



# Light Sterile Neutrinos in Cosmology



Based on [SG et al., arxiv:1507.08204]

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17<sup>th</sup> Lomonosov Conference, Moscow - August 24, 2015

- 1 Light Sterile Neutrino
  - Oscillations in the 3+1 Neutrino Model
  - Parameterization in Cosmology
- 2 Effects on Cosmology
  - Effects from Relativistic Neutrinos
  - Effects from non-Relativistic Neutrinos
- 3 Constraints on Light Sterile Neutrino Properties
  - CMB Constraints
  - Tensions: CMB vs local measurements
- 4 Open issues

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# Summary on Neutrino Oscillations

Neutrino oscillations: analogous to CKM mixing for quarks, with

$$\nu_\alpha = \sum_{k=1}^3 U_{\alpha k} \nu_k \quad (\alpha = e, \mu, \tau)$$

$\nu_\alpha$  flavour eigenstates,  $U_{\alpha k}$  PMNS mixing matrix,  $\nu_k$  mass eigenstates

Current knowledge of the active  $\nu$  mixing

( $\Delta m_{ji}^2 = m_j^2 - m_i^2$ ,  $\theta_{ij}$  mixing angles):

$$\begin{aligned} \Delta m_{SOL}^2 &= (7.53 \pm 0.18) \cdot 10^{-5} \text{ eV}^2 &&= \Delta m_{21}^2 \\ \Delta m_{ATM}^2 &= (2.44 \pm 0.06) \cdot 10^{-3} \text{ eV}^2(\text{NH}) &&= |\Delta m_{32}^2| \simeq |\Delta m_{31}^2| \\ &= (2.52 \pm 0.07) \cdot 10^{-3} \text{ eV}^2(\text{IH}) \end{aligned}$$

$$\sin^2(2\theta_{12}) = 0.846 \pm 0.021$$

$$\sin^2(2\theta_{23}) = 0.999_{-0.018}^{+0.001}(\text{NH}) - 1.000_{-0.017}^{+0.000}(\text{IH})$$

$$\sin^2(2\theta_{13}) = (9.3 \pm 0.8) \cdot 10^{-2}$$

[PDG - Olive et al. (2014)]

Further details:

[C. Giunti talk]

CP violation through phase  $\delta$  (still unknown) possible only if  $\sin \theta_{13} \neq 0$ .

# Neutrino Oscillation Anomalies

Observed oscillation anomalies in Short BaseLine (SBL) experiments.

A short review: [SG et al., 2015]

- *LSND*: search for  $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ , with  $L/E = 0.4 \div 1.5$  m/MeV. Observed a  $3.8\sigma$  excess of  $\bar{\nu}_e$  events [Aguilar et al., 2001]
- *MiniBooNE*: search for  $\nu_\mu \rightarrow \nu_e$  and  $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ , with  $L/E = 0.2 \div 2.6$  m/MeV. No  $\nu_e$  excess detected, but  $\bar{\nu}_e$  excess observed at  $2.8\sigma$  [MiniBooNE, 2013]
- *Reactor anomaly*: re-evaluation of the expected anti-neutrino flux  $\Rightarrow$  excess of  $\bar{\nu}_e$  events compared to predictions ( $\sim 3\sigma$ ) with  $L < 100$  m [Azabajan et al, 2012]
- *Gallium anomaly*: calibration of GALLEX and SAGE Gallium solar neutrino experiments give a  $2.7\sigma$  anomaly (disappearance of  $\nu_e$ ) [Giunti, Laveder, 2011]

Possible explanation:

Additional squared mass difference  $\Delta m_{\text{SBL}}^2 \simeq 1 \text{ eV}^2$

## 3+1 Neutrino Model

$$\text{SBL anomalies} \Rightarrow \Delta m_{\text{SBL}}^2 \simeq 1 \text{ eV}^2$$



Existence of an additional neutrino degree of freedom,  
mass around 1 eV, no weak interaction  $\Rightarrow$  *light, sterile neutrino* ( $LS\nu$ )



3 active ( $m_i \ll 1 \text{ eV}$ ) + 1 sterile ( $m_s \simeq 1 \text{ eV}$ )  $\nu$  scenario

We must update our mixing paradigm:

$$\nu_\alpha = \sum_{k=1}^{3+1} U_{\alpha k} \nu_k \quad (\alpha = e, \mu, \tau, s)$$

$\nu_s$  is mainly  $\nu_4$ :

$$m_s \simeq m_4 \simeq \sqrt{\Delta m_{41}^2} \simeq \sqrt{\Delta m_{\text{SBL}}^2}$$

Active  $\nu$ :

$$\sum m_{\nu, \text{active}} \simeq 0$$

Sterile  $\nu$ :

$$0.82 \leq m_s^2/\text{eV}^2 \leq 2.19 \quad (3\sigma)$$

[Giunti et al, 2013]

# Relativistic Sterile Neutrino: Effective Number $N_{\text{eff}}$

Radiation energy density  $\rho_r$  in the early Universe:

$$\rho_r = \left[ 1 + \frac{7}{8} \left( \frac{4}{11} \right)^{4/3} N_{\text{eff}} \right] \rho_\gamma = [1 + 0.2271 N_{\text{eff}}] \rho_\gamma$$

$\rho_\gamma$  photon energy density,  $7/8$  is for fermions,  $(4/11)^{4/3}$  due to photon reheating after neutrino decoupling

- $N_{\text{eff}} \rightarrow$  all the radiation contribution not given by photons
- $N_{\text{eff}} \simeq 1$  correspond to a single family of active neutrino, in equilibrium in the early Universe
- Active neutrinos:  $N_{\text{eff}} = 3.046$  [Mangano et al., 2005] due to not instantaneous decoupling for the neutrinos
- additional  $LS\nu$  contributes with  $\Delta N_{\text{eff}} = N_{\text{eff}} - 3.046$ :

$$\Delta N_{\text{eff}} = \frac{\rho_s^{\text{rel}}}{\rho_\nu} = \left[ \frac{7}{8} \frac{\pi^2}{15} T_\nu^4 \right]^{-1} \frac{1}{\pi^2} \int dp p^3 f_s(p) \quad [\text{Acero et al., 2009}]$$

$\rho_\nu$  energy density for one active neutrino species,  $\rho_s^{\text{rel}}$  energy density of  $LS\nu$  when relativistic,

$p$  neutrino momentum,  $f_s(p)$  momentum distribution,  $T_\nu = (4/11)^{1/3} T_\gamma$

# Non-Relativistic Sterile Neutrino: Effective Mass $m_s^{\text{eff}}$

$m_s \simeq 1 \text{ eV} \rightarrow \nu_s$  is non-relativistic today ( $T_\nu \propto 10^{-4} \text{ eV}$ )

LS $\nu$  density parameter today:

$$\omega_s = \Omega_s h^2 = \frac{\rho_s}{\rho_c} h^2 = \frac{h^2}{\rho_c} \frac{m_s}{\pi^2} \int dp p^2 f_s(p) \quad [\text{Acero et al., 2009}]$$

$\rho_s$  energy density of non-relativistic LS $\nu$ ,  $\rho_c$  critical density and  $h$  reduced Hubble parameter

Alternatively:

$$m_s^{\text{eff}} = 94.1 \text{ eV} \omega_s \quad [\text{Planck 2013 Results, XVI}]$$

The factor (94.1 eV) is the same for the active neutrinos:

$$\omega_{\nu, \text{active}} = \sum_{\text{active}} m_\nu / (94.1 \text{ eV})$$

$$\text{If } f_s(p) = f_{\text{active}}(p), m_s^{\text{eff}} \equiv m_s$$



## The momentum distribution function $f_s(p)$

$\Delta N_{\text{eff}}$ ,  $m_s^{\text{eff}}$  depend on the momentum distribution function  $f_s(p)$ .

LS $\nu$  relativistic at decoupling  $\Rightarrow f_s(p)$  independent of  $m_s$ .

Active neutrinos decoupled at  $T_\nu \simeq 1$  MeV. Is the same for LS $\nu$ ?

Oscillations + sterile  $\Rightarrow$  LS $\nu$  decouples not later than active neutrinos.

### Production mechanism?

Thermal production (TH):

temperature  $T_s = \alpha T_\nu$

$$f_s(p) = \frac{1}{e^{p/T_s} + 1}$$

$\Downarrow$

$$\Delta N_{\text{eff}} = \alpha^4$$

$$\omega_s = \alpha^3 m_s / (94.1 \text{ eV})$$

$\Downarrow$

$$m_s^{\text{eff}} = \alpha^3 m_s = \Delta N_{\text{eff}}^{3/4} m_s$$

Non-thermal production:

[Dodelson, Widrow 1993] (DW) model

$$f_s(p) = \frac{\beta}{e^{p/T_\nu} + 1}$$

$\Downarrow$

$$\Delta N_{\text{eff}} = \beta$$

$$\omega_s = \beta m_s / (94.1 \text{ eV})$$

$\Downarrow$

$$m_s^{\text{eff}} = \beta m_s = \Delta N_{\text{eff}} m_s$$

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# Additional Radiation in the Early Universe

$$\rho_r = [1 + 0.2271 N_{\text{eff}}] \rho_\gamma$$

$$H^2 = 8\pi G \rho_T / 3$$

$N_{\text{eff}}$  controls the expansion rate  $H$  in the early Universe, during radiation dominated phase

influence on

Big Bang Nucleosynthesis:  
production of light nuclei

matter-radiation equality

expansion rate at  
CMB decoupling

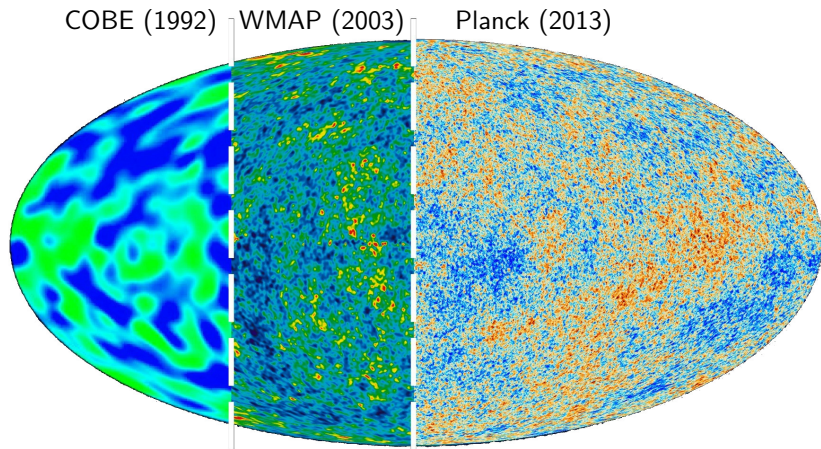
CMB+BBN+ $^4\text{He}$ ,  $D$  abundances:

$\Delta N_{\text{eff}} \lesssim 0.2$  at 95% CL

[Cyburt et al., 2015]

# Cosmic Microwave Background (CMB)

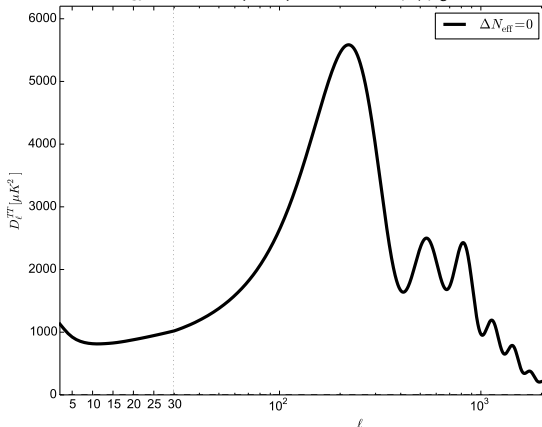
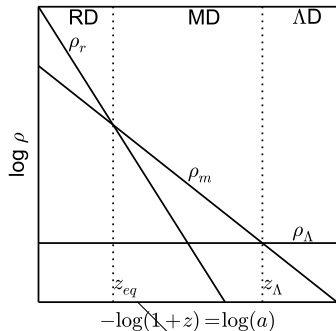
Anisotropies at the level of  $10^{-5}$ : very high precision measurements are needed.  
Improvement of the CMB experiments in 20 years:



# Additional Radiation: Effects on the CMB

Starting configuration:

RD: Radiation Dominated, MD: Matter Dominated,  $\Lambda$ D: Dark Energy Dominated;  $(1+z) = a^{-1}$ ;  $\omega_i = \rho_i / \rho_C$

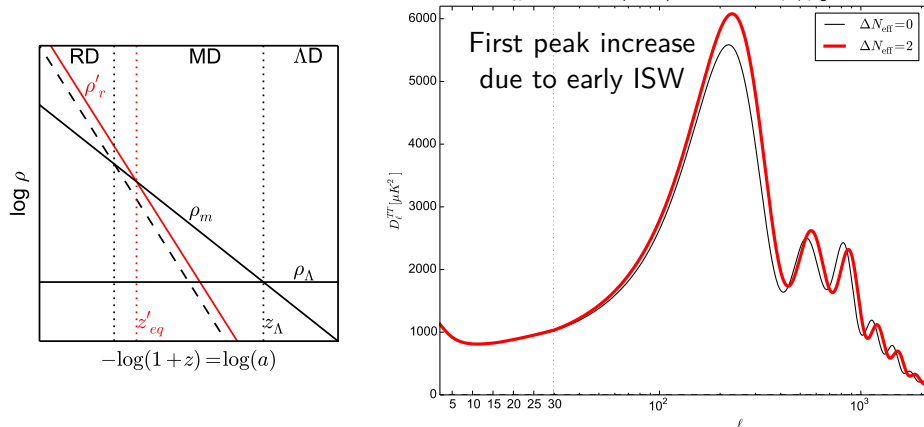


$$1 + z_{eq} = \frac{\omega_m}{\omega_r} = \frac{\omega_m}{\omega_\gamma} \frac{1}{1 + 0.2271 N_{eff}}$$

# Additional Radiation: Effects on the CMB

If we increase  $N_{\text{eff}}$ , all the other parameters fixed:

RD: Radiation Dominated, MD: Matter Dominated,  $\Lambda$ D: Dark Energy Dominated;  $(1+z) = a^{-1}$ ;  $\omega_i = \rho_i/\rho_C$

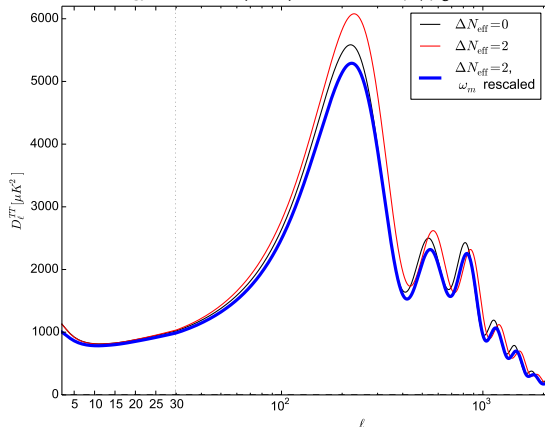
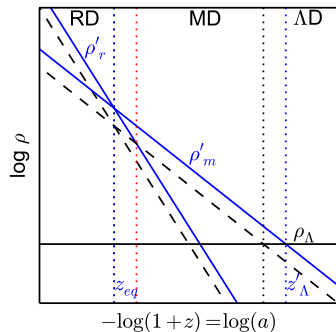


At  $z_{\text{CMB}}$ : higher  $H \propto \rho_r \Rightarrow$  smaller comoving sound horizon  $r_s \propto H^{-1}$   
 $\Rightarrow$  decrease of the angular scale of the acoustic peaks  $\theta_s = r_s/D_A$   
 $\Rightarrow$  shift of the peaks at higher  $\ell$

# Additional Radiation: Effects on the CMB

If we increase  $N_{\text{eff}}$ , plus  $\omega_m$  to fix  $z_{\text{eq}}$ :

RD: Radiation Dominated, MD: Matter Dominated,  $\Lambda$ D: Dark Energy Dominated;  $(1+z) = a^{-1}$ ;  $\omega_i = \rho_i / \rho_C$

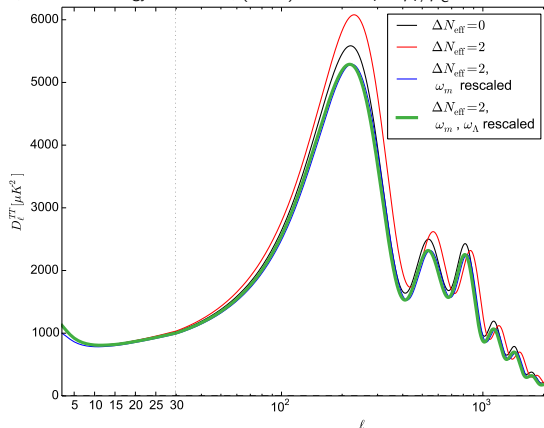
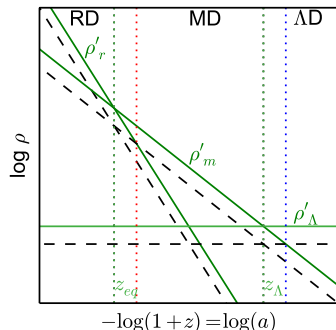


- Contribution from early ISW effect restored (first peak)
- different slope of the Sachs-Wolfe plateau, peak positions, envelope of high- $\ell$  peaks  $\Rightarrow$  due to later  $z_{\Lambda}$

# Additional Radiation: Effects on the CMB

If we increase  $N_{\text{eff}}$ , plus  $\omega_m, \omega_\Lambda$  to fix  $z_{\text{eq}}, z_\Lambda$ :

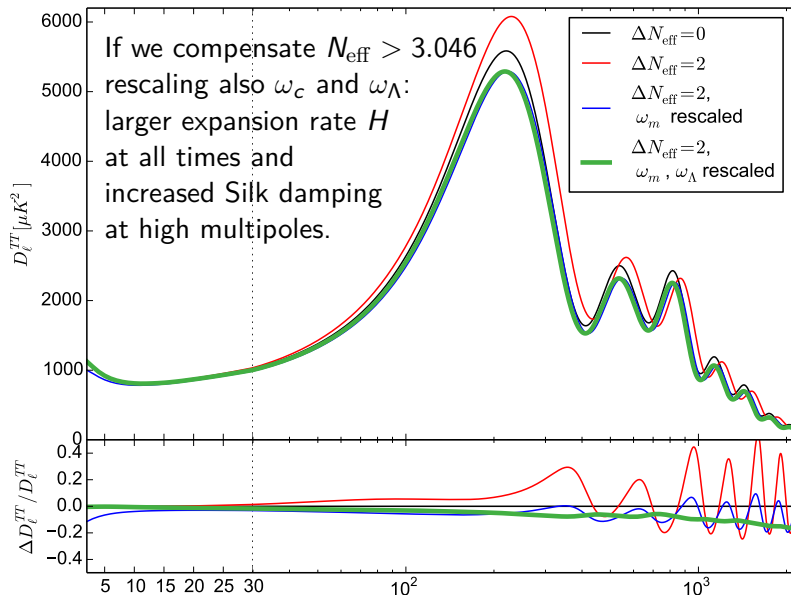
RD: Radiation Dominated, MD: Matter Dominated,  $\Lambda$ D: Dark Energy Dominated;  $(1+z) = a^{-1}$ ;  $\omega_i = \rho_i / \rho_C$



- peak positions recovered;
- slope of the Sachs-Wolfe plateau recovered;
- peak amplitude not recovered!



# Additional Radiation: Effects on the CMB



# Impact of additional non-relativistic neutrinos on the CMB

$\omega_m = \omega_b + \omega_c$  in the early Universe if  $LS\nu$  is relativistic



$1 + z_{\text{eq}} = (\omega_b + \omega_c)/\omega_r$   
independent of  $m_s$

$\omega_m^0 = \omega_b^0 + \omega_c^0 + \omega_\nu^0$  today

$LS\nu$  relativistic at recombination  
affects late time evolution only

Effects can be compensated with a variation in  $\Omega_\Lambda$

Degeneracies with spatial curvature are also possible

small effects on the SW plateau (cosmic variance!)

CMB is not the best direction

# Free-streaming - I

Massive neutrino



damping in the perturbations due to free-streaming length  $\lambda_{FS}$

Relativistic neutrinos

velocity  $v_s \simeq c$

$$\lambda_{FS}/a \propto (aH)^{-1} \propto t^{1/3} \text{ (MD)}$$

Non-relativistic neutrinos

$$\langle v_s \rangle = \frac{\int p^2 dp f(p) p/m_s}{\int p^2 dp f(p)} \propto \frac{\Delta N_{\text{eff}}}{\omega_s}$$

$$\lambda_{FS}/a \propto (a^2 H)^{-1} \propto t^{-1/3} \text{ (MD)}$$

⇒ Maximum  $\lambda_{FS}/a$  at the time of non-relativistic transition.



$$\text{Corresponds to } k_{\text{nr}} \simeq 0.0178 \Omega_m^{1/2} \left( \frac{T_\nu}{T_s} \right)^{1/2} \left( \frac{m_s}{1 \text{ eV}} \right)^{1/2} h \text{ Mpc}^{-1}$$

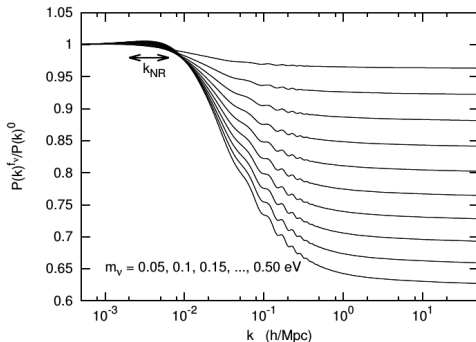
# Free-streaming - II

Damping occurs for all  $k \gtrsim k_{\text{nr}}$

[Neutrino Cosmology, Lesgourgues et al.]  
(fixed  $h, \omega_m, \omega_b, \omega_\Lambda$ )

Plot:  $\frac{P_{m_\nu > 0}(k)}{P_{m_\nu = 0}(k)}$

- top to bottom:  $m_\nu = 0.05$  eV to  $m_\nu = 0.5$  eV
- $\Delta P/P \simeq -7(8)(m_\nu/1 \text{ eV})$



Expected constraints from future surveys:

- Planck CMB + DES:  $\sigma(m_\nu) \simeq 0.04\text{--}0.06$  eV [Font-Ribera et al., 2014]
- Planck CMB + Euclid:  $\sigma(m_\nu) \simeq 0.03$  eV [Audren et al., 2013]

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# CMB and SBL joint constraints

CMB:

(no SBL)

+SBL (DW)

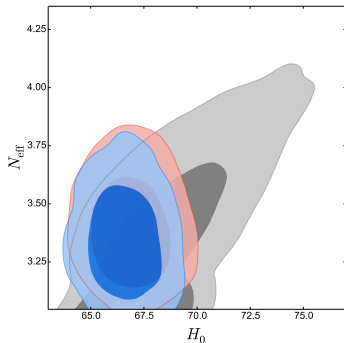
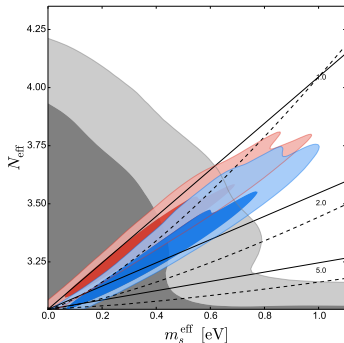
+SBL (TH)

solid lines: (DW)

$$m_s^{\text{eff}} = m_s \Delta N_{\text{eff}}$$

dashed lines: (TH)

$$m_s^{\text{eff}} = m_s (\Delta N_{\text{eff}})^{3/4}$$



- $\nu_s$  as Warm Dark Matter (WDM):  $N_{\text{eff}} \simeq 3.046$ , large  $m_s^{\text{eff}}$  (large  $m_s$ );
- SBL prior:  $m_s \simeq 1.2$  eV, but

$N_{\text{eff}} \simeq 4$  ( $\nu_s$  thermalized as  $\nu_{\text{SM}}$ ) disfavoured;

- (DW), (TH) models give similar results ( $N_{\text{eff}}$  slightly higher in (DW)).

## $\sigma_8$ Tension: cluster counts, weak lensing

Cosmic shear: distortion of distant galaxy images by gravitational lensing of LSS  
 $\Rightarrow$  sensitive to non-linear matter density along the line of sight/amplitude of matter power spectrum.

Assuming  $\Lambda$ CDM model:

$\sigma_8$ : rms fluctuation in total matter (baryons + CDM + neutrinos) in  $8h^{-1}$  Mpc spheres, today;  
 $\Omega_m$ : total matter density today divided by the critical density

CFHTLenS weak lensing data alone  
[Heymans et al., 2013] (68% CL):

$$\sigma_8(\Omega_m/0.27)^{0.46 \pm 0.02} = 0.774 \pm 0.04$$

CMB results

[Planck 2013] (68% CL):

$$\sigma_8(\Omega_m/0.27)^{0.46} = 0.89 \pm 0.03$$

2 $\sigma$  discrepancy!

Similar results from cluster counts:

Planck SZ Cluster Counts

[Planck 2013 Results XX] (68% CL):

$$\sigma_8(\Omega_m/0.27)^{0.3} = 0.764 \pm 0.025$$

Planck + WMAP pol + ACT/SPT

[Planck 2013] (68% CL):

$$\sigma_8(\Omega_m/0.27)^{0.3} = 0.87 \pm 0.02$$

3 $\sigma$  discrepancy!

Qualitatively similar results from *SPT* clusters, *Chandra Cluster Cosmology Project*.

Solving  $\sigma_8$  Tension with the sterile neutrinoCMB+ $H_0$ +BAO+LSS:

(no SBL)

+SBL (DW)

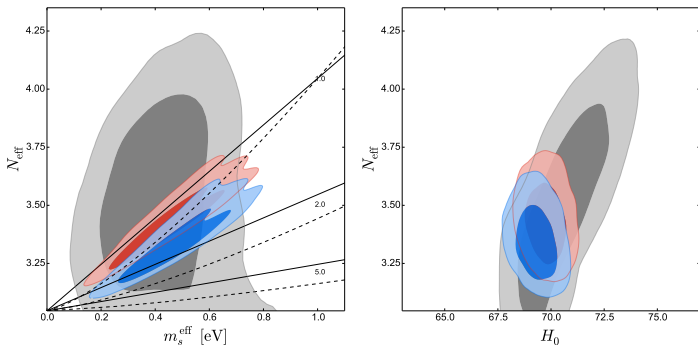
+SBL (TH)

solid lines: (DW)

$$m_s^{\text{eff}} = m_s \Delta N_{\text{eff}}$$

dashed lines: (TH)

$$m_s^{\text{eff}} = m_s (\Delta N_{\text{eff}})^{3/4}$$



- LSS results give preference towards non-zero  $m_s^{\text{eff}} \rightarrow$  non-zero  $m_s$ : smaller  $\sigma_8$  from LSS can be addressed with massive  $\nu_s$  (due to free streaming);
- no SBL prior:  $N_{\text{eff}}$  constraints almost unchanged;
- with SBL prior: preference for  $N_{\text{eff}} > 3.046$  at more than  $2\sigma$ ;
- with SBL prior:  $N_{\text{eff}} \simeq 4$  still hardly disfavoured



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# The Hubble Parameter $H_0$ Tension

[Cuesta et al., 2014] 68% CL error bars

Hubble parameter today:

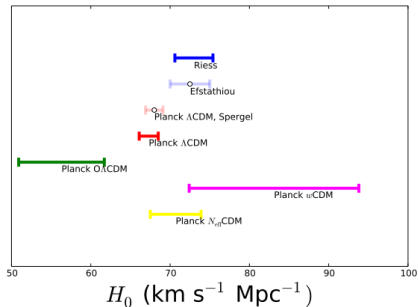
$$v = H_0 d, \text{ with}$$

$$H_0 = H(z = 0)$$

Local measurements:

$H(z = 0)$ , local and independent on evolution

(model independent, systematics?)



CMB measurements

(probe  $z \simeq 1100$ ):

$H_0$  from the cosmological evolution

(model dependent, well controlled systematics)

(HST Cepheids)

[Riess et al., 2011] (SNe Ia calibrated distance):

$$H_0 = 73.8 \pm 2.4 \text{ Km s}^{-1} \text{ Mpc}^{-1}$$

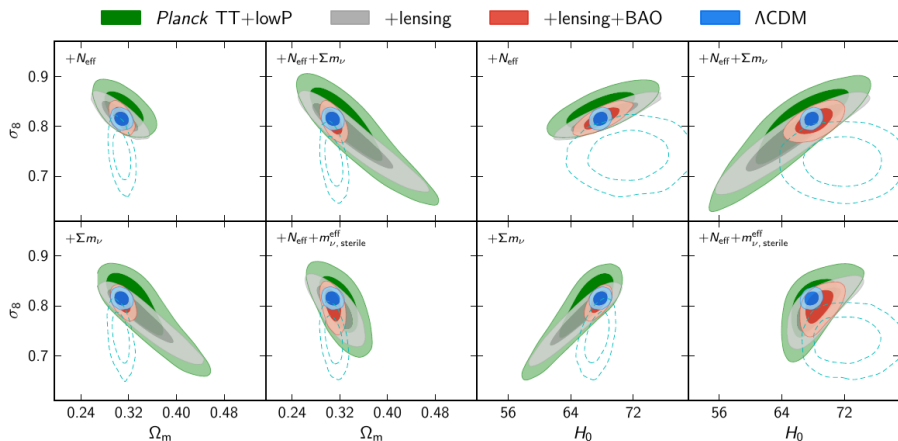
[Efstathiou 2013] (NGC 4258 calibrated distance):

$$H_0 = 70.6 \pm 3.3 \text{ Km s}^{-1} \text{ Mpc}^{-1}$$

( $\Lambda$ CDM - CMB data only)

$$[\text{Planck 2013}]: H_0 = 67.3 \pm 1.2 \text{ Km s}^{-1} \text{ Mpc}^{-1}$$

$$[\text{Planck 2015}]: H_0 = 67.27 \pm 0.66 \text{ Km s}^{-1} \text{ Mpc}^{-1}$$

Solving both  $\sigma_8$  and  $H_0$  Tension?

$H_0$  increases  $\Rightarrow \sigma_8$  increases (and viceversa)!

The correlations do not help.

# Incomplete Thermalization

Many probes constrain  $\Delta N_{\text{eff}} < 1$ . Do we need

- a mechanism to suppress oscillations and full thermalization of  $\nu_s$ ?
- to compensate  $\Delta N_{\text{eff}} = 1$  with additional mechanisms in Cosmology?

Some ideas: (see references in [SG et al., arxiv:1507.08204])

- large lepton asymmetry [Foot et al., 1995; Mirizzi et al., 2012; many more]
- new neutrino interactions [Bento et al., 2001; Dasgupta et al., 2014; Hannestad et al., 2014; Saviano et al., 2014; many more]
- entropy production after neutrino decoupling [Ho et al., 2013]
- very low reheating temperature [Gelmini et al., 2004; Smirnov et al., 2006]
- time varying dark energy components [Giusarma et al., 2012]
- larger expansion rate at the time of  $\nu_s$  production [Rehagen et al., 2014]
- freedom in the Primordial Power Spectrum (PPS) of scalar perturbations from inflation compensate damping due to  $N_{\text{eff}} \neq 3.046$  [SG et al., 2015]

# Conclusions

- Short BaseLine (SBL) oscillations suggest the presence of an additional neutrino, sterile, with  $m_s \simeq m_4 \simeq 1$  eV;
- $\nu_s$  can have measurable effects on cosmological observables;
- CMB measurements give strong constraints on  $\nu_s$  properties;
- tension between CMB observations and local observations;
  - ▶ unaccounted systematics?
  - ▶ new physics?
- sterile neutrinos suggested by SBL oscillation anomalies can help solving the tensions,
  - ▶ but problems in producing them with small  $\Delta N_{\text{eff}}$  (preferred by cosmology);
- further investigation and/or new ideas needed!

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Thank you for the attention