



# Stefano Gariazzo

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## Neutrino mass ordering: current status

*What Bayesian model comparison  
and available data can tell us*

- 1 Basics of Bayesian statistics
  - Bayes' theorem
  - Bayesian model comparison
- 2 Constraining the neutrino mass ordering
  - Introducing the problem
  - Comparing models and mass orderings
- 3 Constraining the neutrino masses
- 4 Conclusions

## 1 Basics of Bayesian statistics

- Bayes' theorem
- Bayesian model comparison

## 2 Constraining the neutrino mass ordering

- Introducing the problem
- Comparing models and mass orderings

## 3 Constraining the neutrino masses

## 4 Conclusions

## Bayes' theorem

Basic rule to deal with Bayesian probability!

given hypothesis  $H$ , data  $d$ , some information  $I$  (true):

Bayes theorem:

$$p(H|d, I) = \frac{p(d|H, I) p(H|I)}{p(d|I)}$$

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sampling distribution of  
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or "Bayesian evidence",

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model comparison

## Bayesian evidence

“Bayesian evidence” or “Marginal likelihood”

$$Z = p(d|\mathcal{M}) = \sum_H p(d|H, I) p(H|I)$$

sum over different (discrete) hypothesis  
(given that  $I$  is true)

## Bayesian evidence

“Bayesian evidence” or “Marginal likelihood”

$$Z = p(d|\mathcal{M}) = \int_{\Omega_{\mathcal{M}}} p(d|\theta, \mathcal{M}) p(\theta|\mathcal{M}) d\theta$$

integrate over all possible (continuous) parameters of model  $\mathcal{M}$   
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Model posterior:

$$p(\mathcal{M}_i|d) \propto p(\mathcal{M}_i) Z_i$$

given model prior  $p(\mathcal{M}_i)$

proportional to  
constant that  
depends only on data



Posterior odds of  $\mathcal{M}_1$  versus  $\mathcal{M}_2$ :

$$\frac{p(\mathcal{M}_1|d)}{p(\mathcal{M}_2|d)} = B_{1,2} \frac{p(\mathcal{M}_1)}{p(\mathcal{M}_2)}$$

Bayes factor:

$$B_{1,2} = \frac{Z_1}{Z_2} \Rightarrow \ln B_{1,2} = \ln Z_1 - \ln Z_2$$

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if priors are the same [ $p(\mathcal{M}_1) = p(\mathcal{M}_2)$ ],  
 $B_{1,2}$  tells which one is preferred:

$B_{1,2} > 1$  ( $\ln B_{1,2} > 0$ )

$\mathcal{M}_1$  preferred

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$\mathcal{M}_1$  preferred

$$B_{1,2} < 1 \quad (\ln B_{1,2} < 0)$$

$\mathcal{M}_2$  preferred

$|B_{1,2}|$  tells the odds in favor of preferred model

# Jeffreys' scale

odds in favor of the preferred model:

$$|B_{1,2}| : 1$$

strength of preference according to Jeffreys' scale:

$ \ln B_{1,2} $	Odds	probability	strength of evidence
$< 1.0$	$\lesssim 3 : 1$	$< 0.750$	inconclusive
$\in [1.0, 2.5]$	$(3 - 12) : 1$	$< 0.923$	weak
$\in [2.5, 5.0]$	$(12 - 150) : 1$	$< 0.993$	moderate
$> 5.0$	$> 150 : 1$	$> 0.993$	strong

odds & strength always valid

probability correct given equal priors and that only two models are possible (see e.g. neutrino mass ordering: normal OR inverted)

## Occam's razor

what the Bayesian model comparison tells us?

Best model strikes optimum balance between

Quality of fit

Predictivity

Occam's razor

the simplest theory that fits data is preferred

model with more parameters  $\longrightarrow$  better fit (usually)

$\longmapsto$  are all the parameters needed?

Bayes factor penalizes unnecessarily complex models!

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Bayesian evidence depends on priors!

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Bayes factor DOES NOT penalize models with parameters that are unconstrained by the data

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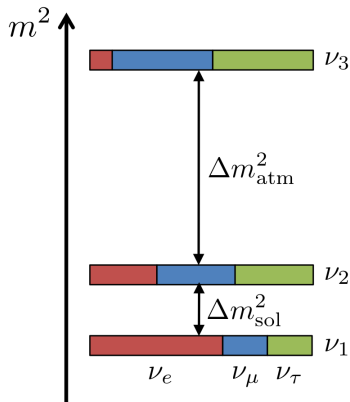


# Neutrino masses

## Normal ordering (NO)

$$m_1 < m_2 < m_3$$

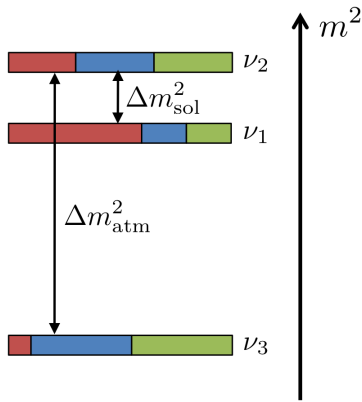
$$\sum m_k \gtrsim 0.06 \text{ eV}$$



## Inverted ordering (IO)

$$m_3 < m_1 < m_2$$

$$\sum m_k \gtrsim 0.1 \text{ eV}$$

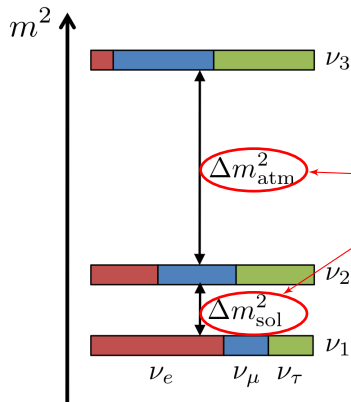


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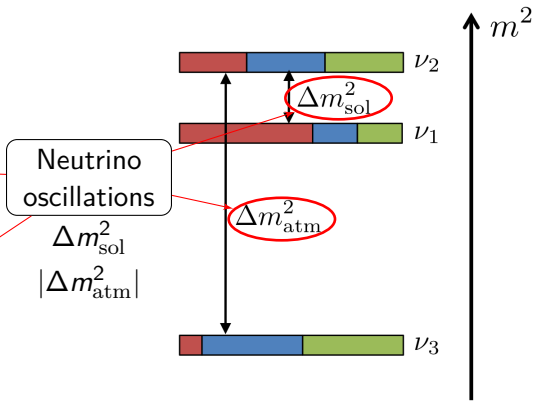
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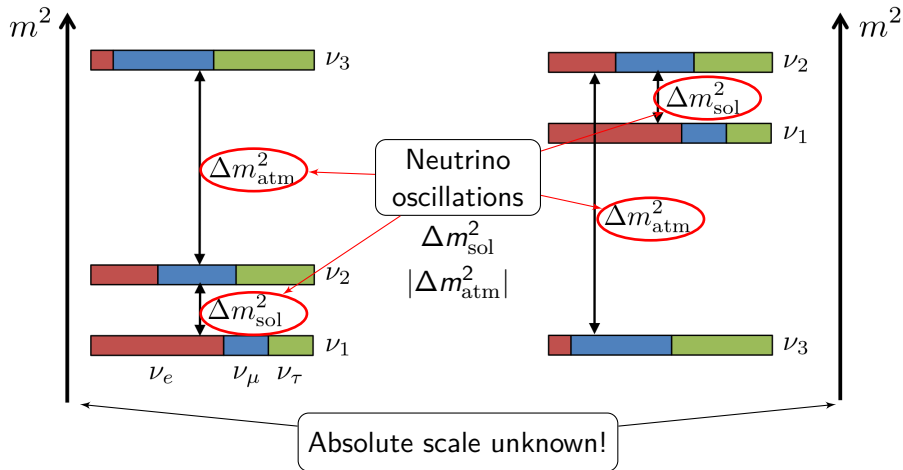
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Constrain mass ordering by constraining  $\sum m_k$

## Constraining the absolute scale of neutrino masses

Neutrino effects on  $\beta$  decay endpoint

Mainz/Troitsk limits,  $m_{\nu_e} \lesssim 2$  eV

Katrin, (expected)  $m_{\nu_e} \lesssim 0.2$  eV

$$m_{\nu_e}^2 = \sum_k |U_{ek}|^2 m_k^2$$

$U_{ek}$  mixing matrix

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(if neutrino is Majorana)

Constraints from neutrinoless double beta decay

Measure  $T_{1/2}^{0\nu}$ , convert into  $m_{\beta\beta}$  using  $m_e/m_{\beta\beta} = \mathcal{M}'^{\nu} \sqrt{G_{0\nu} T_{1/2}^{0\nu}}$

and then use  $m_{\beta\beta} = \left| \sum_k e^{i\alpha_k} U_{ek}^2 m_k \right|$

$\alpha_k$  Majorana phases

$m_e$  electron mass,  
 $G_{0\nu}$  phase space,  
 $\mathcal{M}'^{\nu}$  matrix element

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## Cosmological constraints from neutrino mass effects

Mainly through CMB effects and free streaming

Currently can constrain  $\sum_k m_k \lesssim 0.1X$  eV, single  $m_k$  in future?

## Can current data tell us the neutrino mass ordering?

- [Hannestad, Schwetz, 2016]: extremely weak (2:1, 3:2) preference for NO (cosmology + [Bergstrom et al., 2015] neutrino oscillation fit) Bayesian approach;
- [Gerbino et al, 2016]: extremely weak (up to 3:2) preference for NO (cosmology only), Bayesian approach;
- [Simpson et al., 2017]: strong preference for NO (cosmological limits on  $\sum m_\nu$  + constraints on  $\Delta m_{21}^2$  and  $|\Delta m_{31}^2|$ ) Bayesian approach;
- [Capozzi et al., 2017]:  $2\sigma$  preference for NO (cosmology + [Capozzi et al., 2016, updated 2017] neutrino oscillation fit) frequentist approach;
- [Caldwell et al., 2017] very mild indication for NO (cosmology + neutrinoless double-beta decay + [Esteban et al., 2016] readapted oscillation results) Bayesian approach;
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## Neutrino oscillations

full  $\chi^2 = -2 \log \mathcal{L}_{\text{osc}}$   
from global fit

[de Salas et al, 2017]

### Neutrino mixing

Parameter	Prior
$\sin^2 \theta_{12}$	0.1 – 0.6
$\sin^2 \theta_{13}$	0.00 – 0.06
$\sin^2 \theta_{23}$	0.25 – 0.75

Masses: see later!

$\beta\beta 0\nu$  data

Likelihood approximations as in [Caldwell et al, 2017], from [Gerda, 2017] (Ge), [KamLAND-Zen, 2016], [EXO-200, 2014] (Xe)

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$\beta\beta 0\nu$ parameters		Neutrino mixing	
Parameter	Prior	Parameter	Prior
$\alpha_2$	$0 - 2\pi$	$\sin^2 \theta_{12}$	$0.1 - 0.6$
$\alpha_3$	$0 - 2\pi$	$\sin^2 \theta_{13}$	$0.00 - 0.06$
$\mathcal{M}_{76\text{Ge}}^{0\nu}$	$4.07 - 4.87$	$\sin^2 \theta_{23}$	$0.25 - 0.75$
$\mathcal{M}_{136\text{Xe}}^{0\nu}$	$2.74 - 3.45$		

Masses: see later!

# Parameterizations, priors and data

[Gariazzo et al., in preparation]

Cosmological data

Full CMB temperature and polarization spectra from [Planck, 2015], working with  $\Lambda$ CDM model as basis

$\beta\beta 0\nu$  data

Likelihood approximations as in [Caldwell et al, 2017], from [Gerda, 2017] (Ge), [KamLAND-Zen, 2016], [EXO-200, 2014] (Xe)

Neutrino oscillations

full  $\chi^2 = -2 \log \mathcal{L}_{\text{osc}}$  from global fit [de Salas et al, 2017]

Cosmological		$\beta\beta 0\nu$ parameters		Neutrino mixing	
Parameter	Prior	Parameter	Prior	Parameter	Prior
$\omega_b$	0.019 – 0.025	$\alpha_2$	0 – $2\pi$	$\sin^2 \theta_{12}$	0.1 – 0.6
$\omega_c$	0.095 – 0.145	$\alpha_3$	0 – $2\pi$	$\sin^2 \theta_{13}$	0.00 – 0.06
$\Theta_s$	1.03 – 1.05	$\mathcal{M}_{76\text{Ge}}^{0\nu}$	4.07 – 4.87	$\sin^2 \theta_{23}$	0.25 – 0.75
$\tau$	0.01 – 0.4	$\mathcal{M}_{136\text{Xe}}^{0\nu}$	2.74 – 3.45		
$n_s$	0.885 – 1.04				
$\log(10^{10} A_s)$	2.5 – 3.7				

Masses: see later!

[Simpson et al, 2017]

using  $m_1, m_2, m_3$  (A)

[Caldwell et al, 2017]

using  $m_{\text{lightest}}, \Delta m_{21}^2, |\Delta m_{31}^2|$  (B)

intuition says: (B) is closer to observable quantities! Better than (A)?

Should we use linear or logarithmic priors on  $m_k$  ( $m_{\text{lightest}}$ )?

Can data help to select (A) or (B), linear or log?

# Modeling neutrino masses

[Simpson et al, 2017]

using  $m_1, m_2, m_3$  (A)

[Caldwell et al, 2017]

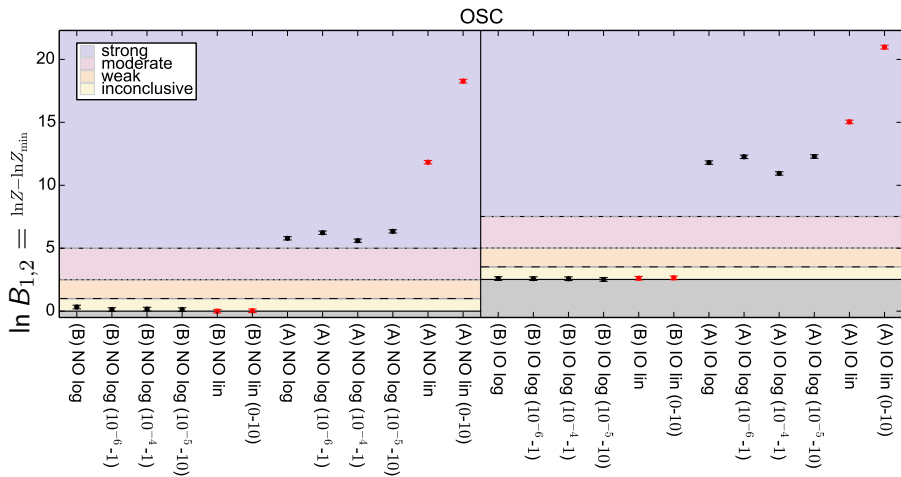
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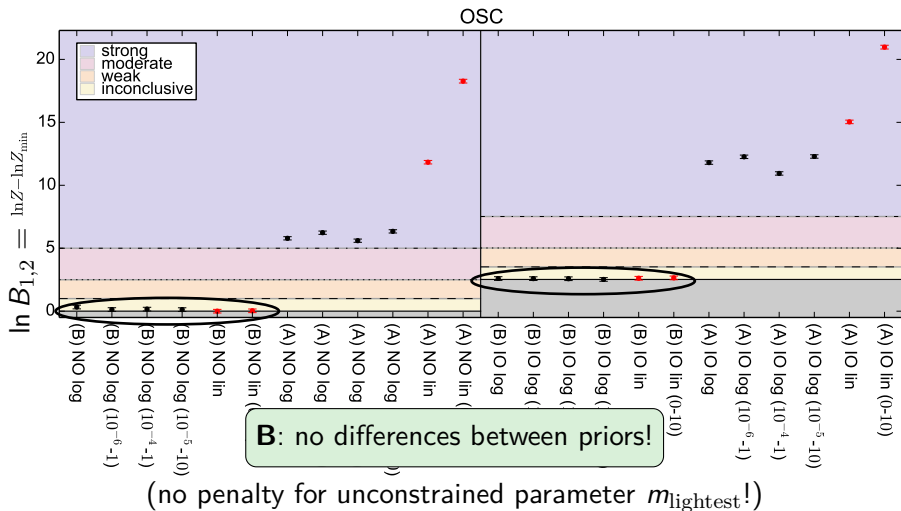
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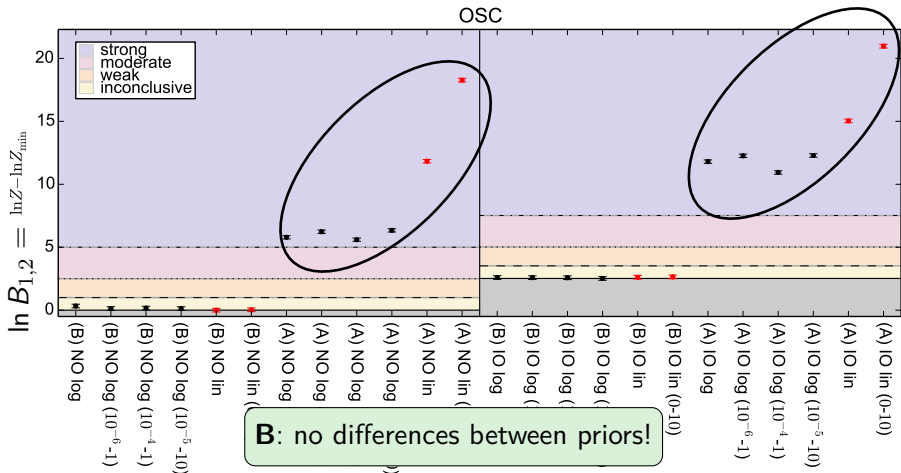
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Model A			Model B		
Parameter	Prior	Range	Parameter	Prior	Range
$m_1/\text{eV}$	linear log	0 – 1 $10^{-5} - 1$	$m_{\text{lightest}}/\text{eV}$	linear log	0 – 1 $10^{-5} - 1$
$m_2/\text{eV}$	linear log	0 – 1 $10^{-5} - 1$	$\Delta m_{21}^2/\text{eV}^2$	linear	$5 \times 10^{-5} - 10^{-4}$
$m_3/\text{eV}$	linear log	0 – 1 $10^{-5} - 1$	$ \Delta m_{31}^2 /\text{eV}^2$	linear	$1.5 \times 10^{-3} - 3.5 \times 10^{-3}$





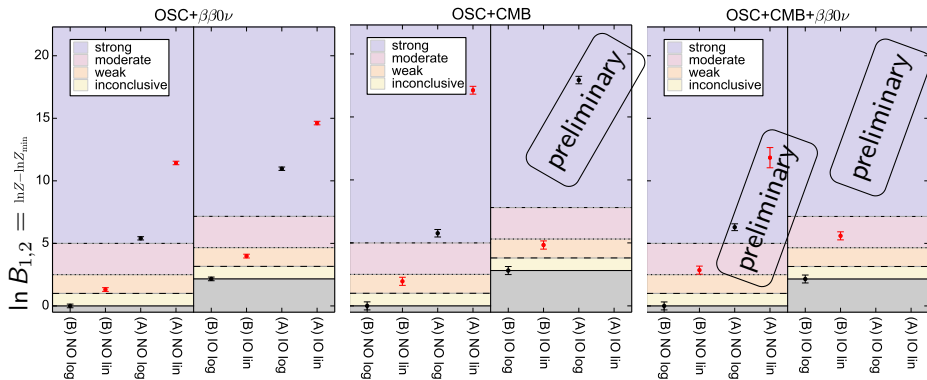


(no penalty for unconstrained parameter  $m_{\text{lightest}}$ !)

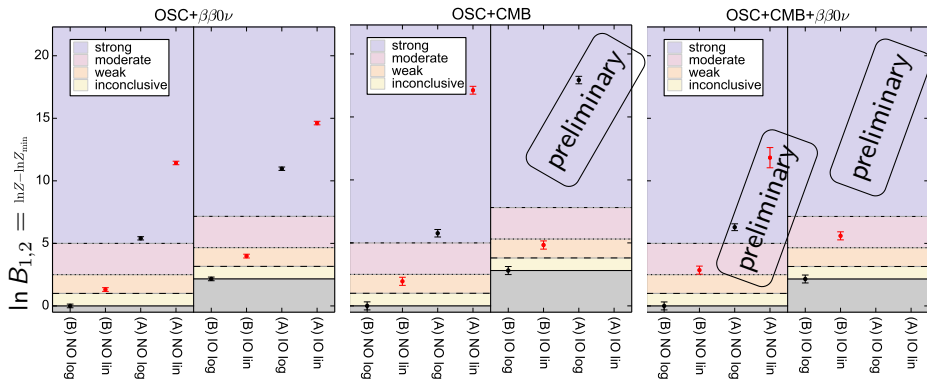
**A: always strongly disfavored!**

(waste of parameter space, no unconstrained parameter due to  $\Delta m_{11}^2$ !)



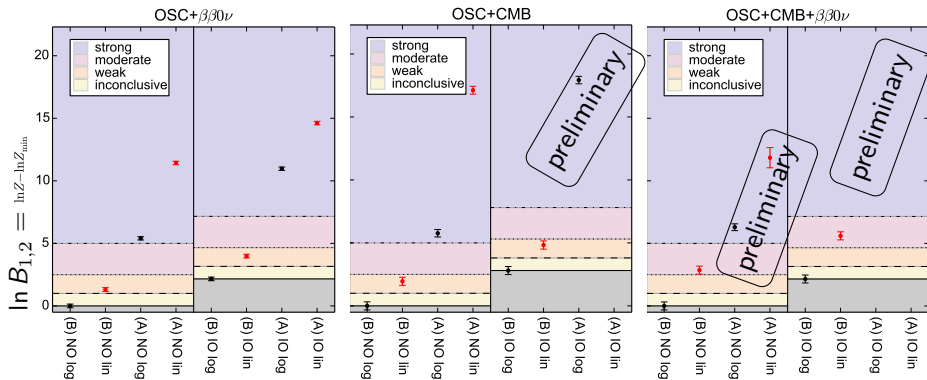


compare **linear** versus **logarithmic**



compare **linear** versus **logarithmic**

$\ln B_{\log, \text{lin}}$  prefers **log** priors:  
weakly-to-moderately for model B, moderately-to-strongly for model A

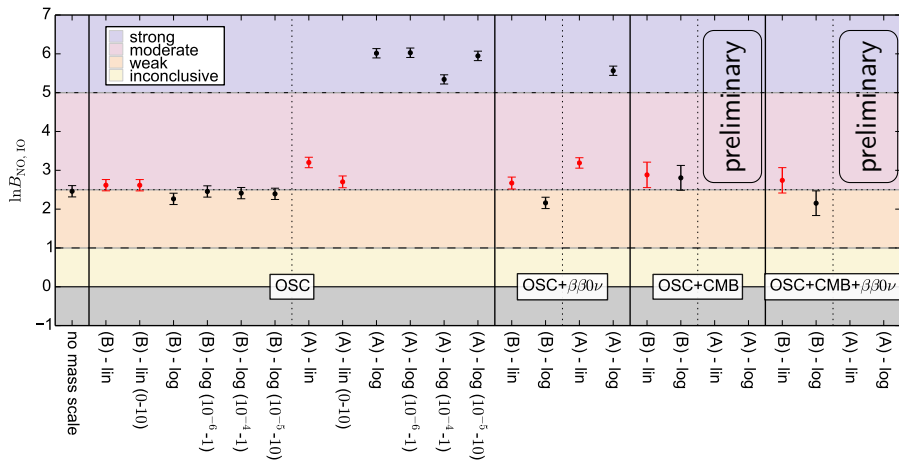


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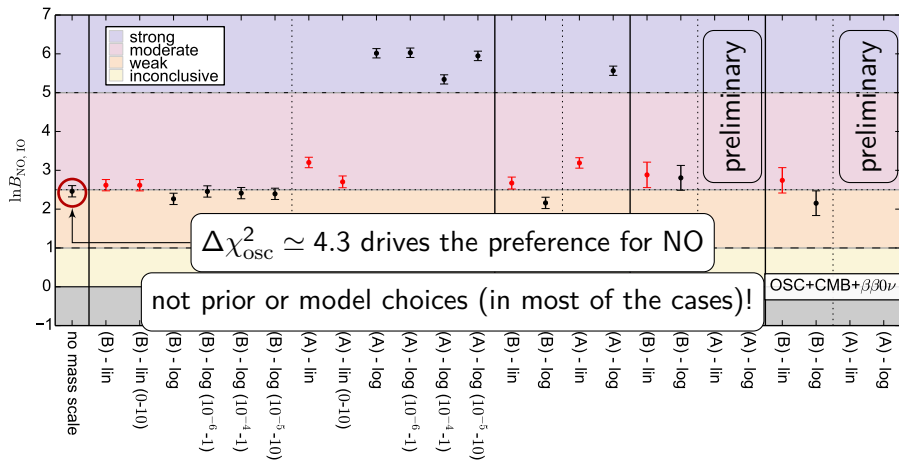
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summary: model B, log priors is better!

# Comparing the mass orderings

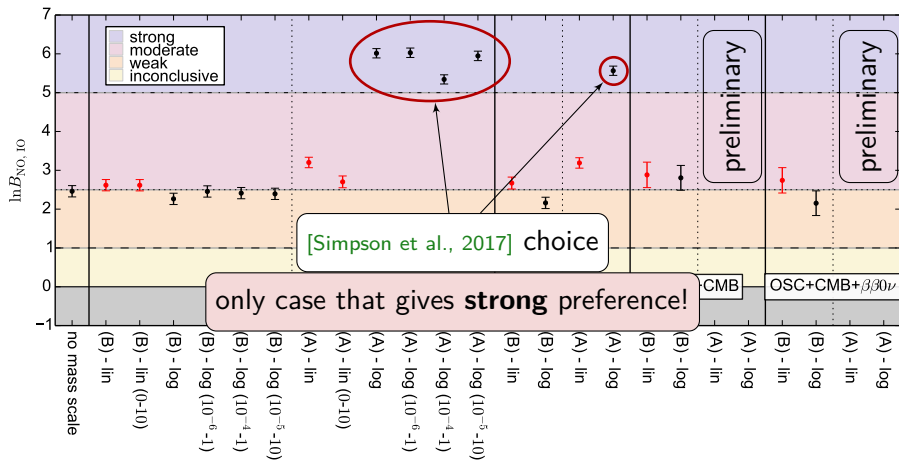


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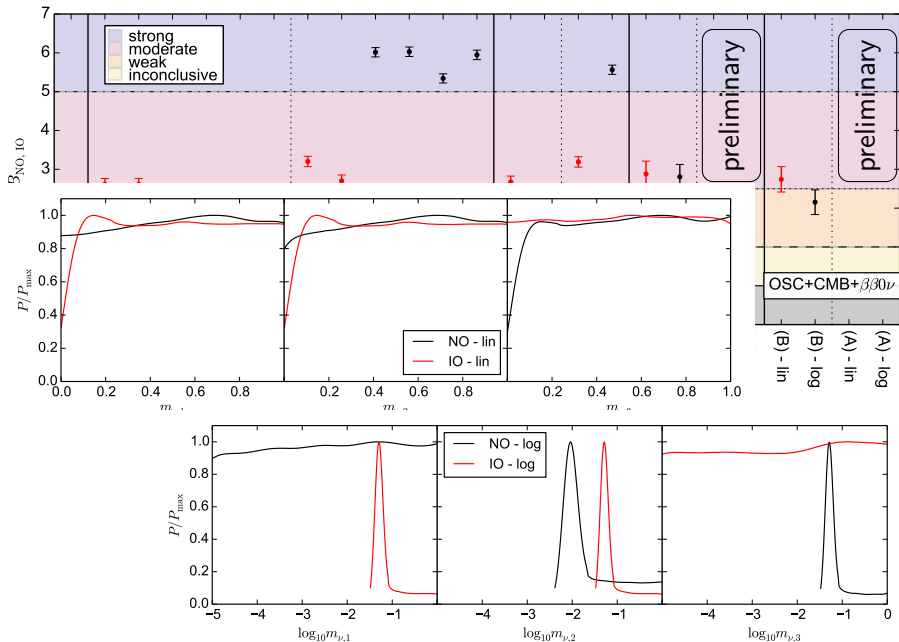
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[Gariazzo et al., in preparation]



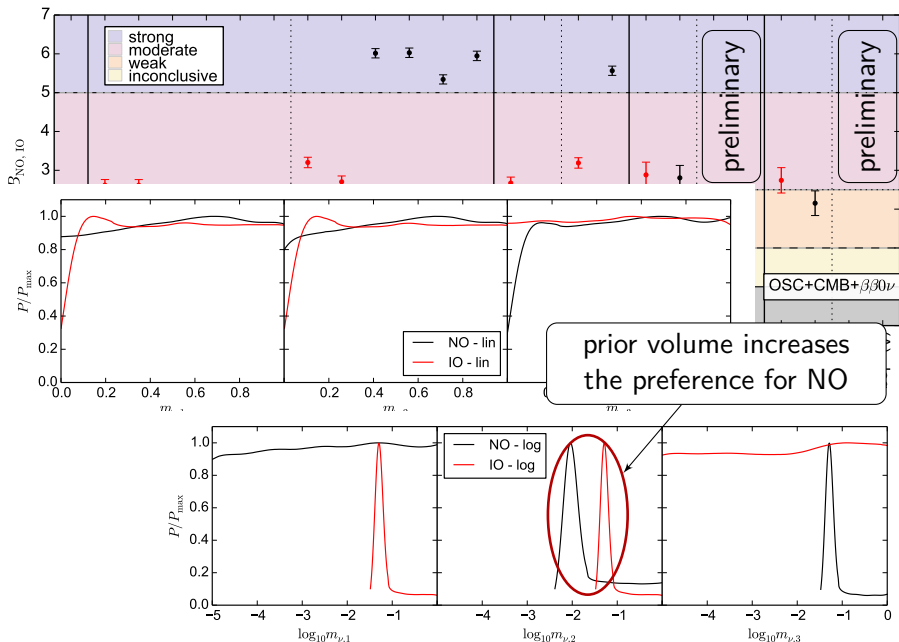
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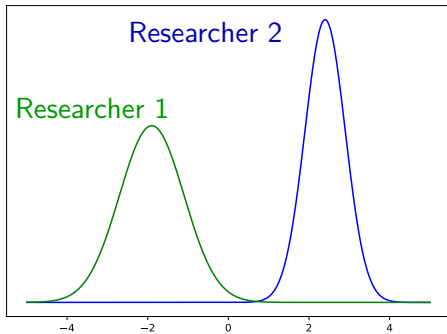


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# Bayes theorem in action

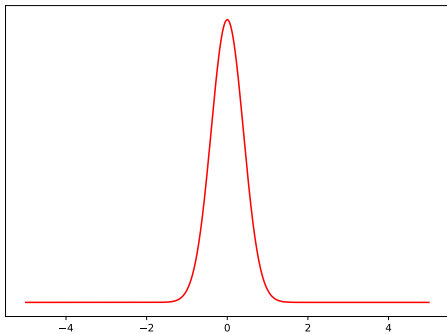
$$\text{posterior} = \frac{\text{likelihood} \times \text{prior}}{\text{evidence}}$$

Prior



What each researcher knew  
before the experiment

Likelihood

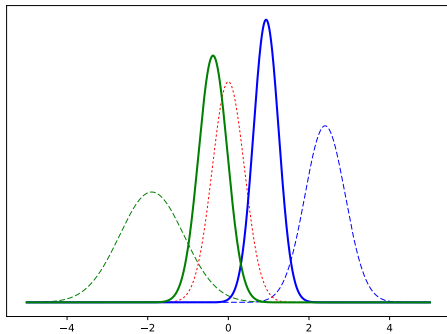


The result of the experiment

# Bayes theorem in action

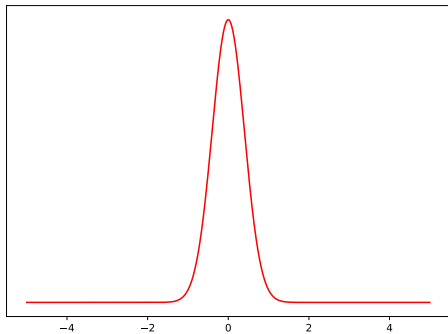
$$\text{posterior} = \frac{\text{likelihood} \times \text{prior}}{\text{evidence}}$$

Posterior



What each researcher knows after the experiment

Likelihood



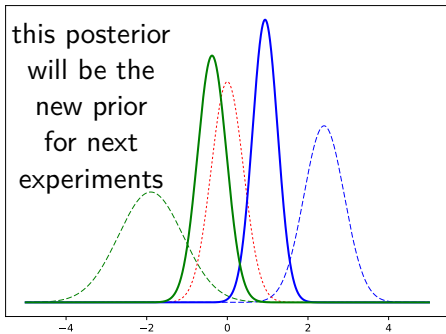
The result of the experiment

Posterior depends on prior!

# Bayes theorem in action

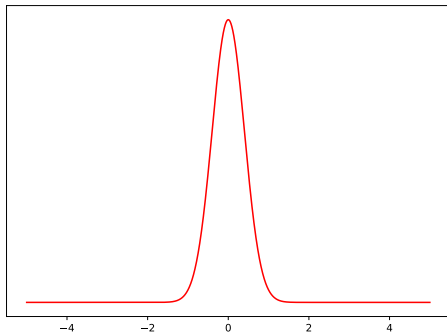
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Posterior



What each researcher knows after the experiment

Likelihood



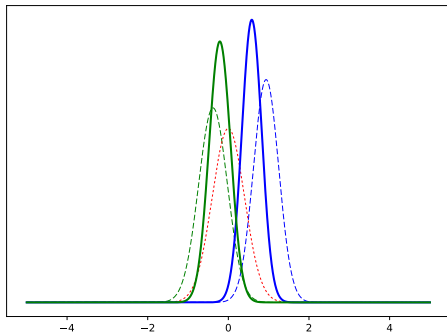
The result of the experiment

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# Bayes theorem in action

$$\text{posterior} = \frac{\text{likelihood} \times \text{prior}}{\text{evidence}}$$

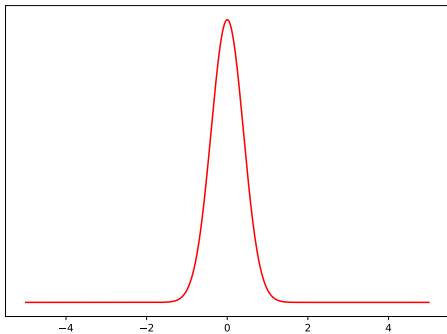
Posterior



What each researcher knows  
after the second experiment

Remember:  
 $\sigma_N^2 = \sigma^2/N$

Likelihood



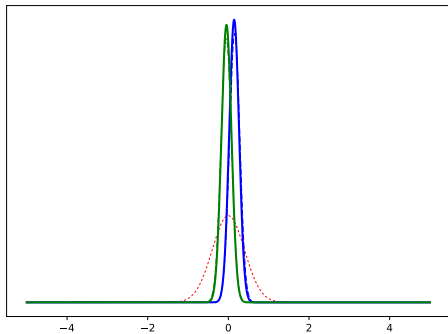
The result of the experiment

Posterior depends on prior!

# Bayes theorem in action

$$\text{posterior} = \frac{\text{likelihood} \times \text{prior}}{\text{evidence}}$$

Posterior

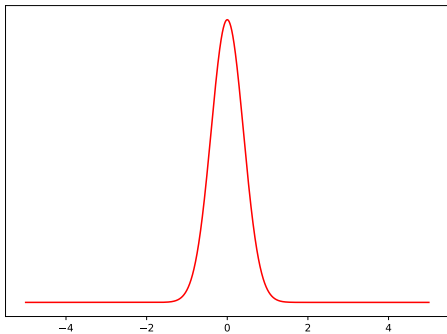


What each researcher  
knows after 10 experiments

Remember:

$$\sigma_N^2 = \sigma^2/N$$

Likelihood

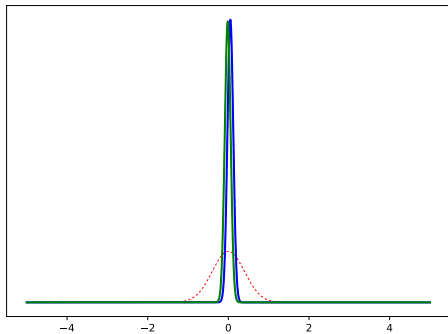


The result of the experiment

## Bayes theorem in action

$$\text{posterior} = \frac{\text{likelihood} \times \text{prior}}{\text{evidence}}$$

Posterior

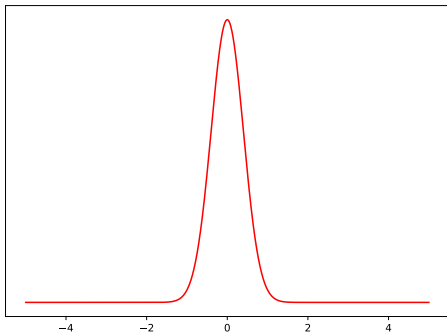


What each researcher  
knows after 30 experiments

Remember:

$$\sigma_N^2 = \sigma^2/N$$

Likelihood

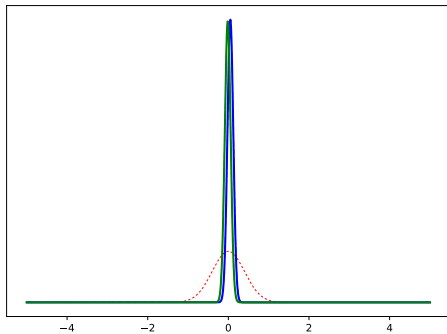


The result of the experiment

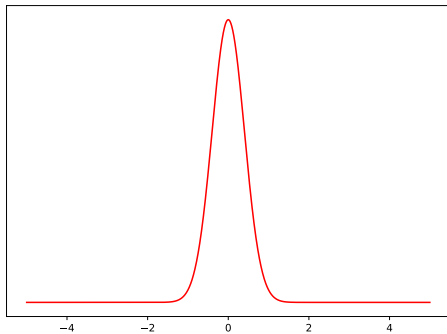
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$$\text{posterior} = \frac{\text{likelihood} \times \text{prior}}{\text{evidence}}$$

Posterior



Likelihood



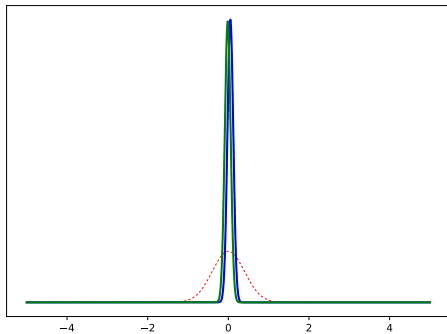
Knowledge converges using information from experiments



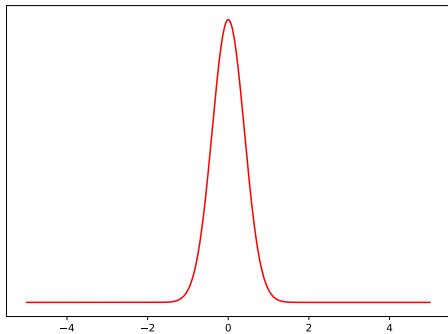
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Posterior



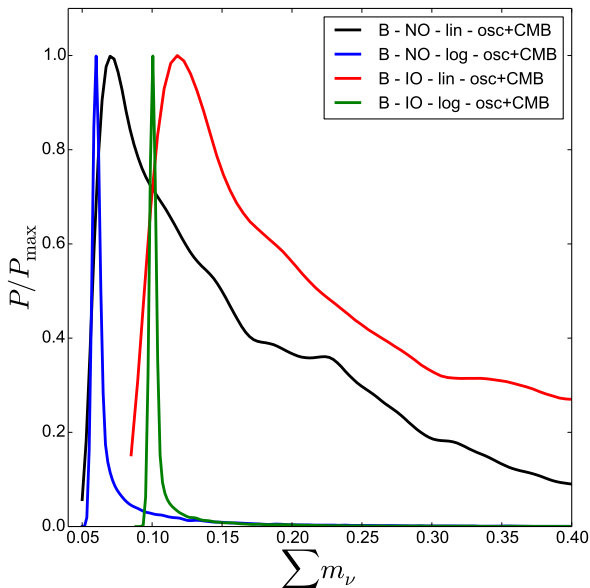
Likelihood



Knowledge converges using information from experiments

Prior dependence (subjectivity) only if not enough information in data!

# The role of priors: $\sum m_\nu$



showing **model B only**  
(only 1 parameter changes)

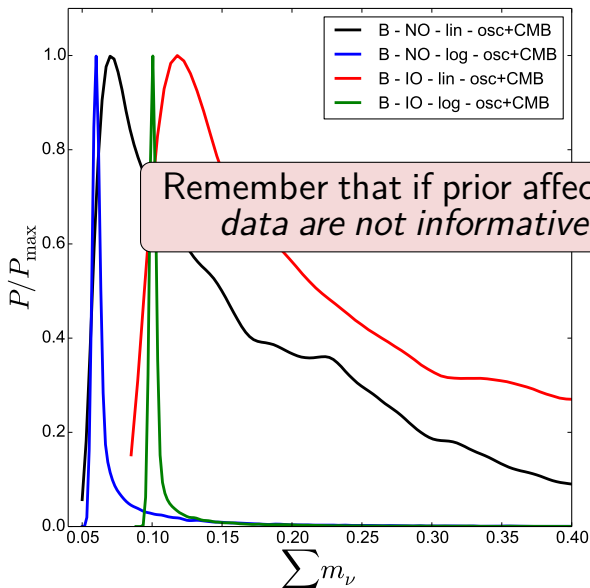
same for model A,  
but **amplified** (3  
parameters change!)

**logarithmic** prior  
corresponds to  
 $1/m_k$  probability!

↓  
more importance  
to smaller masses

↓  
limits closer to  
minimum allowed  
value of  $\sum m_\nu$

# The role of priors: $\sum m_\nu$



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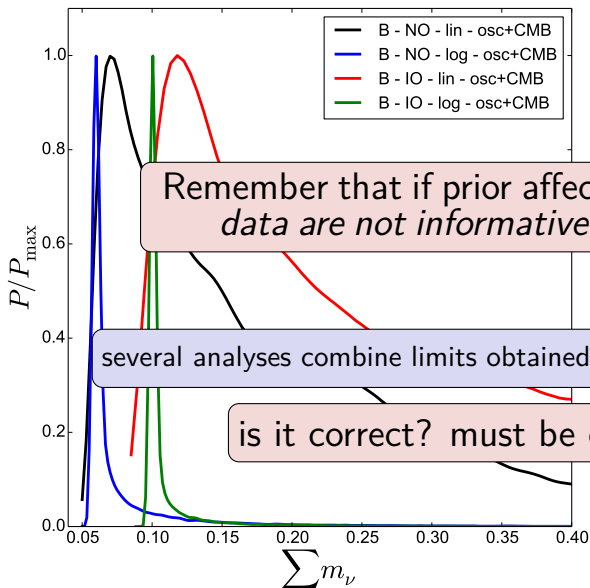
same for model A,  
changed (3 change!)

logarithmic prior  
corresponds to  
 $1/m_k$  probability!

↓  
more importance  
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↓  
limits closer to  
minimum allowed  
value of  $\sum m_\nu$

# The role of priors: $\sum m_\nu$



showing **model B only**  
(only 1 parameter changes)

same for model A,  
changed (3 change!)

Remember that if prior affects posterior,  
*data are not informative enough!*

several analyses combine limits obtained with different priors

is it correct? must be careful!

logarithmic prior  
corresponds to

more importance  
smaller masses

limits closer to  
minimum allowed  
value of  $\sum m_\nu$

- 1 Basics of Bayesian statistics
  - Bayes' theorem
  - Bayesian model comparison
- 2 Constraining the neutrino mass ordering
  - Introducing the problem
  - Comparing models and mass orderings
- 3 Constraining the neutrino masses
- 4 Conclusions

# Conclusions

## Bayesian model comparison

1

through Bayesian evidence/Bayes factor  
to **robustly test models**/priors against data

2

Be **careful with** the effects of prior  
(or of **other subjective choices**)  
on the results of your calculations

3

**data** only **weakly/moderately** prefer normal  
versus inverted neutrino mass ordering

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Thank you for the attention!