

Ultra-Fast Silicon Detectors

Hartmut Sadrozinski, Abe Seiden (UCSC)
Nicolo Cartiglia (INFN Torino)

Ultra-Fast Silicon Detectors (UFSD)

provide in the same detector and readout chain

- ultra-fast timing resolution [10's of ps]
- precision location information [10's of μm]

(N.B. a time resolution ≈ 50 ps would already be competitive with SiPM)

2 questions:

- can they work: signal, capacitance, collection time vs. thickness
- will they work: required gain and E-field, fast readout

Disclaimer: data are still coming in, so conclusions and extrapolation are tenuous!

Motivation



- Up to now, semiconductor sensors have supplied precision data only for the 3 space dimensions (diodes, strips, pixels, even “3D”), while the time dimension has had limited accuracy (e.g. to match the beam structure in the accelerator).
- We believe that being able to resolve the time dimension with ps accuracy would open up completely new applications not limited to HEP
- An example in HEP are forward physics projects at the LHC, like the AFP. Scattered protons are tracked from stations 100’s of meters downstream back to the interaction region and the z-vertex is determined from the timing information. So good position resolution and excellent timing are required, which at the moment is done with different detector technologies (pixels and Micro-channel plates).
- Proposal: Combined-function pixel detector will collect electrons from thin n-on-p pixel sensors read out with short shaping time electronics
- Charge multiplication with gain g increases the collected signal
- Need very fast pixel readout

UFSD Pixel Collected Charge

Signal = thickness*EPM

Collection time = thickness/vsat (vsat = 80 μ m/ns)

BackPlane

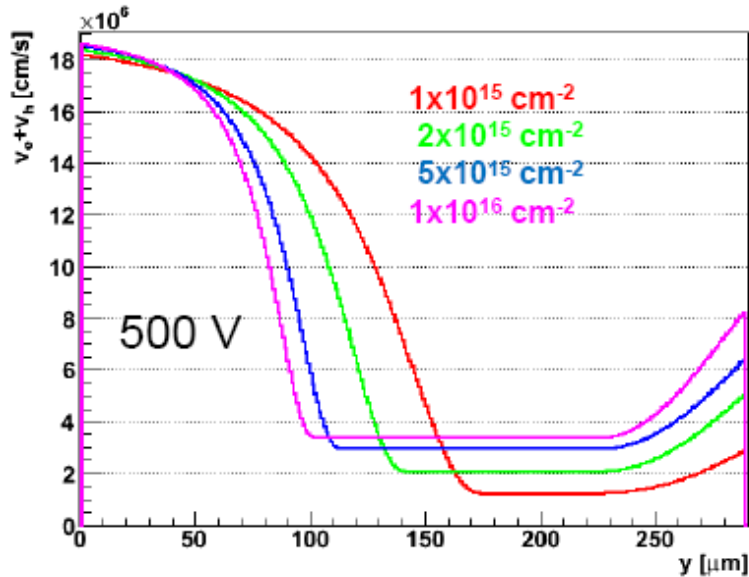
Thickness [μ m]	Capacitance [fF]	Signal [# of e-]	Coll. Time [ps]	Gain req. for 2000 e	
0.1	2500	8.3	1.3	241.0	Realistic gain & cap
1	250	83	12.5	24.1	
2	125	166	25.0	12.0	
5	50	415	62.5	4.8	
10	25	830	125.0	2.4	
20	13	1660	250.0	1.2	Good time resolution
100	2.5	8300	1250.0	0.2	
300	0.8	24900	3750.0	0.1	

For thickness > 5 μ m, Capacitance to the backplane $C_b \ll C_{int}$

For thickness = 2 μ m, $C_b \sim \frac{1}{2}$ of C_{int} , and we might need bipolar (SiGe)?

Fast charge collection: Drift Velocity and E-Field

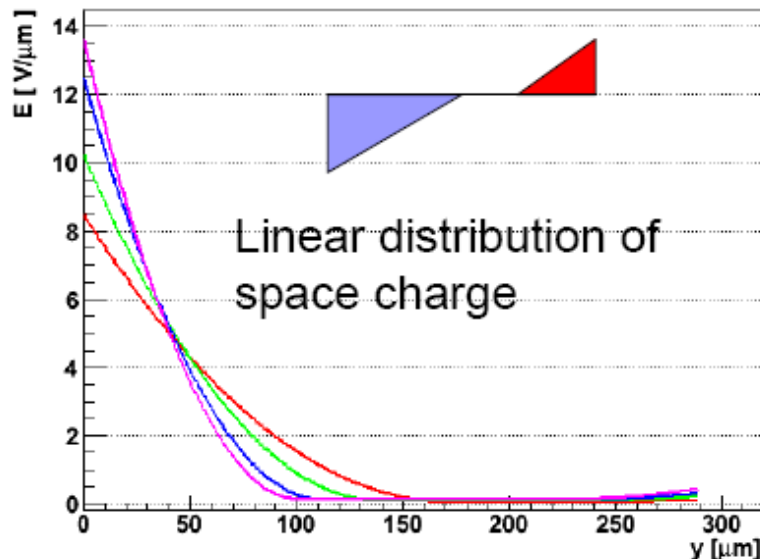
Gregor Kramberger, 19th RD50 Workshop



As long as E-field $E > 20\text{kV/cm}$
drift velocity $v_{\text{drift}} = v_{\text{sat}} \approx 80\mu\text{m/ns}$

Full Depletion Voltage V_{FD} [V]

Res. Rho [kohm-cm]	20	10	5	2
0.01	453	113	28	4.5
0.1	45	11	2.8	0.45
1	4.5	1.1	0.28	0.045
10	0.45	0.11	0.028	0.0045



Bias Voltage for $E_{\text{min}} = 20\text{kV/cm}$

Res. Rho [kohm-cm]	20	10	5	2
0.01		135	40	9
0.1	90	35	14	5
1	50	21	11	4
10	40	20	10	4

Benefits of Gain in Detectors

⊕ Charge multiplication (CM) in silicon sensors (discovered by RD50 institutions) might have applications beyond off-setting charge lost due to trapping during the drift of electrons or holes.

⊕ Charge multiplication makes silicon sensors similar to drift chambers (DC) or Gas Micro-strip Detectors (GMSD), where a modest number of created charges drift to the sense wire, are amplified there (by factors of $> 10^4$) and are then used for fast timing.

⊕ We propose considering silicon detectors for simultaneous precision position and fast timing measurements.

⊕ Recall our experience with DC:

Need to balance the need of high E-field around a wire to have charge multiplication with the need to keep the E-field low to prevent breakdown (wire diameter, field shaping wires,..) and give proper drift field

The drifting electrons contribute mainly to the collected charge after they have undergone charge multiplication, which means that the pulse develops (in principle) in a short time.

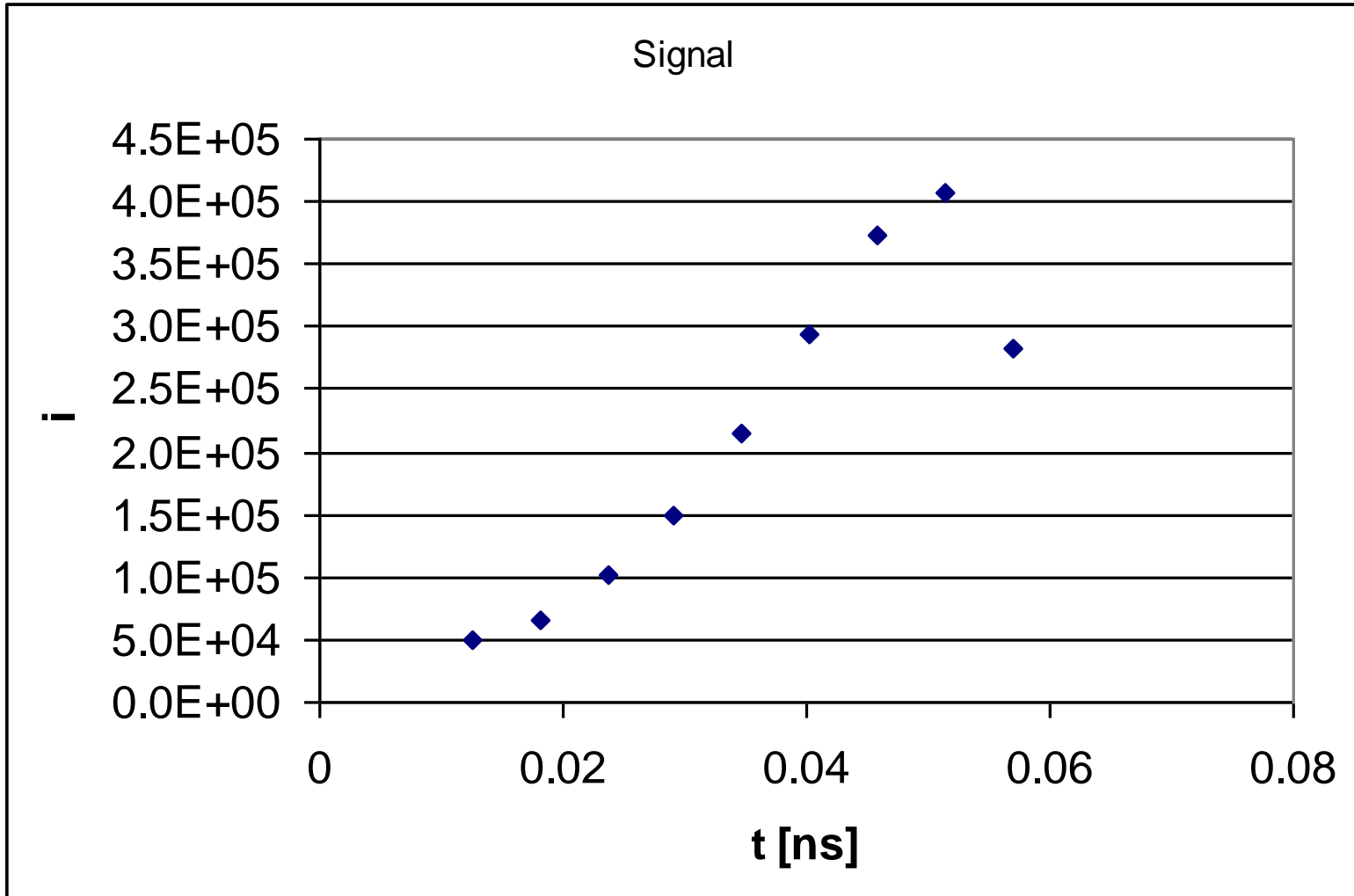
But the very large charge density in the “plasma cloud” prevents the electron to move until the ions have drifted away! So electron signal given by hole dynamics!?

⊕ **Basic question for use of CM**

What field strength do we need?

Can the sensor geometry and doping profile engineered so that the amplification field can be kept high, but just below the breakdown field?

Estimate of pulse shape from Ramo



Timing with falling edge?

Impact ionization - Overstraeten model

➤ In high electric fields electrons and holes acquire enough energy to initiate sizeable charge multiplication

Electrons $> \sim 300$ kV/cm

Holes $> \sim 400$ kV/cm

$$\alpha_{e,h}(\mathcal{E}) = \alpha_{\infty e,h} \exp(-b_{e,h}/|\mathcal{E}|).$$

➤ $300 \text{ kV/cm} < E < 400 \text{ kV/cm}$

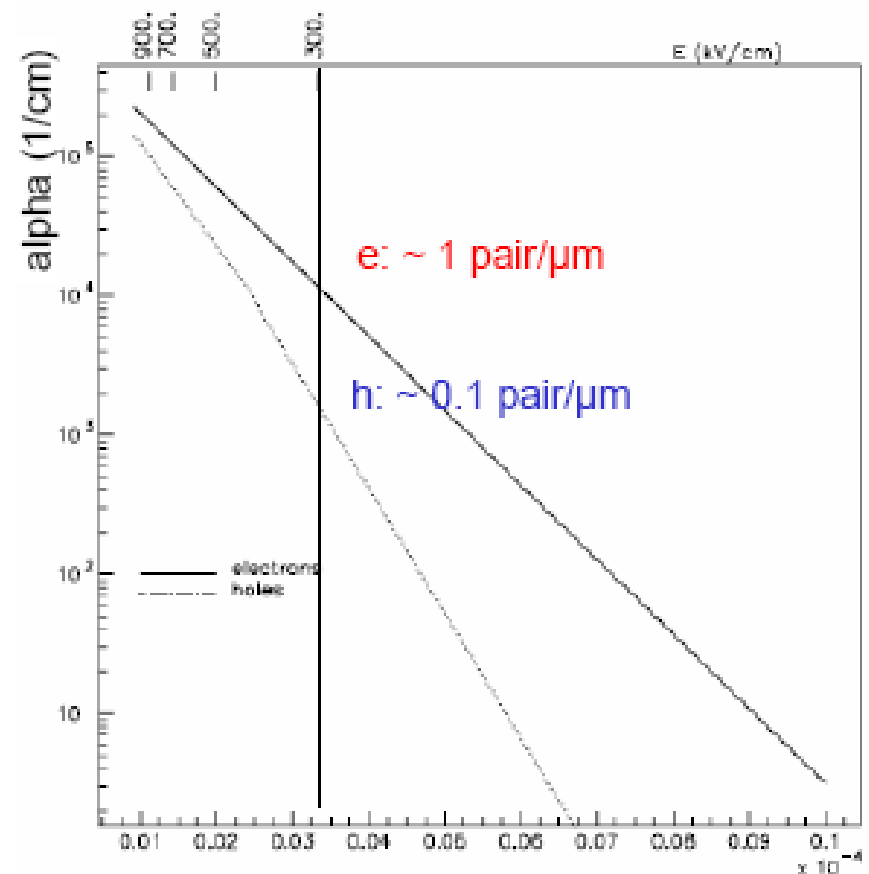
Only electrons ionize

Linear mode

Gain < 10

$N(l) = \exp(\alpha l)$ with l =path length in high electric field

Overstraeten impact ionization model



Anna Macchiolo, 16th RD50 Workshop Barcelona

$1/E$ (cm/V)

“Avalanche” needs distance to develop, but here $l \leq 10 \mu\text{m}$!

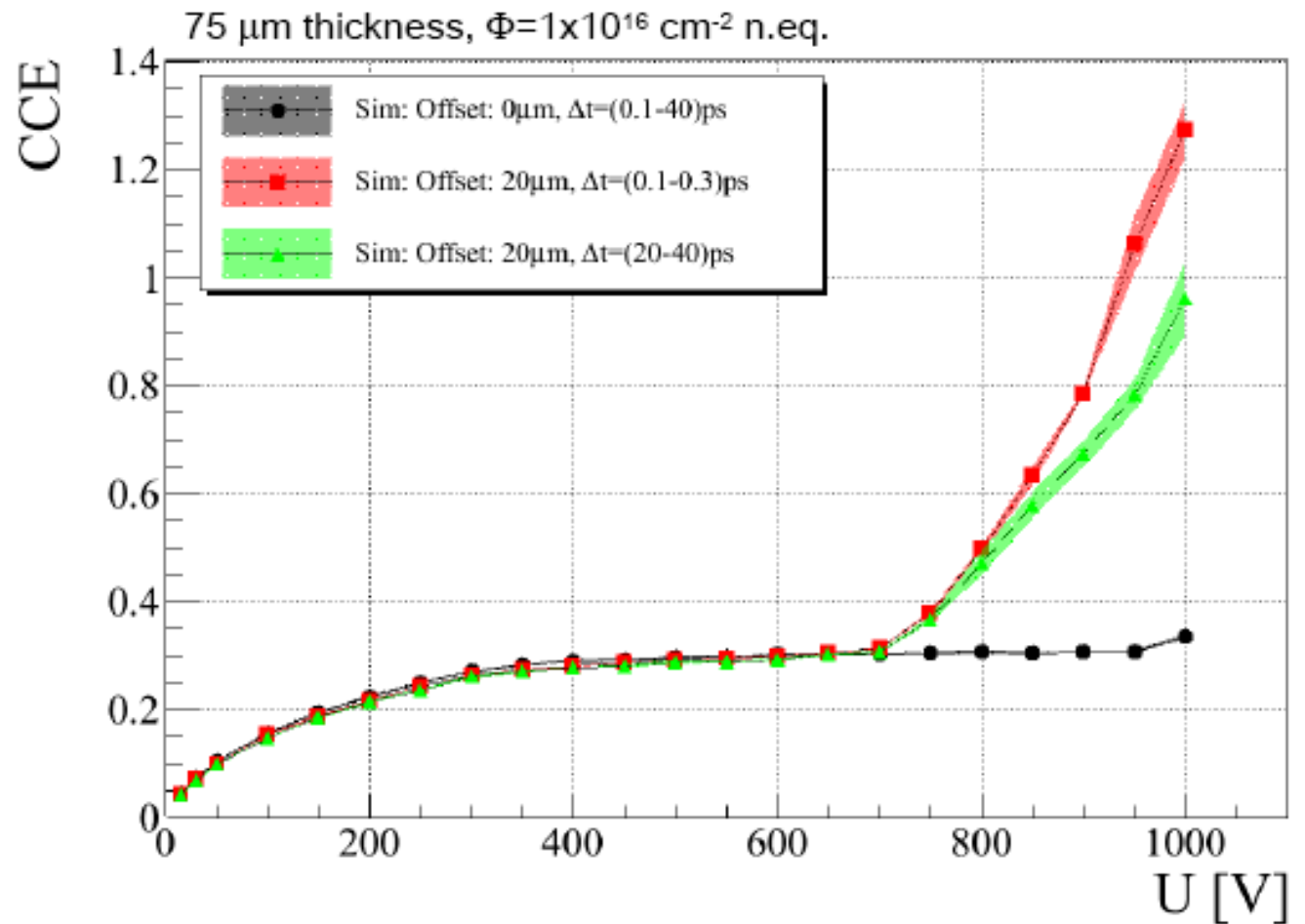
Require max and min E-field close to breakdown field (270kV/cm)

Bias Voltages to reach maximum E-field large over-voltage required (diode)

Res. Rho [kohm-cm]	Thickness [um]							
	20				10			
	VFD	V	E _{max}	E _{min}	VFD	V	E _{max}	E _{min}
0.01	453				113	157	270	44
0.1	45	495	270	225	11	259	270	248
1	4.5	535	270	265	1.1	269	270	268
10	0.45	540	270	270	0.11	270	270	270

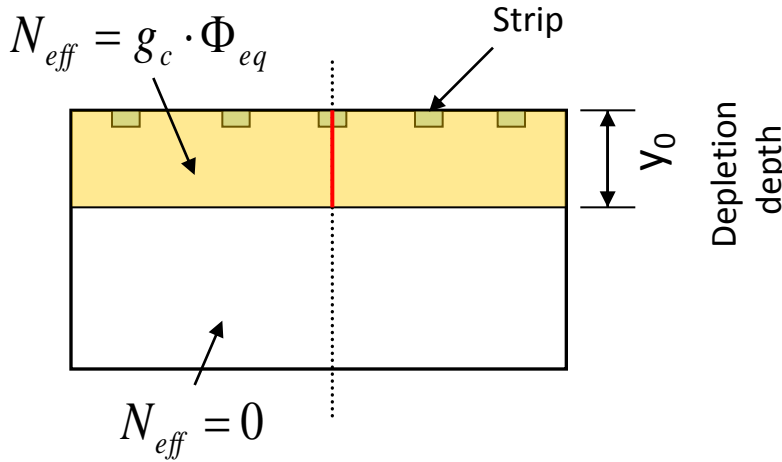
Res. Rho [kohm-cm]	Thickness [um]							
	5				2			
	VFD	V	E _{max}	E _{min}	VFD	V	E _{max}	E _{min}
0.01	28	107	270	157	4.5	50	270	225
0.1	2.8	132	270	258	0.45	54	270	265
1	0.28	135	270	268	0.045	54	270	270
10	0.028	135	270	270	0.0045	54	270	270

E-Field across a pixel



Anna Macchiolo, 16th RD50 Workshop Barcelona

Simplified simulation of strip sensors, real E-field variation not as bad because the hole current injection stabilizes the field but even 30% difference (Elena) is bad



At $V_{bias} < V_{fd}$:

y_0 chosen so that: $V_{bias} = \int \vec{E} d\vec{r}$

Example of electric field:
 $\Phi = 1.6 \cdot 10^{15} \text{ n/cm}^2$, $U = 900 \text{ V}$

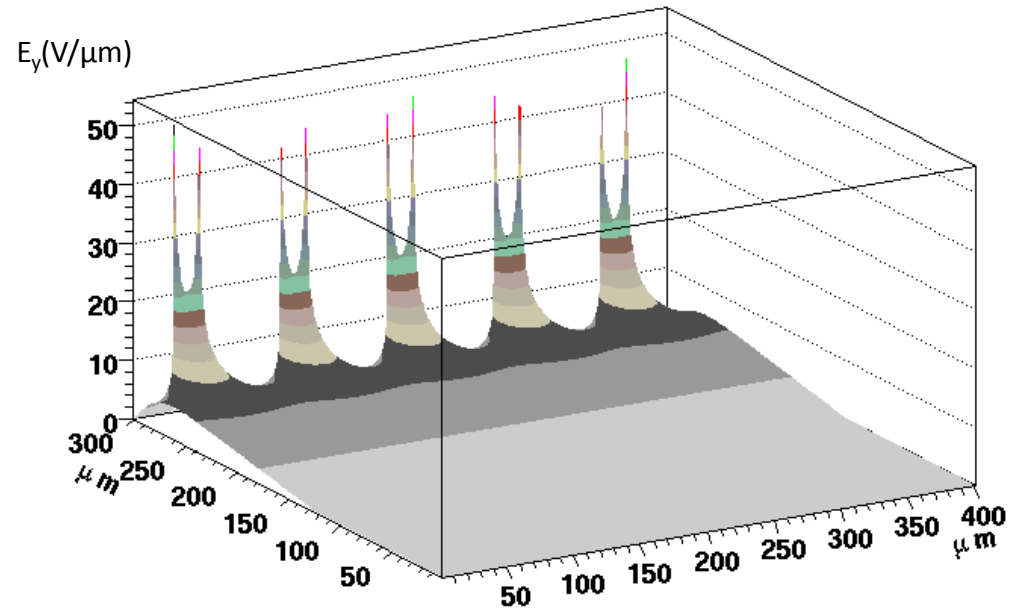
Solve numerically:

$$\nabla^2 V(\vec{r}) = -\frac{e_0 N_{eff}(\vec{r})}{\epsilon_{Si} \epsilon_0}$$

Drift charge through the field,
 get signal using Ramo's theorem:

$$I = q \vec{E}_w(\vec{r}) \cdot \vec{v}(\vec{r})$$

(simulation programs written by G. Kramberger)



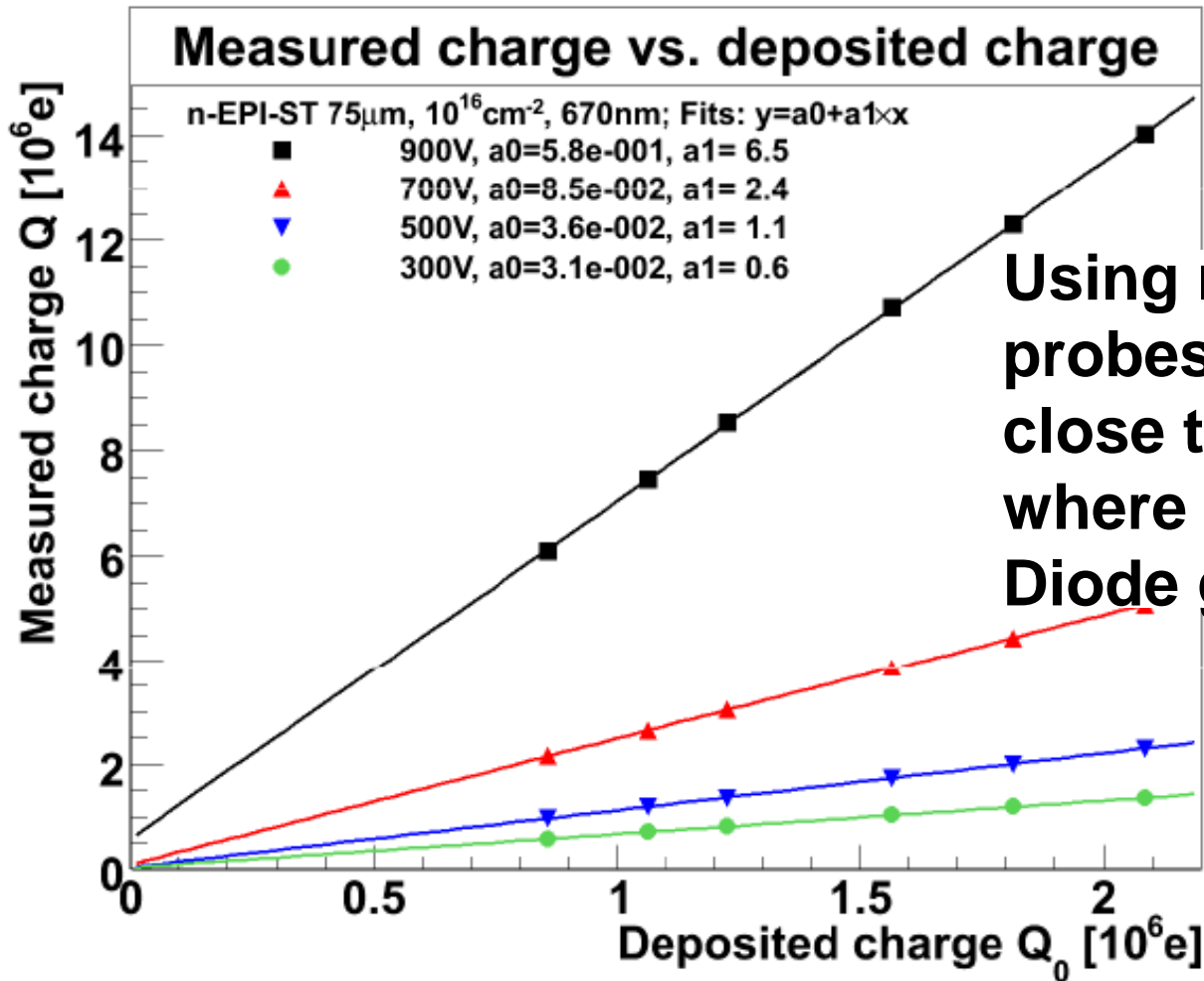
E-Field across a pixel will vary , resulting in different gains in the center and the edges.

Estimated e-Yield vs. E-Field and Thickness

Bias V	E_{max}	E_{min}	Thickness [um]		
			10	5	2
100%	270	270	51k	3k	0.4k
90%	243	243	11k	1.5k	0.3k
80%	217	217	3.2k	0.8k	0.2k

To make use of E-field close to breakdown requires uniformity across pixel

Epi, short drift distances and planar diode gives $g = 6.5$



Using red laser and α 's probes E-field and gain close to the junction, where it counts.
Diode gives uniformity

What about fast readout:

- CERN fixed-target experiment (NA62) needs very fast pixel sensors: Gigatracker (GTK)
- Prototype CFD system (INFN Torino) has ~ 100 ps resolution, predicted to be 30 ps in next iteration.
- Optimized for $200\mu\text{m}$ sensors and hole collection (?), could it be re-designed for electron collection from $2 - 10\mu\text{m}$ sensors?

Summary and more questions:

- Thin pixel sensors with moderate gain could give time resolution between 10 and 100 ps.
- Epi diodes seem perform adequately ($g=6.5$). How can we engineer a pixel sensor to have a uniform field across the pixel like a planar diode? Graded doping concentration? Special deep implants? Trenches?
- How much can we over-deplete pixel sensors to raise E_{\min} ?
- Do we need irradiated sensors, or can we use low-ohmic sensors with field shaping?
- Charge multiplication increases after long annealing (Igor, Marco), could this be employed?
- Start making thin epi strip sensors
- Extend simulations to very thin sensors!
- Need to collaborate with chip designers.

Back-up

E-Field, Depletion Voltage, Resistivity, Thickness

- Assumptions:
 - Desired signal charge 2000 e-
 - Observed Gain:
 - irradiated 3D sensors: $g \approx 3.5$ (Freiburg),
 - irradiated epi diodes $g = 6.6$ (Hamburg)
 - gain > 20 unrealistic?
 - N-on-p pixels (collect electrons)
 - Pixel area $A = 50 \text{ um} \times 50 \text{ um}$
 - Electron yield EPM = 80 e-/um
 - Saturated drift velocity $v_{\text{sat}} = 180 \text{ um/ns}$
 - Breakdown field $E_B \approx 270 \text{ kV/cm}$
 - Full depletion Voltage $V_{\text{FD}} = \text{cons} \cdot d^2 / \rho$
 cons = 0.0113, d in um, rho in kOhm-cm
 - Max Efield at depletion $E_{\text{max}} = 2 \cdot V_{\text{FD}} / d$
 - Max Efield at Bias Voltage V $E_{\text{max}} = (V + V_{\text{FD}}) / d$
 - Min Efield at Bias Voltage V $E_{\text{min}} = (V - V_{\text{FD}}) / d$
 - Capacitance to backplane $C_b = A / d$
 - Capacitance to neighbors $C_{\text{int}} \approx 200 \text{ fF}$