINGS?



RESCATTERINGS WITH HADRONS (π) MIGHT
HELP TO DECREASE $|\vec{P}_{\text{rel}}|$ IN THE DD^* PAIR

A Esposito et al. J.Mod.Phys. 4 (2013) 1569
 A Guerrieri et al. PRD 90 (2014) 034003
 C Bignamini et al PRL 103 (2009) 162001

RESCATTERINGS?



- THE MOST PROBABLE DO* CONFIGURATIONS HAVE HIGH $|\vec{P}_{rel}|$
- THIS IS MORE AND MORE VISIBLE INCREASING THE CUT IN $|\vec{P}_{\perp}|$
- THE FEED-DOWN OBTAINED BY RESCATTERING ON 1,2,3 π IS NEGLIGIBLE,