Stefano Gariazzo

IFIC, Valencia (ES)
CSIC – Universitat de Valencia

gariazzo@ific.uv.es
http://ific.uv.es/~gariazzo/

Horizon 2020
European Union funding
for Research & Innovation

Light sterile neutrinos: oscillations and cosmology

Matter to the Deepest, Chorzów (PL), 1–6/09/2019

Three Neutrino Oscillations

$$\nu_\alpha = \sum_{k=1}^{3} U_{\alpha k} \nu_k \quad (\alpha = e, \mu, \tau)$$

$U_{\alpha k}$ described by 3 mixing angles θ_{12}, θ_{13}, θ_{23} and one CP phase $\delta_{\rm CP}$

Current knowledge of the 3 active ν mixing: [de Salas et al. (2018)]

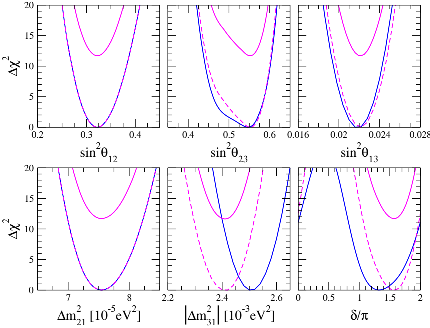
NO: Normal Ordering, $m_1 < m_2 < m_3$

IO: Inverted Ordering, $m_3 < m_1 < m_2$

$$\begin{aligned} \Delta m^2_{21} &= (7.55^{+0.20}_{-0.16}) \cdot 10^{-5}\ \mathrm{eV}^2 \\ |\Delta m^2_{31}| &= (2.50 \pm 0.03) \cdot 10^{-3}\ \mathrm{eV}^2\ (\mathrm{NO}) \\ &= (2.42^{+0.03}_{-0.04}) \cdot 10^{-3}\ \mathrm{eV}^2\ (\mathrm{IO}) \end{aligned}$$

$$\begin{aligned} \sin^2(\theta_{12}) &= 0.320^{+0.020}_{-0.016} \\ \sin^2(\theta_{13}) &= 0.0216^{+0.008}_{-0.007}\ (\mathrm{NO}) \\ &= 0.0222^{+0.007}_{-0.008}\ (\mathrm{IO}) \\ \sin^2(\theta_{23}) &= 0.547^{+0.020}_{-0.030}\ (\mathrm{NO}) \\ &= 0.551^{+0.018}_{-0.030}\ (\mathrm{IO}) \end{aligned}$$

First hints for $\delta_{\rm CP} \simeq 3/2\pi$



see also: `http://globalfit.astroparticles.es`

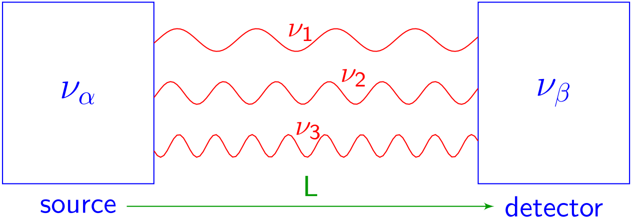
Two types of neutrinos

flavor neutrinos ν_α

$|\nu_\alpha\rangle = U_{\alpha k}|\nu_k\rangle$

massive neutrinos ν_k

$$|\nu(t=0)\rangle = |\nu_\alpha\rangle = U_{\alpha 1}|\nu_1\rangle + U_{\alpha 2}|\nu_2\rangle + U_{\alpha 3}|\nu_3\rangle$$



$$|\nu(t>0)\rangle = |\nu_\beta\rangle = U_{\alpha 1}\, e^{-iE_1 t}|\nu_1\rangle + U_{\alpha 2}\, e^{-iE_2 t}|\nu_2\rangle + U_{\alpha 3}\, e^{-iE_3 t}|\nu_3\rangle \neq |\nu_\alpha\rangle$$

$E_k^2 = p^2 + m_k^2 \longleftarrow$ define $\longrightarrow t = L$

$$P_{\nu_\alpha\to\nu_\beta}(L) = |\langle\nu_\alpha|\nu(L)\rangle|^2 = \sum_{k,j} U_{\beta k} U^*_{\alpha k} U^*_{\beta j} U_{\alpha j} \exp\left(-i\frac{\Delta m^2_{kj} L}{2E}\right)$$

$$\Delta m^2_{ij} = m_i^2 - m_j^2$$

A large family

In principle, previous discussion is valid for N neutrinos

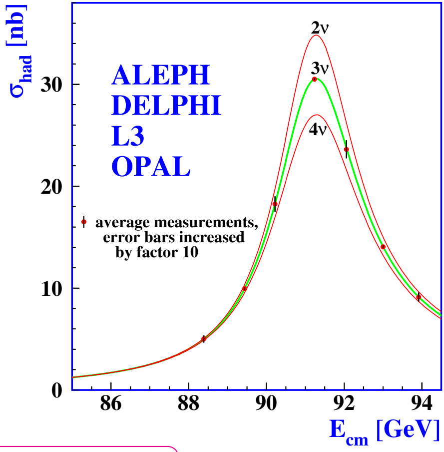
only constraint: there are exactly three flavor neutrinos in the SM

[LEP, Phys. Rept. 427 (2006) 257, arXiv:hep-ex/0509008]

$$N_\nu^{(Z)} = 2.9840 \pm 0.0082$$

through the measurement of the Z resonance

$$e^+e^- \longrightarrow Z \longrightarrow \sum_{a=e,\mu,\tau} \nu_a \bar{\nu}_a$$



neutrinos $\alpha > 3$ must be sterile

sterile neutrino = SM singlet: no couplings with other SM particles

A large family

In principle, previous discussion is valid for N neutrinos

$N \times N$ mixing matrix, N flavor neutrinos, N massive neutrinos

$$
\begin{pmatrix} |\nu_e\rangle \\ |\nu_\mu\rangle \\ |\nu_\tau\rangle \\ |\nu_{s_1}\rangle \\ \ldots \end{pmatrix}
=
\begin{pmatrix}
U_{e1} & U_{e2} & U_{e3} & U_{e4} & \vdots \\
U_{\mu1} & U_{\mu2} & U_{\mu3} & U_{\mu4} & \\
U_{\tau1} & U_{\tau2} & U_{\tau3} & U_{\tau4} & \\
U_{s_1 1} & U_{s_1 2} & U_{s_1 3} & U_{s_1 4} & \\
\ldots & & & & \ddots
\end{pmatrix}
\begin{pmatrix} |\nu_1\rangle \\ |\nu_2\rangle \\ |\nu_3\rangle \\ |\nu_4\rangle \\ \ldots \end{pmatrix}
$$

A large family

In principle, previous discussion is valid for N neutrinos

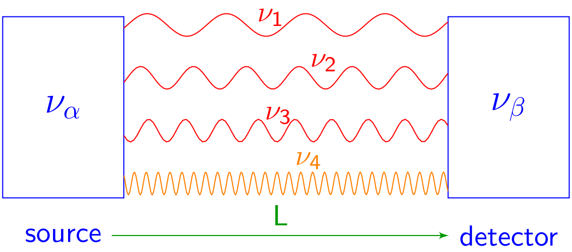
$N \times N$ mixing matrix, N flavor neutrinos, N massive neutrinos

$$\begin{pmatrix} |\nu_e\rangle \\ |\nu_\mu\rangle \\ |\nu_\tau\rangle \\ |\nu_{s_1}\rangle \\ \dots \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & U_{e4} & \vdots \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} & U_{\mu 4} & \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} & U_{\tau 4} & \\ U_{s_1 1} & U_{s_1 2} & U_{s_1 3} & U_{s_1 4} & \\ \dots & & & & \ddots \end{pmatrix} \begin{pmatrix} |\nu_1\rangle \\ |\nu_2\rangle \\ |\nu_3\rangle \\ |\nu_4\rangle \\ \dots \end{pmatrix}$$

Our case will be 3 (active)+1 (sterile), a perturbation of 3 neutrinos case

$m_4 \gg m_\ell$
$\ell = 1, 2, 3$

$\Delta m^2_{41} \simeq 1\ \text{eV}^2 \simeq \Delta m^2_{42} \simeq \Delta m^2_{43}$



6 angles

3 Dirac CP phases

3 Majorana CP phases

Short BaseLine (SBL)

$$P_{\nu_\alpha\to\nu_\beta}(L) = |\langle\nu_\alpha|\nu(L)\rangle|^2 = \sum_{k,j} U_{\beta k} U^*_{\alpha k} U^*_{\beta j} U_{\alpha j} \exp\left(-i\frac{\Delta m^2_{kj} L}{2E}\right)$$

If $m_4 \gg m_\ell$, faster oscillations

ν_4 oscillations are averaged in most neutrino oscillation experiments

Effect of 4th neutrino only visible as global normalization

Short BaseLine (SBL) oscillations: $\frac{\Delta m^2_{41} L}{E} \simeq 1$

At SBL, oscillations due to Δm^2_{21} and $|\Delta m^2_{31}|$ do not develop

Short BaseLine (SBL)

$$P_{\nu_\alpha\to\nu_\beta}(L) = |\langle\nu_\alpha|\nu(L)\rangle|^2 = \sum_{k,j} U_{\beta k}U^*_{\alpha k}U^*_{\beta j}U_{\alpha j}\exp\left(-i\frac{\Delta m^2_{kj}L}{2E}\right)$$

If $m_4 \gg m_\ell$, faster oscillations

ν_4 oscillations are averaged in most neutrino oscillation experiments

Effect of 4th neutrino only visible as global normalization

Short BaseLine (SBL) oscillations: $\frac{\Delta m^2_{41}L}{E} \simeq 1$

At SBL, oscillations due to Δm^2_{21} and $|\Delta m^2_{31}|$ do not develop

APPearance ($\alpha \neq \beta$)

source of $\overset{(-)}{\nu}_\alpha$

detect different flavor $\overset{(-)}{\nu}_\beta$

Short BaseLine (SBL)

$$P_{\nu_\alpha \to \nu_\beta}(L) = |\langle \nu_\alpha | \nu(L) \rangle|^2 = \sum_{k,j} U_{\beta k} U^*_{\alpha k} U^*_{\beta j} U_{\alpha j} \exp\left(-i \frac{\Delta m^2_{kj} L}{2E}\right)$$

If $m_4 \gg m_\ell$, faster oscillations

ν_4 oscillations are averaged in most neutrino oscillation experiments

Effect of 4th neutrino only visible as global normalization

Short BaseLine (SBL) oscillations: $\frac{\Delta m^2_{41} L}{E} \simeq 1$

At SBL, oscillations due to Δm^2_{21} and $|\Delta m^2_{31}|$ do not develop

APPearance ($\alpha \neq \beta$)

source of $\overset{(-)}{\nu}_\alpha$

detect different flavor $\overset{(-)}{\nu}_\beta$

DISappearance

source of $\overset{(-)}{\nu}_\alpha$

$\overset{(-)}{\nu}_s$ goes undetected

Short BaseLine (SBL)

$$P_{\nu_\alpha \to \nu_\beta}(L) = |\langle \nu_\alpha | \nu(L) \rangle|^2 = \sum_{k,j} U_{\beta k} U^*_{\alpha k} U^*_{\beta j} U_{\alpha j} \exp\left(-i \frac{\Delta m^2_{kj} L}{2E} \right)$$

If $m_4 \gg m_\ell$, faster oscillations

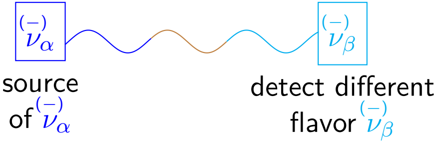
ν_4 oscillations are averaged in most neutrino oscillation experiments

Effect of 4th neutrino only visible as global normalization

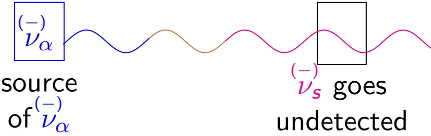
Short BaseLine (SBL) oscillations: $\frac{\Delta m^2_{41} L}{E} \simeq 1$

At SBL, oscillations due to Δm^2_{21} and $|\Delta m^2_{31}|$ do not develop

APPearance ($\alpha \neq \beta$)



DISappearance



CP violation cannot be observed in SBL experiments!

New mixings in the 3+1 scenario

4 × 4 mixing matrix: $\begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & U_{e4} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} & U_{\mu 4} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} & U_{\tau 4} \\ U_{s_1 1} & U_{s_1 2} & U_{s_1 3} & U_{s_1 4} \end{pmatrix}$

New mixings in the 3+1 scenario

4×4 mixing matrix: $\begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & U_{e4} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} & U_{\mu 4} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} & U_{\tau 4} \\ U_{s_1 1} & U_{s_1 2} & U_{s_1 3} & U_{s_1 4} \end{pmatrix}$ ϑ_{14} ϑ_{24} ϑ_{34}

New mixings in the 3+1 scenario

4 × 4 mixing matrix: $\begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & U_{e4} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} & U_{\mu4} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} & U_{\tau4} \\ U_{s_1 1} & U_{s_1 2} & U_{s_1 3} & U_{s_1 4} \end{pmatrix}$ ϑ_{14} ϑ_{24} ϑ_{34}

DISappearance

$$P^{\rm SBL}_{\overset{(-)}{\nu}_\alpha \to \overset{(-)}{\nu}_\alpha} \simeq 1 - \sin^2 2\vartheta_{\alpha\alpha} \sin^2\left(\frac{\Delta m^2_{41} L}{4E}\right)$$

$$\sin^2 2\vartheta_{\alpha\alpha} = 4|U_{\alpha4}|^2(1-|U_{\alpha4}|^2)$$

$\overset{(-)}{\nu}_e \to \overset{(-)}{\nu}_e$

reactor

gallium

$$|U_{e4}|^2 = \sin^2\vartheta_{14}$$

$\overset{(-)}{\nu}_\mu \to \overset{(-)}{\nu}_\mu$

accelerator

atmospheric

$$|U_{\mu4}|^2 = \cos^2\vartheta_{14}\sin^2\vartheta_{24}$$

New mixings in the 3+1 scenario

4 × 4 mixing matrix: $\begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & U_{e4} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} & U_{\mu4} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} & U_{\tau4} \\ U_{s_1 1} & U_{s_1 2} & U_{s_1 3} & U_{s_1 4} \end{pmatrix}$ ϑ_{14}, ϑ_{24}, ϑ_{34}

DISappearance

$$P^{\mathrm{SBL}}_{\overset{(-)}{\nu}_\alpha \to \overset{(-)}{\nu}_\alpha} \simeq 1 - \sin^2 2\vartheta_{\alpha\alpha} \sin^2\left(\frac{\Delta m^2_{41} L}{4E}\right)$$

$$\sin^2 2\vartheta_{\alpha\alpha} = 4|U_{\alpha4}|^2(1-|U_{\alpha4}|^2)$$

$\overset{(-)}{\nu}_e \to \overset{(-)}{\nu}_e$ — reactor, gallium

$$|U_{e4}|^2 = \sin^2\vartheta_{14}$$

$\overset{(-)}{\nu}_\mu \to \overset{(-)}{\nu}_\mu$ — accelerator, atmospheric

$$|U_{\mu4}|^2 = \cos^2\vartheta_{14}\sin^2\vartheta_{24}$$

APPearance

$$P^{\mathrm{SBL}}_{\overset{(-)}{\nu}_\alpha \to \overset{(-)}{\nu}_\beta} \simeq \sin^2 2\vartheta_{\alpha\beta} \sin^2\left(\frac{\Delta m^2_{41} L}{4E}\right)$$

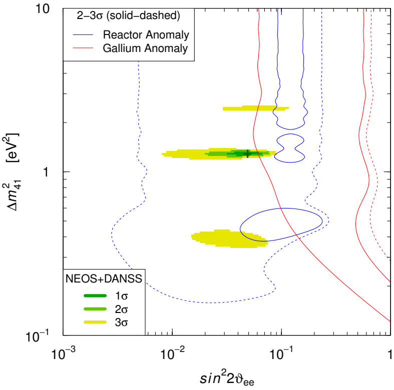
$$\sin^2 2\vartheta_{\alpha\beta} = 4|U_{\alpha4}|^2|U_{\beta4}|^2$$

$\overset{(-)}{\nu}_\mu \to \overset{(-)}{\nu}_e$ — LSND, MiniBooNE, KARMEN, OPERA, ...

$$\sin^2 2\vartheta_{e\mu} = 4|U_{e4}|^2|U_{\mu4}|^2$$

quadratically suppressed!

for small $|U_{e4}|^2$, $|U_{\mu4}|^2$


2–3σ (solid–dashed)
Reactor Anomaly
Gallium Anomaly
NEOS+DANSS
1σ
2σ
3σ
10
1
10^{-1}
Δm^2_{41} [eV^2]
10^{-3}
10^{-2}
10^{-1}
1
$sin^2 2\vartheta_{ee}$

Gallium anomaly

[SAGE, 2006][Laveder, 2007][Giunti&Laveder, 2011]

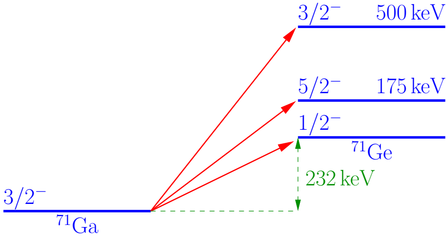
$L \simeq 1.9$ m $L \simeq 0.6$ m

Gallium radioactive source experiments: GALLEX and SAGE

ν_e sources: $e^- + {}^{51}\text{Cr} \rightarrow {}^{51}\text{V} + \nu_e$ $e^- + {}^{37}\text{Ar} \rightarrow {}^{37}\text{Cl} + \nu_e$

$E \simeq 0.75$ MeV $E \simeq 0.81$ MeV

In the detector: $\nu_e + {}^{71}\text{Ga} \rightarrow {}^{71}\text{Ge} + e^-$



cross sections of
the transitions from

[Krofcheck et al., PRL 55 (1985) 1051]

[Frekers et al., PLB 706 (2011) 134]

Gallium anomaly

[SAGE, 2006][Laveder, 2007][Giunti&Laveder, 2011]

$L \simeq 1.9$ m $\quad L \simeq 0.6$ m

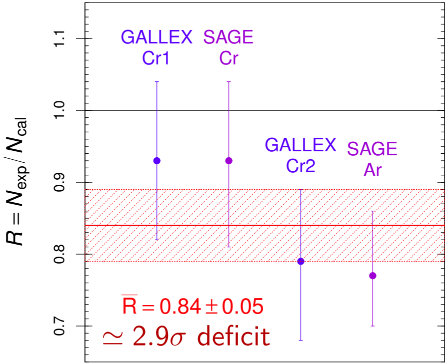
Gallium radioactive source experiments: GALLEX and SAGE

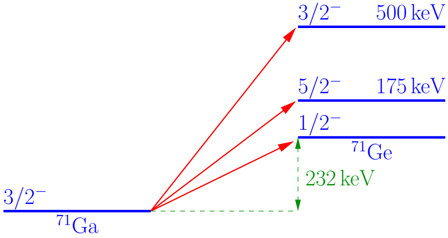
ν_e sources: $e^- + {}^{51}\text{Cr} \to {}^{51}\text{V} + \nu_e$ $\quad e^- + {}^{37}\text{Ar} \to {}^{37}\text{Cl} + \nu_e$

$E \simeq 0.75$ MeV $\quad E \simeq 0.81$ MeV

In the detector: $\nu_e + {}^{71}\text{Ga} \to {}^{71}\text{Ge} + e^-$

Test detection of solar ν_e





cross sections of the transitions from

[Krofcheck et al., PRL 55 (1985) 1051]

[Frekers et al., PLB 706 (2011) 134]

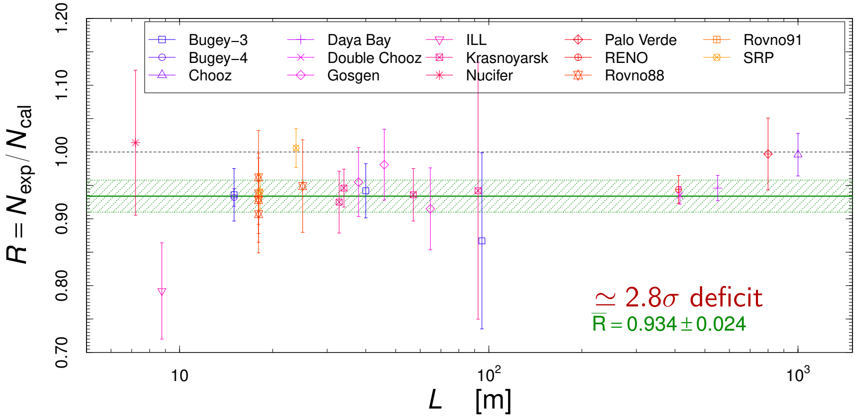
Reactor Antineutrino Anomaly (RAA)

[PRD 83 (2011) 073006]

2011: new reactor $\bar{\nu}_e$ fluxes by Huber and Mueller+ (HM)

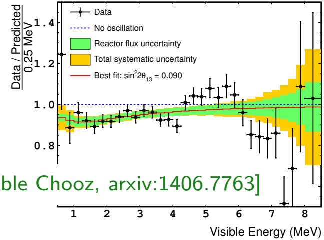
[Huber, PRC 84 (2011) 024617] [Mueller et al., PRC 83 (2011) 054615]

Previous reactor rates evaluated with new fluxes ⇒ deficit

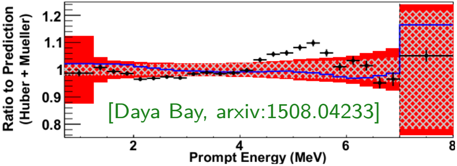


Suppression at detector due to active-sterile oscillations?

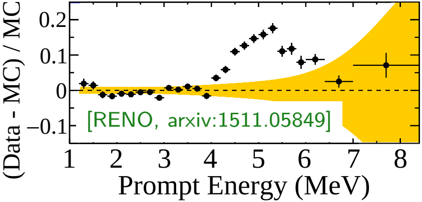
Can we trust the HM fluxes?



[Double Chooz, arxiv:1406.7763]



[Daya Bay, arxiv:1508.04233]



[RENO, arxiv:1511.05849]

2014:

bump in the spectrum around 5 MeV!

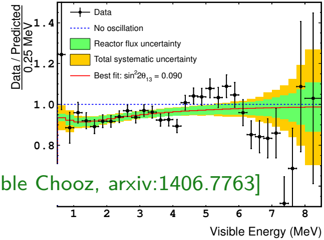
cannot be explained by SBL oscillations

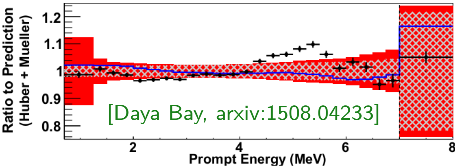
(averaged at the observed distances)

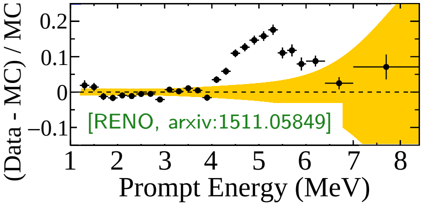
many attempts of possible explanations,

how to clarify the issue?

Can we trust the HM fluxes?







2014:

bump in the spectrum around 5 MeV!

cannot be explained by SBL oscillations

(averaged at the observed distances)

many attempts of possible explanations, how to clarify the issue?

Model independent information!

(i.e. take ratio of spectra at different distances)

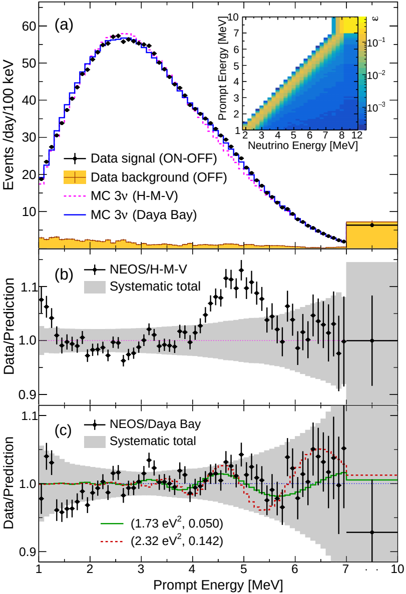
$$\Phi_1 = \Phi_0(E) f(L_1, E) \quad \Phi_2 = \Phi_0(E) f(L_2, E)$$

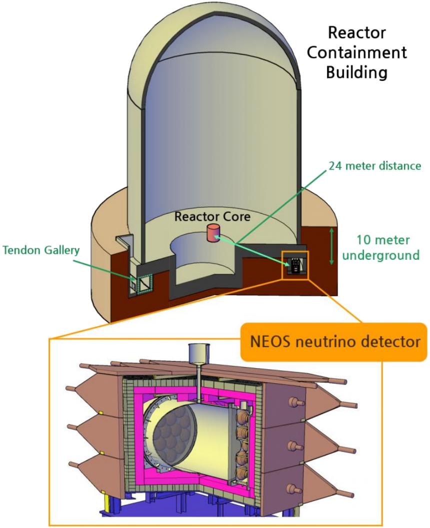
$$\Phi_1/\Phi_2 = f(L_1, E)/f(L_2, E)$$

NEOS

[PRL 118 (2017) 121802]

Single detector experiment

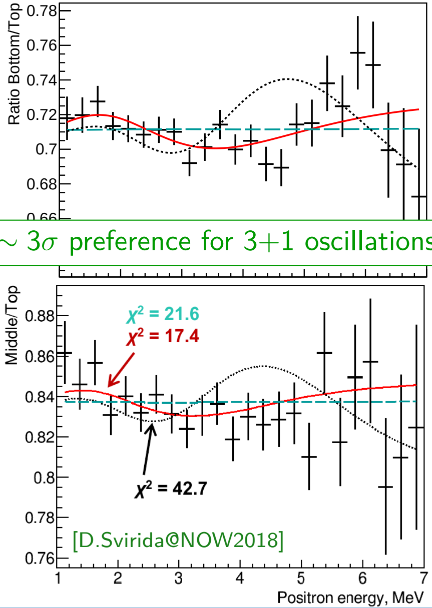




Ratio to DayaBay measurement to be model independent

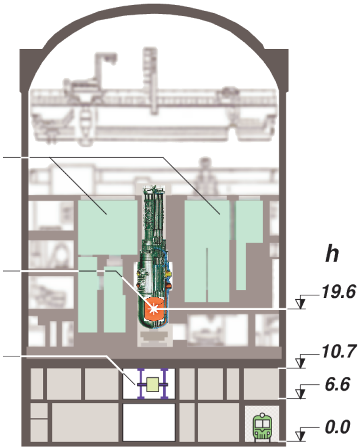
DANSS

Single movable detector



$\sim 3\sigma$ preference for 3+1 oscillations

[D.Svirida@NOW2018]

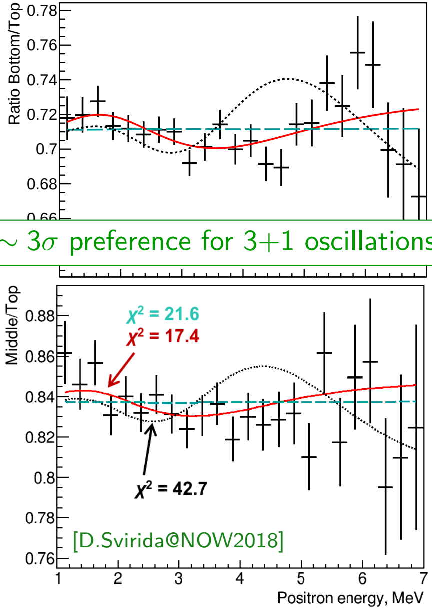


Detector can be at $\sim$ 10.5, $\sim$ 11.5 or $\sim$ 12.5 m from reactor core

DANSS

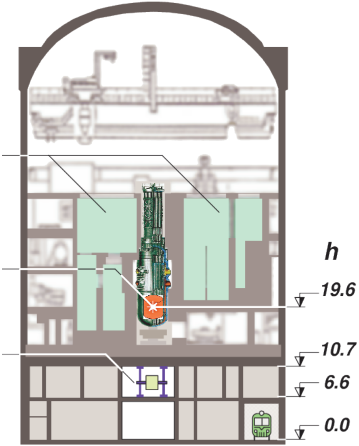
[PLB 787 (2018) 56]

Single movable detector



$\sim 3\sigma$ preference for 3+1 oscillations

[D.Svirida@NOW2018]



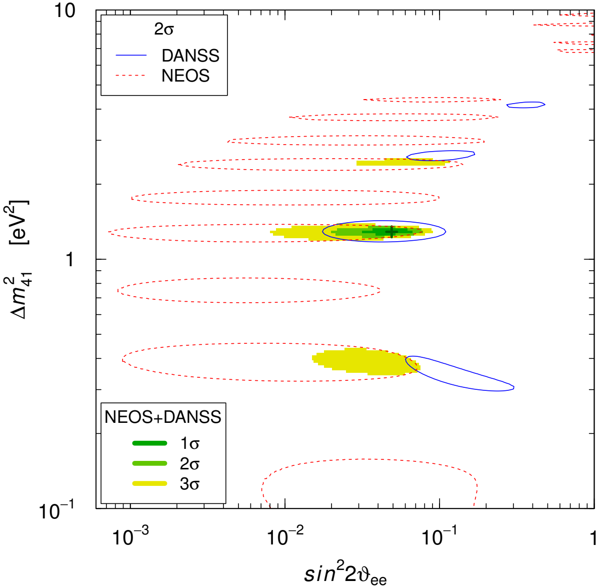
Detector can be at $\sim$ 10.5, $\sim$ 11.5 or $\sim$ 12.5 m from reactor core

see later for 2019 update!

Model-independent fit of $\overset{(-)}{\nu}_e$ DIS (2018)

[SG+, PLB 782 (2018) 13]

NEOS + DANSS



The NEOS and DANSS region perfectly overlap at

$$\Delta m^2_{41} \simeq 1.3 \text{ eV}^2$$

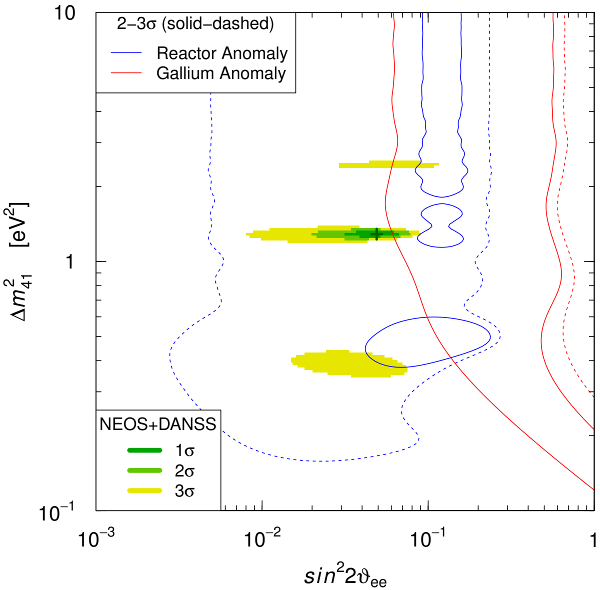
$$\sin^2 2\vartheta_{ee} \simeq 0.05$$

$$\sin^2 \vartheta_{14} \simeq 0.01$$

Model-independent fit of $\overset{(-)}{\nu}_e$ DIS (2018)

[SG+, PLB 782 (2018) 13]

DANSS + NEOS + RAA + Gallium

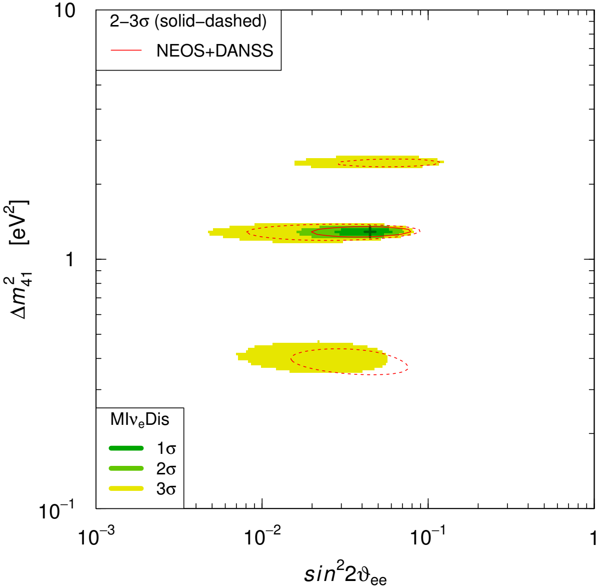


DANSS + NEOS do not agree with Gallium and RAA

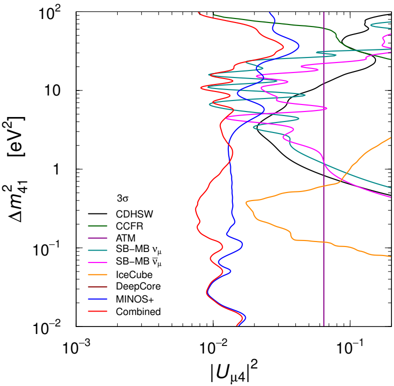
Model-independent fit of $\overset{(-)}{\nu}_e$ DIS (2018)

[SG+, PLB 782 (2018) 13]

All data:



Fit dominated by DANSS + NEOS


10^2
10
1
10^{-1}
10^{-2}
Δm^2_{41} [eV^2]
10^{-3}
10^{-2}
10^{-1}
$|U_{\mu 4}|^2$
3σ
CDHSW
CCFR
ATM
SB–MB ν_μ
SB–MB $\bar{\nu}_\mu$
IceCube
DeepCore
MINOS+
Combined

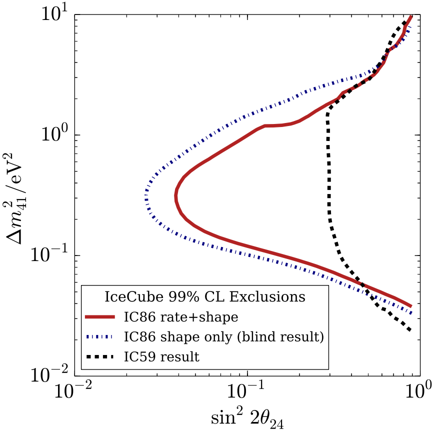
IceCube and DeepCore

IceCube

$\mathcal{O}(10\text{ km}) \lesssim L \lesssim \mathcal{O}(10^4\text{ km})$

$\sim 2 \times 10^4$ High energy μ events

320 GeV $<$ E $<$ 20 TeV



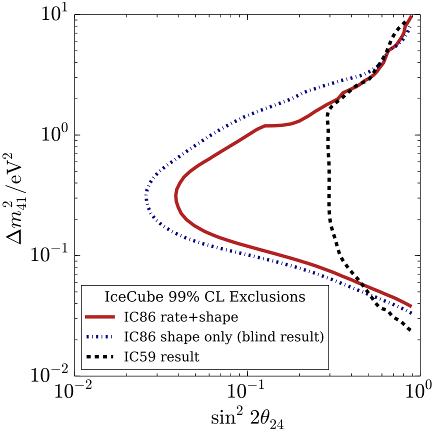
[PRL 117 (2016) 071801]

IceCube and DeepCore

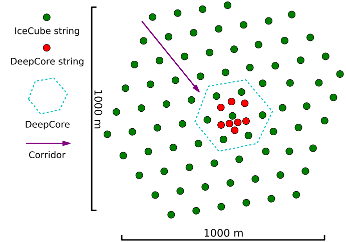
IceCube $\mathcal{O}(10\text{ km}) \lesssim L \lesssim \mathcal{O}(10^4\text{ km})$ DeepCore

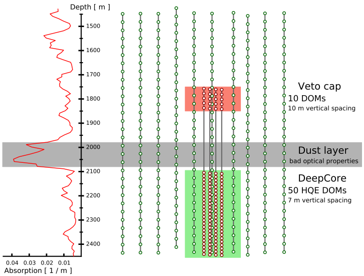
$\sim 2 \times 10^4$ High energy μ events

$320\text{ GeV} < E < 20\text{ TeV}$



[PRL 117 (2016) 071801]





IceCube and DeepCore

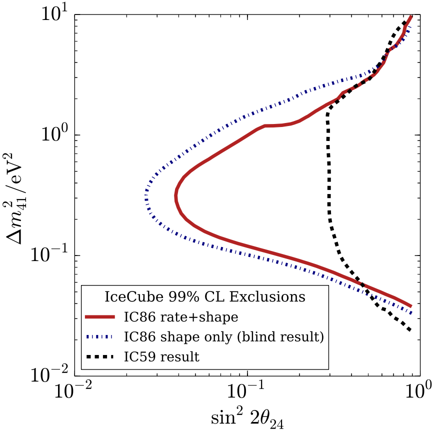
IceCube $\mathcal{O}(10\text{ km}) \lesssim L \lesssim \mathcal{O}(10^4\text{ km})$ DeepCore

$\sim 2 \times 10^4$ High energy μ events

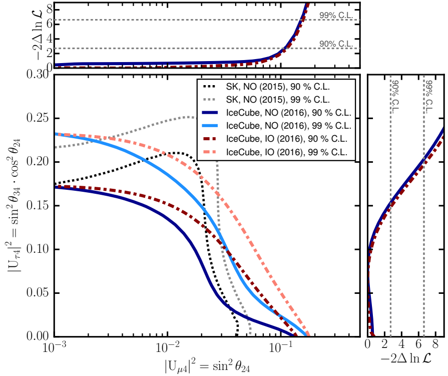
$320\text{ GeV} < E < 20\text{ TeV}$

$\sim 5 \times 10^3$ tracklike events

$6\text{ GeV} \lesssim E \lesssim 60\text{ GeV}$



[PRL 117 (2016) 071801]



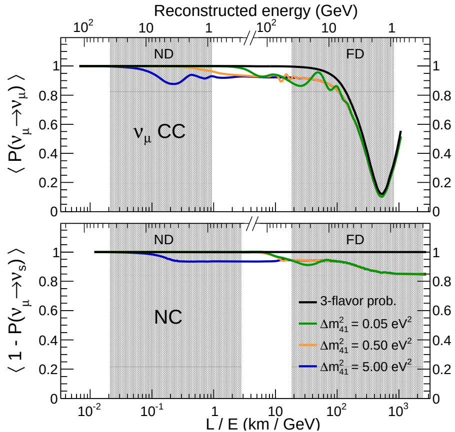
[PRD 95 (2017) 112002]

Both also constrain $|U_{\tau 4}|^2$

MINOS & MINOS+

Near (ND, $\simeq$ 500 m) and far (FD, $\simeq$ 800 km) detector

1 GeV $\lesssim E \lesssim$ 40 GeV, peak at 3 GeV



[PRL 117 (2016) 151803]:

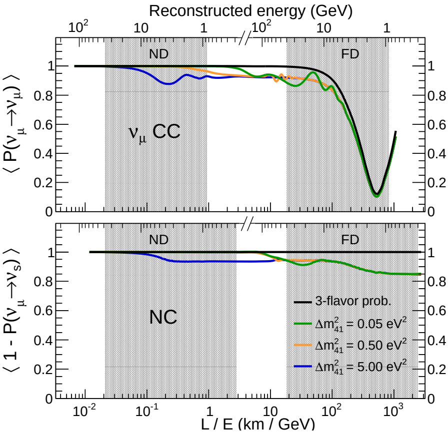
far-to-near ratio

[PRL 122 (2019) 091803]:

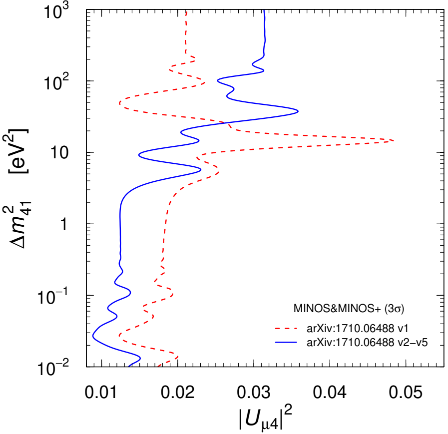
full two-detectors fit

Near (ND, $\simeq$ 500 m) and far (FD, $\simeq$ 800 km) detector

1 GeV $\lesssim E \lesssim$ 40 GeV, peak at 3 GeV



MINOS+ update:



[PRL 117 (2016) 151803]:

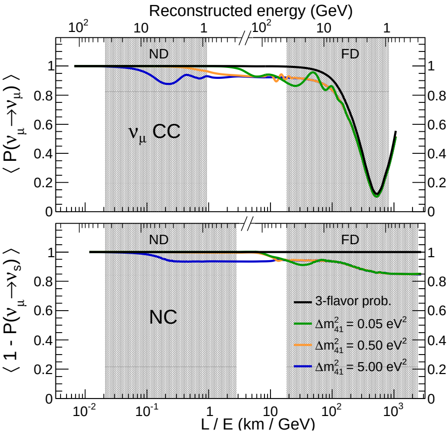
far-to-near ratio

[PRL 122 (2019) 091803]:

full two-detectors fit

Near (ND, $\simeq$ 500 m) and far (FD, $\simeq$ 800 km) detector

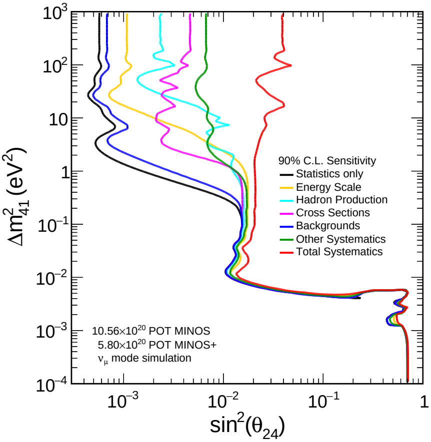
1 GeV $\lesssim E \lesssim$ 40 GeV, peak at 3 GeV



[PRL 117 (2016) 151803]: far-to-near ratio

[PRL 122 (2019) 091803]: full two-detectors fit

Systematics:

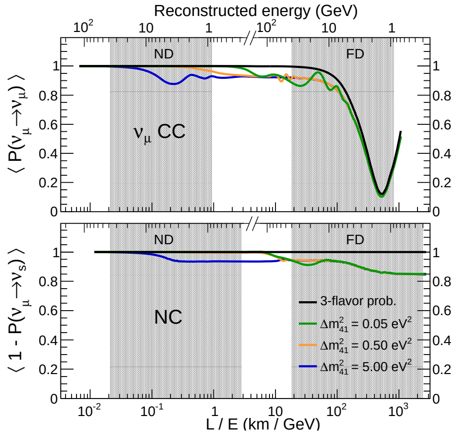


[PRL 122 (2019) 091803]

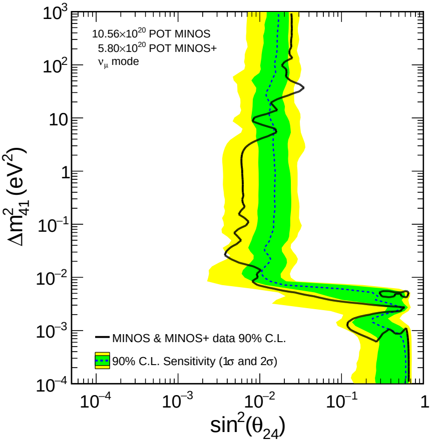
MINOS & MINOS+

Near (ND, $\simeq$ 500 m) and far (FD, $\simeq$ 800 km) detector

1 GeV $\lesssim E \lesssim$ 40 GeV, peak at 3 GeV



Sensitivity and exclusion limit:



[PRL 117 (2016) 151803]:

far-to-near ratio

[PRL 122 (2019) 091803]:

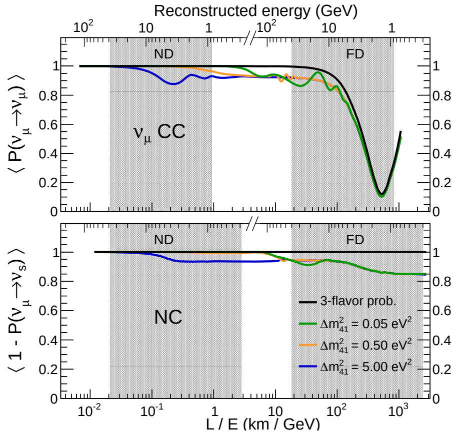
full two-detectors fit

[PRL 122 (2019) 091803]

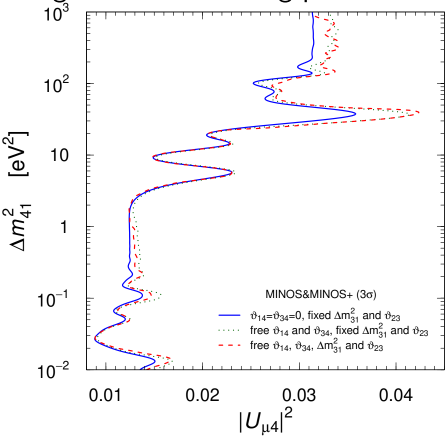
MINOS & MINOS+

Near (ND, $\simeq$ 500 m) and far (FD, $\simeq$ 800 km) detector

1 GeV $\lesssim E \lesssim$ 40 GeV, peak at 3 GeV



Marginalize over mixing parameters:



[PRL 117 (2016) 151803]: far-to-near ratio

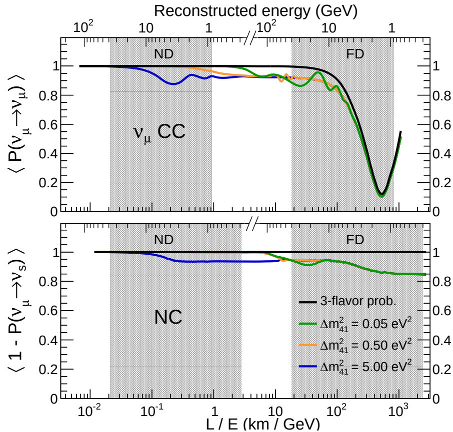
[PRL 122 (2019) 091803]: full two-detectors fit

[SG+, in preparation]

MINOS & MINOS+

Near (ND, $\simeq$ 500 m) and far (FD, $\simeq$ 800 km) detector

$1\text{ GeV} \lesssim E \lesssim 40\text{ GeV}$, peak at 3 GeV



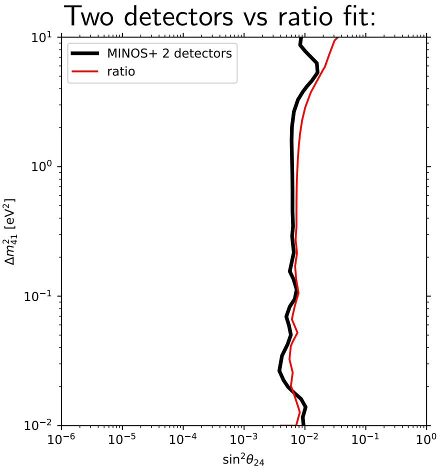
[PRL 117 (2016) 151803]:

far-to-near ratio

[PRL 122 (2019) 091803]:

full two-detectors fit

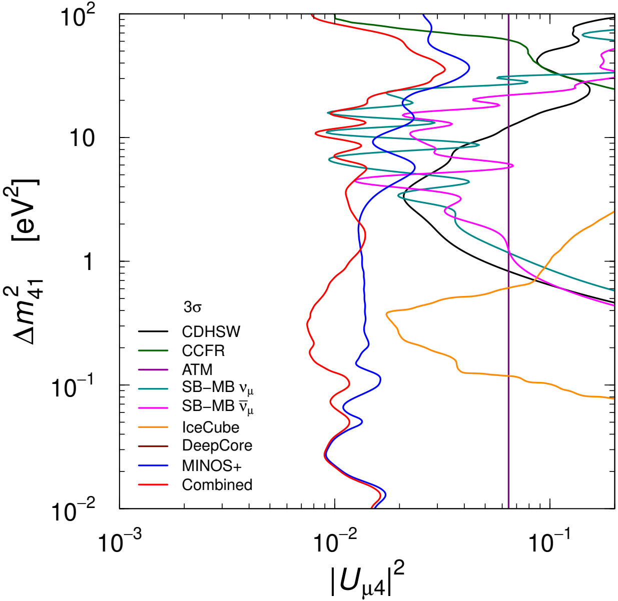
Two detectors vs ratio fit:



[SG+, in preparation]

Global fit of $\overset{(-)}{\nu}_\mu$ DIS

[SG+, in preparation]

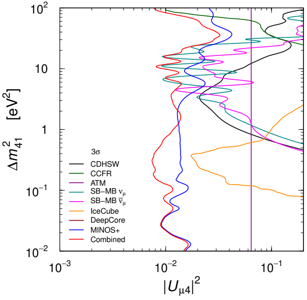


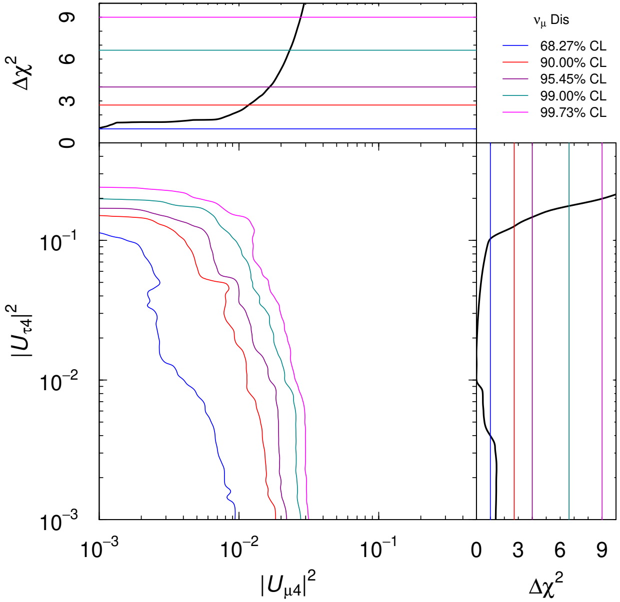
MINOS+ dominates at small Δm^2_{41}

IceCube important at $\Delta m^2_{41} \simeq 0.2$ eV^2

Global fit of $\overset{(-)}{\nu}_\mu$ DIS

[SG+, in preparation]





1 *Neutrino Oscillations - Some theory*

2 *Electron (anti)neutrino disappearence*

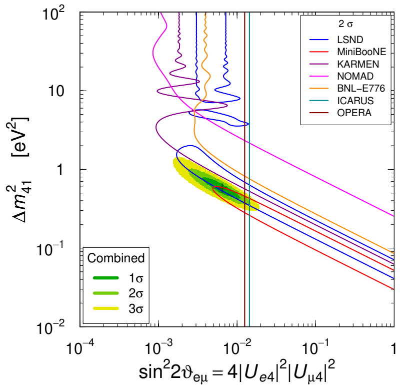
3 *Muon (anti)neutrino disappearence*

4 ***Electron (anti)neutrino appearence***

5 *Global fit*

6 *Cosmology*

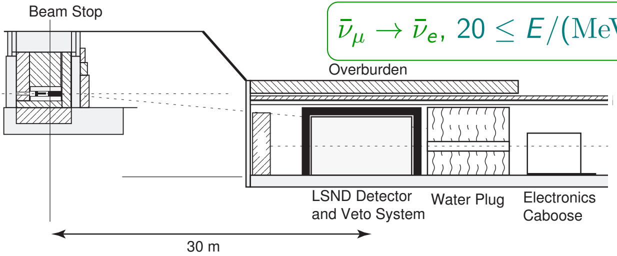
7 *Conclusions*



LSND

[PRD 64 (2001) 112007]

$\bar{\nu}_\mu \to \bar{\nu}_e$, $20 \leq E/(\mathrm{MeV}) \leq 52.8$



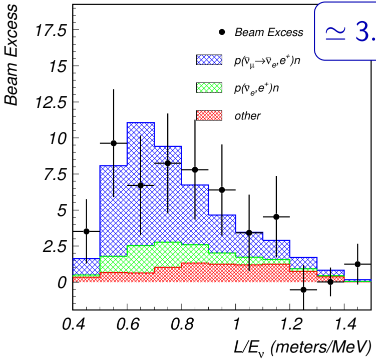
well known source of $\bar{\nu}_\mu$:

$$p + \mathrm{target} \to \pi^+ \xrightarrow{\mathrm{at\ rest}} \mu^+ + \nu_\mu$$

$$\mu^+ \xrightarrow{\mathrm{at\ rest}} e^+ + \nu_e + \bar{\nu}_\mu$$

$L \simeq 30$ m

$$\bar{\nu}_e + p \longrightarrow n + e^+$$



$\simeq 3.8\sigma$

No signal seen in KARMEN ($L \simeq 18$ m)

[PRD 65 (2002) 112001]

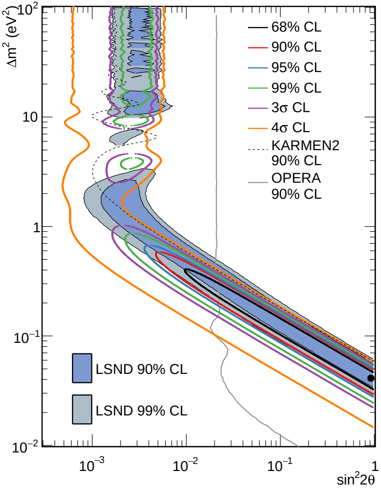
MiniBooNE

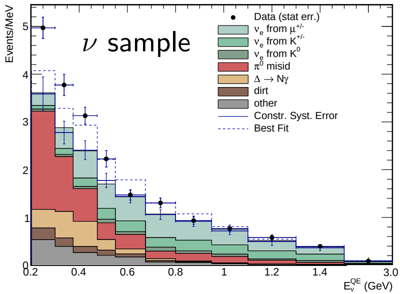
[PRL 121 (2018) 221801]

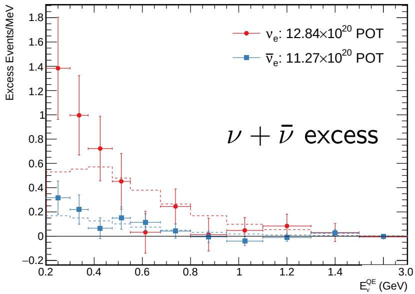
purpose: check LSND signal

$L \simeq 541$ m, 200 MeV $\leq E \lesssim 3$ GeV

no money, no near detector







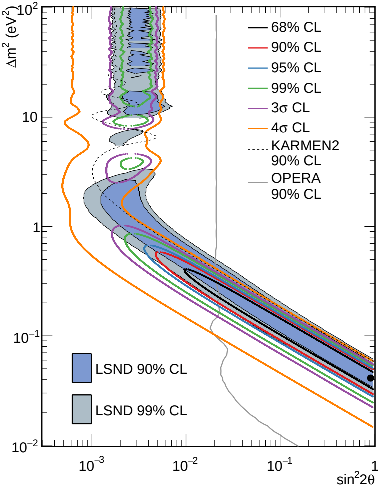
MiniBooNE

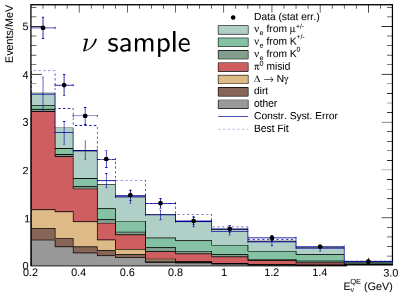
[PRL 121 (2018) 221801]

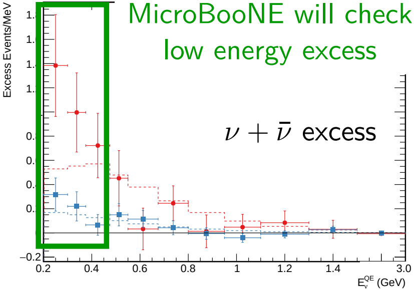
purpose: check LSND signal

$L \simeq 541$ m, 200 MeV $\leq E \lesssim$ 3 GeV

no money, no near detector



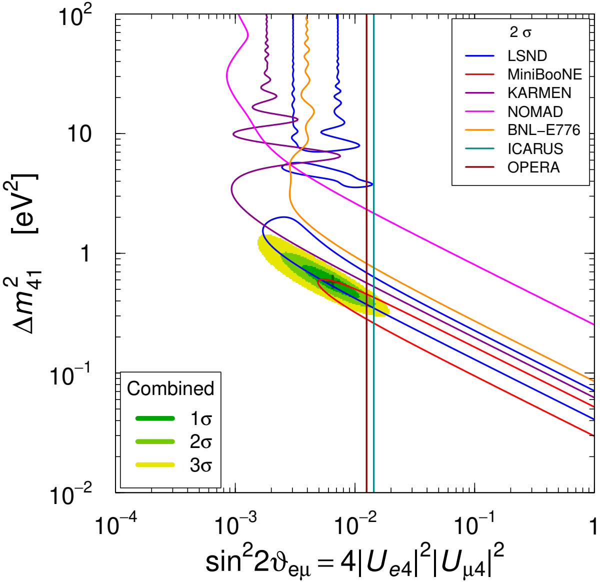




MicroBooNE will check low energy excess

Global fit of $\overset{(-)}{\nu}_\mu \to \overset{(-)}{\nu}_e$ APP

[SG+, in preparation]



with full MiniBooNE data

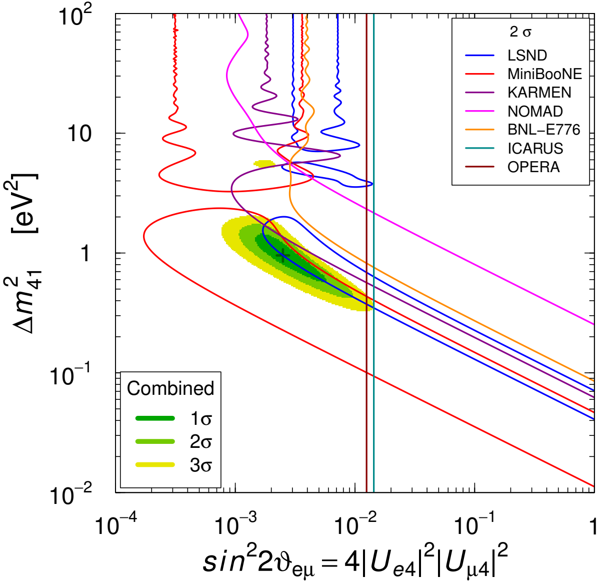
ICARUS and OPERA exclude MiniBooNE best fit

LSND and MiniBooNE only partially in agreement

KARMEN cuts part of LSND region

Global fit of $\overset{(-)}{\nu}_\mu \to \overset{(-)}{\nu}_e$ APP

[SG+, in preparation]

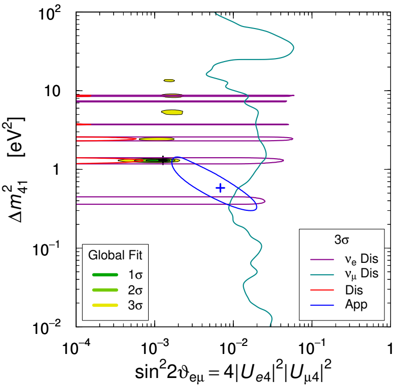


without MiniBooNE low energy bins

ICARUS and OPERA exclude MiniBooNE best fit

LSND and MiniBooNE only partially in agreement

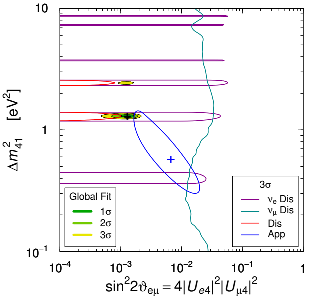
KARMEN cuts part of LSND region


10^2
10
1
10^{-1}
10^{-2}
Δm^2_{41} [eV^2]
10^{-4}
10^{-3}
10^{-2}
10^{-1}
1
$\sin^2 2\vartheta_{e\mu} = 4|U_{e4}|^2|U_{\mu4}|^2$
Global Fit
1σ
2σ
3σ
3σ
ν_e Dis
ν_μ Dis
Dis
App

APP – DIS tension in 2019

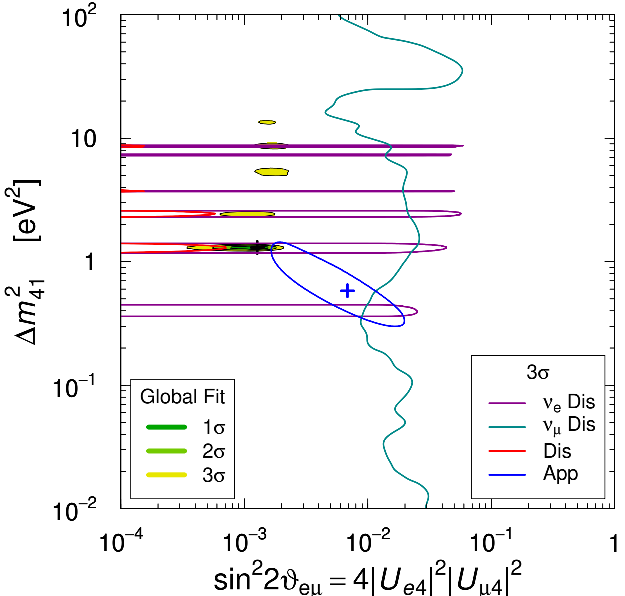
[SG+, in preparation]

Status just after Neutrino 2018:



MINOS+ update, new data including MiniBooNE (all bins)

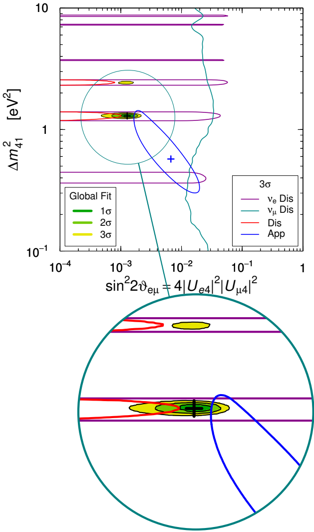
Status in early 2019



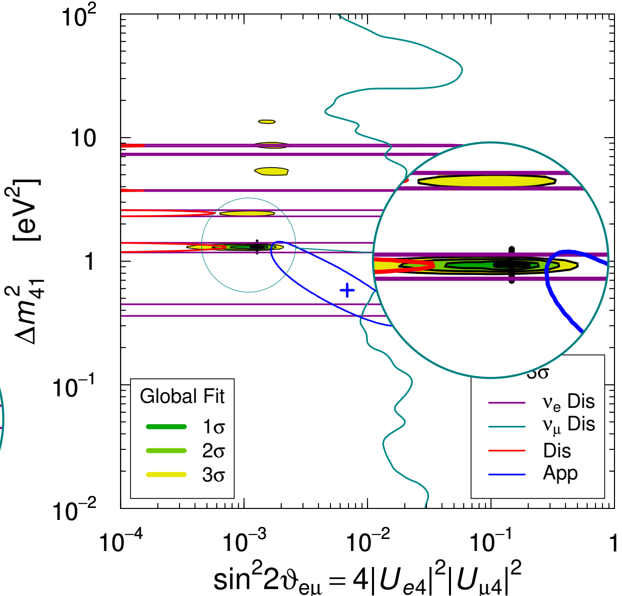
APP – DIS tension in 2019

[SG+, in preparation]

Status just after Neutrino 2018:

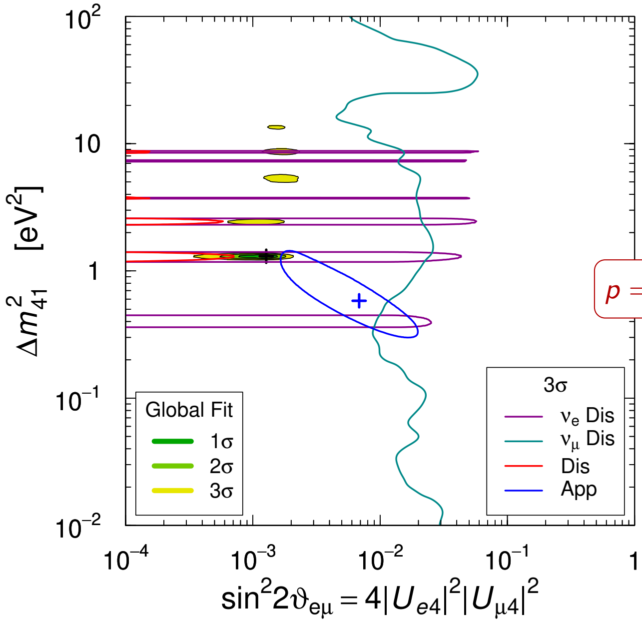


Status in early 2019



APP – DIS tension in 2019

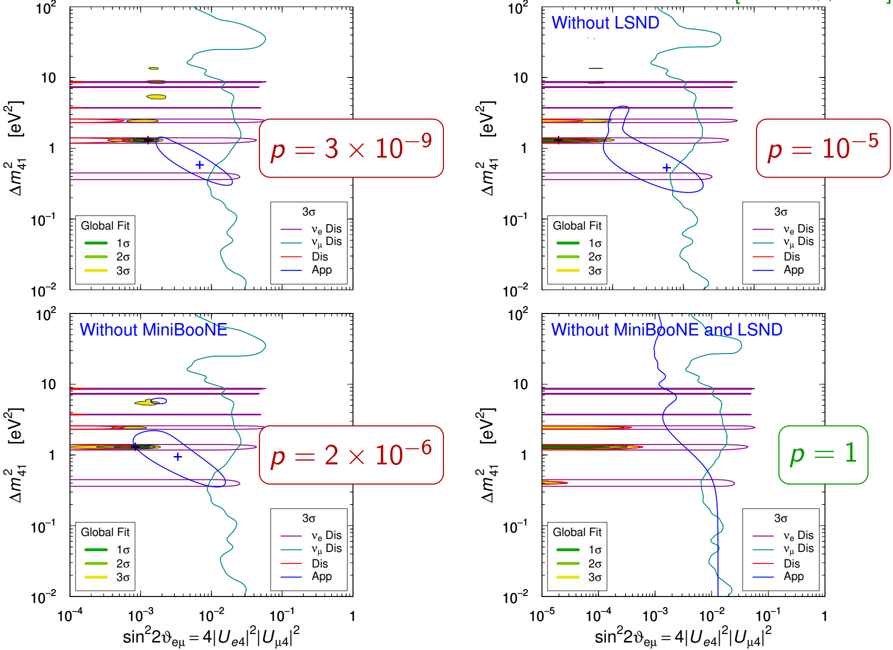
[SG+, in preparation]



$p = 3 \times 10^{-9}$

APP – DIS tension in 2019

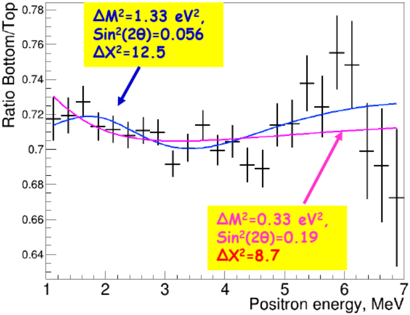
[SG+, in preparation]
see also [Dentler+, 2018]



DANSS in 2019

[Danilov@EPS-HEP, 2019]

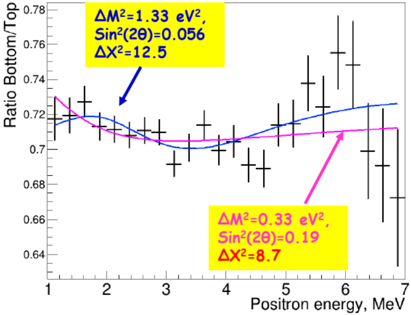
old data



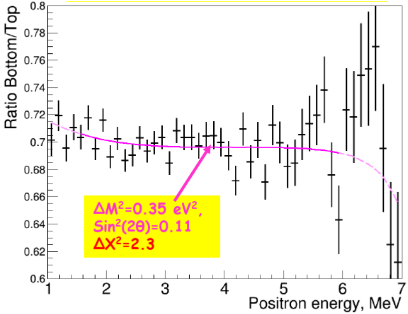
DANSS in 2019

[Danilov@EPS-HEP, 2019]

old data



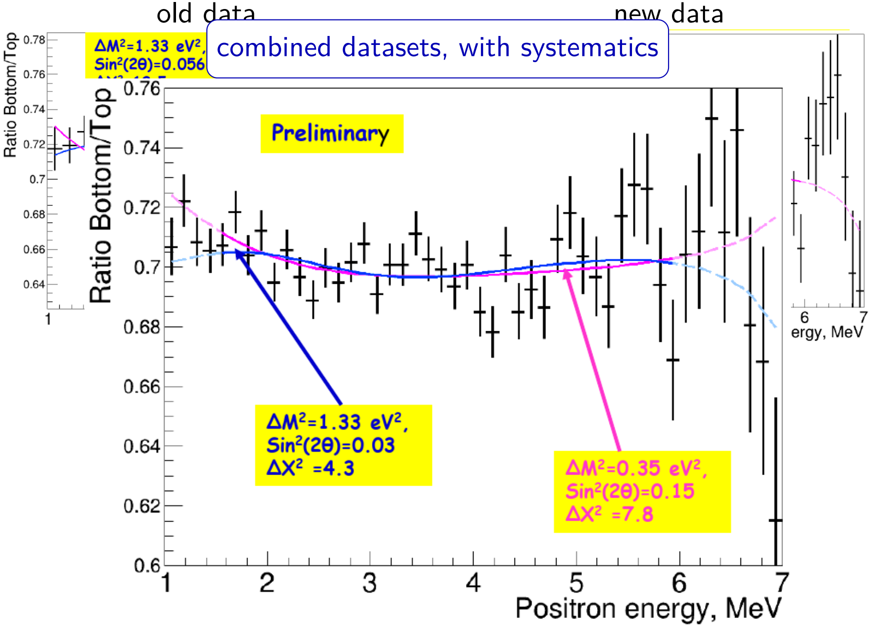
new data



New analysis also considers systematics!

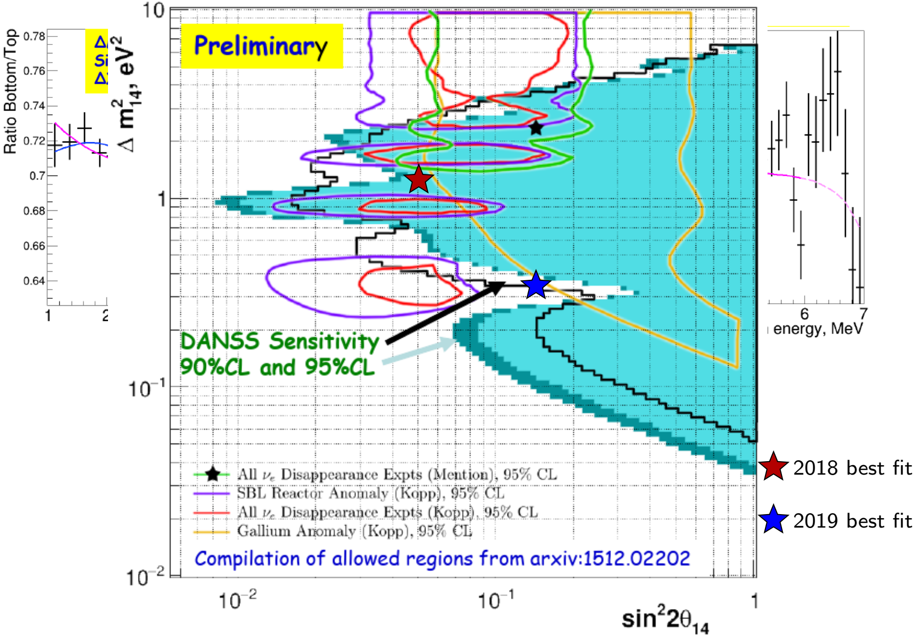
DANSS in 2019

[Danilov@EPS-HEP, 2019]



DANSS in 2019

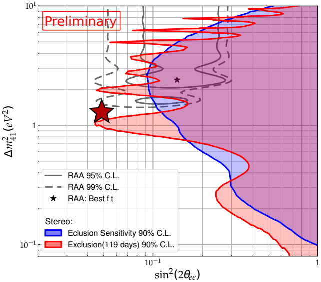
[Danilov@EPS-HEP, 2019]



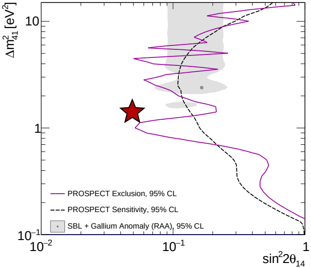
More to come...

★ = 2018 DANSS+NEOS best fit
[SG et al., PLB 782 (2018) 13]

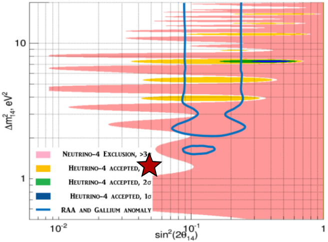
[STEREO, arxiv:1905.11896]



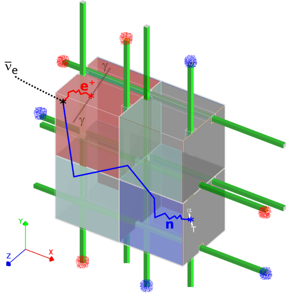
[PROSPECT, PRL 121 (2018) 251802]

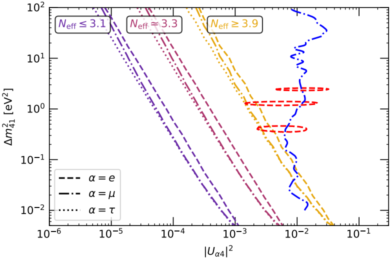


[Neutrino-4, PZETF 109 (2019) 209-218]



[SoLiD, JINST 13 (2018) P09005]




$N_{\rm eff} \leq 3.1$
$N_{\rm eff} \simeq 3.3$
$N_{\rm eff} \geq 3.9$
$\alpha = e$
$\alpha = \mu$
$\alpha = \tau$
Δm^2_{41} [eV^2]
$|U_{\alpha 4}|^2$

ν oscillations in the early universe

[SG+, JCAP 07 (2019) 014]

comoving coordinates: $a = 1/T \quad x \equiv m_e\, a \quad y \equiv p\, a \quad z \equiv T_\gamma\, a \quad w \equiv T_\nu\, a$

density matrix: $\varrho(x,y) = \begin{pmatrix} \varrho_{ee} \equiv f_{\nu_e} & \varrho_{e\mu} & \varrho_{e\tau} & \varrho_{es} \\ \varrho_{\mu e} & \varrho_{\mu\mu} \equiv f_{\nu_\mu} & \varrho_{\mu\tau} & \varrho_{\mu s} \\ \varrho_{\tau e} & \varrho_{\tau\mu} & \varrho_{\tau\tau} \equiv f_{\nu_\tau} & \varrho_{\tau s} \\ \varrho_{se} & \varrho_{s\mu} & \varrho_{s\tau} & \varrho_{ss} \equiv f_{\nu_s} \end{pmatrix}$

$$\frac{\mathrm{d}\varrho(y,x)}{\mathrm{d}x} = \sqrt{\frac{3m_{\mathrm{Pl}}^2}{8\pi\rho_{\mathrm{T}}}} \left\{ -i\frac{x^2}{m_e^3}\left[\frac{\mathbb{M}_{\mathrm{F}}}{2y} - \frac{8\sqrt{2}G_{\mathrm{F}}\, y m_e^6}{3x^6}\left(\frac{\mathbb{E}_\ell}{m_W^2} + \frac{\mathbb{E}_\nu}{m_Z^2}\right), \varrho\right] + \frac{m_e^3 G_{\mathrm{F}}^2}{(2\pi)^3 x^4 y^2}\mathcal{I}(\varrho) \right\}$$

m_{Pl} Planck mass – ρ_T total energy density – $m_{W,Z}$ mass of the W, Z bosons – G_{F} Fermi constant – $[.,.]$ commutator

ν oscillations in the early universe

[SG+, JCAP 07 (2019) 014]

comoving coordinates: $a = 1/T \quad x \equiv m_e\, a \quad y \equiv p\, a \quad z \equiv T_\gamma\, a \quad w \equiv T_\nu\, a$

$$\text{density matrix:} \quad \varrho(x,y) = \begin{pmatrix} \varrho_{ee} \equiv f_{\nu_e} & \varrho_{e\mu} & \varrho_{e\tau} & \varrho_{es} \\ \varrho_{\mu e} & \varrho_{\mu\mu} \equiv f_{\nu_\mu} & \varrho_{\mu\tau} & \varrho_{\mu s} \\ \varrho_{\tau e} & \varrho_{\tau\mu} & \varrho_{\tau\tau} \equiv f_{\nu_\tau} & \varrho_{\tau s} \\ \varrho_{se} & \varrho_{s\mu} & \varrho_{s\tau} & \varrho_{ss} \equiv f_{\nu_s} \end{pmatrix}$$

$$\frac{\mathrm{d}\varrho(y,x)}{\mathrm{d}x} = \sqrt{\frac{3 m_{\rm Pl}^2}{8\pi\rho_{\rm T}}} \left\{ -i \frac{x^2}{m_e^3} \left[\frac{\mathbb{M}_{\rm F}}{2y} - \frac{8\sqrt{2} G_{\rm F}\, y\, m_e^6}{3x^6} \left(\frac{\mathbb{E}_\ell}{m_W^2} + \frac{\mathbb{E}_\nu}{m_Z^2} \right), \varrho \right] + \frac{m_e^3 G_{\rm F}^2}{(2\pi)^3 x^4 y^2} \mathcal{I}(\varrho) \right\}$$

$m_{\rm Pl}$ Planck mass – ρ_T total energy density – $m_{W,Z}$ mass of the W, Z bosons – $G_{\rm F}$ Fermi constant – $[.,.]$ commutator

$$\mathbb{M}_{\rm F} = U \mathbb{M} U^\dagger$$

$$\mathbb{M} = \mathrm{diag}(m_1^2, \ldots, m_N^2)$$

$$U = R^{34} R^{24} R^{14} R^{23} R^{13} R^{12} \qquad \text{e.g. } R^{14} = \begin{pmatrix} \cos\theta_{14} & 0 & 0 & \sin\theta_{14} \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ -\sin\theta_{14} & 0 & 0 & \cos\theta_{14} \end{pmatrix}$$

$$|U|^2 = \begin{pmatrix} \ldots & \ldots & \ldots & \sin^2\theta_{14} \\ \ldots & \ldots & \ldots & \cos^2\theta_{14}\sin^2\theta_{24} \\ \ldots & \ldots & \ldots & \cos^2\theta_{14}\cos^2\theta_{24}\sin^2\theta_{34} \\ \ldots & \ldots & \ldots & \cos^2\theta_{14}\cos^2\theta_{24}\cos^2\theta_{34} \end{pmatrix}$$

ν oscillations in the early universe

[SG+, JCAP 07 (2019) 014]

comoving coordinates: $a = 1/T \quad x \equiv m_e\, a \quad y \equiv p\, a \quad z \equiv T_\gamma\, a \quad w \equiv T_\nu\, a$

density matrix: $$\varrho(x,y) = \begin{pmatrix} \varrho_{ee} \equiv f_{\nu_e} & \varrho_{e\mu} & \varrho_{e\tau} & \varrho_{es} \\ \varrho_{\mu e} & \varrho_{\mu\mu} \equiv f_{\nu_\mu} & \varrho_{\mu\tau} & \varrho_{\mu s} \\ \varrho_{\tau e} & \varrho_{\tau\mu} & \varrho_{\tau\tau} \equiv f_{\nu_\tau} & \varrho_{\tau s} \\ \varrho_{se} & \varrho_{s\mu} & \varrho_{s\tau} & \varrho_{ss} \equiv f_{\nu_s} \end{pmatrix}$$

$$\frac{\mathrm{d}\varrho(y,x)}{\mathrm{d}x} = \sqrt{\frac{3 m_{\mathrm{Pl}}^2}{8\pi\rho_{\mathrm{T}}}} \left\{ -i\frac{x^2}{m_e^3}\left[\frac{\mathbb{M}_{\mathrm{F}}}{2y} - \frac{8\sqrt{2}G_{\mathrm{F}}\, y m_e^6}{3x^6}\left(\frac{\mathbb{E}_\ell}{m_W^2} + \frac{\mathbb{E}_\nu}{m_Z^2}\right), \varrho\right] + \frac{m_e^3 G_{\mathrm{F}}^2}{(2\pi)^3 x^4 y^2}\mathcal{I}(\varrho) \right\}$$

m_{Pl} Planck mass – ρ_T total energy density – $m_{W,Z}$ mass of the W, Z bosons – G_{F} Fermi constant – [., .] commutator

$$\mathbb{M}_{\mathrm{F}} = U \mathbb{M} U^\dagger$$

$\mathbb{E}_\ell = \mathrm{diag}(\rho_e, \rho_\mu, 0, 0)$ $\qquad$ $\mathbb{E}_\nu = S_a \left(\int dy y^3 \varrho\right) S_a$ $\quad$ with $S_a = \mathrm{diag}(1,1,1,0)$

lepton densities $\qquad$ neutrino densities $\qquad$ (only for active neutrinos)

take into account matter effects in oscillations

ν oscillations in the early universe

[SG+, JCAP 07 (2019) 014]

comoving coordinates: $a = 1/T \quad x \equiv m_e\, a \quad y \equiv p\, a \quad z \equiv T_\gamma\, a \quad w \equiv T_\nu\, a$

density matrix: $$\varrho(x,y) = \begin{pmatrix} \varrho_{ee} \equiv f_{\nu_e} & \varrho_{e\mu} & \varrho_{e\tau} & \varrho_{es} \\ \varrho_{\mu e} & \varrho_{\mu\mu} \equiv f_{\nu_\mu} & \varrho_{\mu\tau} & \varrho_{\mu s} \\ \varrho_{\tau e} & \varrho_{\tau\mu} & \varrho_{\tau\tau} \equiv f_{\nu_\tau} & \varrho_{\tau s} \\ \varrho_{se} & \varrho_{s\mu} & \varrho_{s\tau} & \varrho_{ss} \equiv f_{\nu_s} \end{pmatrix}$$

$$\frac{\mathrm{d}\varrho(y,x)}{\mathrm{d}x} = \sqrt{\frac{3m_{\mathrm{Pl}}^2}{8\pi\rho_{\mathrm{T}}}} \left\{ -i\frac{x^2}{m_e^3}\left[\frac{\mathbb{M}_{\mathrm{F}}}{2y} - \frac{8\sqrt{2}G_{\mathrm{F}}\, y m_e^6}{3x^6}\left(\frac{\mathbb{E}_\ell}{m_W^2} + \frac{\mathbb{E}_\nu}{m_Z^2}\right), \varrho\right] + \frac{m_e^3 G_{\mathrm{F}}^2}{(2\pi)^3 x^4 y^2}\mathcal{I}(\varrho) \right\}$$

m_{Pl} Planck mass – ρ_T total energy density – $m_{W,Z}$ mass of the W, Z bosons – G_{F} Fermi constant – $[.,.]$ commutator

$$\mathbb{M}_{\mathrm{F}} = U \mathbb{M} U^\dagger \qquad \mathbb{E}_\ell = \mathrm{diag}(\rho_e, \rho_\mu, 0, 0) \qquad \mathbb{E}_\nu = S_a \left(\int dy y^3 \varrho\right) S_a$$

$\mathcal{I}(\varrho)$ collision integrals

take into account neutrino-electron scattering and pair annihilation

2D integrals over the momentum, take most of the computation time

ν oscillations in the early universe

[SG+, JCAP 07 (2019) 014]

comoving coordinates: $a = 1/T \quad x \equiv m_e\, a \quad y \equiv p\, a \quad z \equiv T_\gamma\, a \quad w \equiv T_\nu\, a$

density matrix: $$\varrho(x,y) = \begin{pmatrix} \varrho_{ee} \equiv f_{\nu_e} & \varrho_{e\mu} & \varrho_{e\tau} & \varrho_{es} \\ \varrho_{\mu e} & \varrho_{\mu\mu} \equiv f_{\nu_\mu} & \varrho_{\mu\tau} & \varrho_{\mu s} \\ \varrho_{\tau e} & \varrho_{\tau\mu} & \varrho_{\tau\tau} \equiv f_{\nu_\tau} & \varrho_{\tau s} \\ \varrho_{se} & \varrho_{s\mu} & \varrho_{s\tau} & \varrho_{ss} \equiv f_{\nu_s} \end{pmatrix}$$

$$\frac{\mathrm{d}\varrho(y,x)}{\mathrm{d}x} = \sqrt{\frac{3m_{\mathrm{Pl}}^2}{8\pi\rho_{\mathrm{T}}}}\left\{-i\frac{x^2}{m_e^3}\left[\frac{\mathbb{M}_{\mathrm{F}}}{2y} - \frac{8\sqrt{2}G_{\mathrm{F}}\,y m_e^6}{3x^6}\left(\frac{\mathbb{E}_\ell}{m_W^2}+\frac{\mathbb{E}_\nu}{m_Z^2}\right),\varrho\right] + \frac{m_e^3 G_{\mathrm{F}}^2}{(2\pi)^3 x^4 y^2}\mathcal{I}(\varrho)\right\}$$

m_{Pl} Planck mass – ρ_T total energy density – $m_{W,Z}$ mass of the W, Z bosons – G_{F} Fermi constant – $[.,.]$ commutator

$$\mathbb{M}_{\mathrm{F}} = U\mathbb{M}U^\dagger \qquad \mathbb{E}_\ell = \mathrm{diag}(\rho_e, \rho_\mu, 0, 0) \qquad \mathbb{E}_\nu = S_a\left(\int dy y^3 \varrho\right) S_a$$

$\mathcal{I}(\varrho)$ collision integrals

from continuity equation

$\dot{\rho} = -3H(\rho + P)$

$$\frac{\mathrm{d}z}{\mathrm{d}x} = \frac{\displaystyle\sum_{\ell=e,\mu}\left[\frac{r_\ell^2}{r}J(r_\ell)\right] + G_1(r) - \frac{1}{2\pi^2 z^3}\int_0^\infty dy\, y^3 \sum_{\alpha=e}^{s}\frac{\mathrm{d}\varrho_{\alpha\alpha}}{\mathrm{d}x}}{\displaystyle\sum_{\ell=e,\mu}\left[r_\ell^2 J(r_\ell) + Y(r_\ell)\right] + G_2(r) + \frac{2\pi^2}{15}}$$

$r = x/z$, $r_\ell = m_\ell/m_e\, r$ $\quad$ $J(r)$, $Y(r)$ from non-relativistic transition of $e^\pm$, $\mu^\pm$

$G_1(r)$ and $G_2(r)$ from electromagnetic corrections

ν oscillations in the early universe

[SG+, JCAP 07 (2019) 014]

comoving coordinates: $a = 1/T \quad x \equiv m_e\, a \quad y \equiv p\, a \quad z \equiv T_\gamma\, a \quad w \equiv T_\nu\, a$

density matrix: $$\varrho(x,y) = \begin{pmatrix} \varrho_{ee} \equiv f_{\nu_e} & \varrho_{e\mu} & \varrho_{e\tau} & \varrho_{es} \\ \varrho_{\mu e} & \varrho_{\mu\mu} \equiv f_{\nu_\mu} & \varrho_{\mu\tau} & \varrho_{\mu s} \\ \varrho_{\tau e} & \varrho_{\tau\mu} & \varrho_{\tau\tau} \equiv f_{\nu_\tau} & \varrho_{\tau s} \\ \varrho_{se} & \varrho_{s\mu} & \varrho_{s\tau} & \varrho_{ss} \equiv f_{\nu_s} \end{pmatrix}$$

$$\frac{\mathrm{d}\varrho(y,x)}{\mathrm{d}x} = \sqrt{\frac{3m_{\mathrm{Pl}}^2}{8\pi\rho_{\mathrm{T}}}}\left\{-i\frac{x^2}{m_e^3}\left[\frac{\mathbb{M}_{\mathrm{F}}}{2y} - \frac{8\sqrt{2}G_{\mathrm{F}}\,y m_e^6}{3x^6}\left(\frac{\mathbb{E}_\ell}{m_W^2}+\frac{\mathbb{E}_\nu}{m_Z^2}\right),\varrho\right] + \frac{m_e^3 G_{\mathrm{F}}^2}{(2\pi)^3 x^4 y^2}\mathcal{I}(\varrho)\right\}$$

m_{Pl} Planck mass – ρ_T total energy density – $m_{W,Z}$ mass of the W, Z bosons – G_{F} Fermi constant – $[.,.]$ commutator

$$\mathbb{M}_{\mathrm{F}} = U \mathbb{M} U^\dagger \qquad \mathbb{E}_\ell = \mathrm{diag}(\rho_e, \rho_\mu, 0, 0) \qquad \mathbb{E}_\nu = S_a\left(\int dy y^3 \varrho\right) S_a$$

$\mathcal{I}(\varrho)$ collision integrals

from continuity equation
$\dot\rho = -3H(\rho + P)$

$$\frac{\mathrm{d}z}{\mathrm{d}x} = \frac{\sum_{\ell=e,\mu}\left[\frac{r_\ell^2}{r}J(r_\ell)\right] + G_1(r) - \frac{1}{2\pi^2 z^3}\int_0^\infty dy\, y^3 \sum_{\alpha=e}^{s}\frac{\mathrm{d}\varrho_{\alpha\alpha}}{\mathrm{d}x}}{\sum_{\ell=e,\mu}\left[r_\ell^2 J(r_\ell) + Y(r_\ell)\right] + G_2(r) + \frac{2\pi^2}{15}}$$

neutrino temperature w: same equation as z, but electrons always relativistic

ν oscillations in the early universe

[SG+, JCAP 07 (2019) 014]

comoving coordinates: $a = 1/T \quad x \equiv m_e\, a \quad y \equiv p\, a \quad z \equiv T_\gamma\, a \quad w \equiv T_\nu\, a$

density matrix: $$\varrho(x,y) = \begin{pmatrix} \varrho_{ee} \equiv f_{\nu_e} & \varrho_{e\mu} & \varrho_{e\tau} & \varrho_{es} \\ \varrho_{\mu e} & \varrho_{\mu\mu} \equiv f_{\nu_\mu} & \varrho_{\mu\tau} & \varrho_{\mu s} \\ \varrho_{\tau e} & \varrho_{\tau\mu} & \varrho_{\tau\tau} \equiv f_{\nu_\tau} & \varrho_{\tau s} \\ \varrho_{se} & \varrho_{s\mu} & \varrho_{s\tau} & \varrho_{ss} \equiv f_{\nu_s} \end{pmatrix}$$

$$\frac{\mathrm{d}\varrho(y,x)}{\mathrm{d}x} = \sqrt{\frac{3m_{\mathrm{Pl}}^2}{8\pi\rho_{\mathrm{T}}}} \left\{ -i\frac{x^2}{m_e^3}\left[\frac{\mathbb{M}_{\mathrm{F}}}{2y} - \frac{8\sqrt{2}G_{\mathrm{F}}\, y m_e^6}{3x^6}\left(\frac{\mathbb{E}_\ell}{m_W^2} + \frac{\mathbb{E}_\nu}{m_Z^2}\right), \varrho\right] + \frac{m_e^3 G_{\mathrm{F}}^2}{(2\pi)^3 x^4 y^2}\mathcal{I}(\varrho) \right\}$$

m_{Pl} Planck mass – ρ_T total energy density – $m_{W,Z}$ mass of the W, Z bosons – G_{F} Fermi constant – $[.,.]$ commutator

$$\mathbb{M}_{\mathrm{F}} = U \mathbb{M} U^\dagger \qquad \mathbb{E}_\ell = \mathrm{diag}(\rho_e, \rho_\mu, 0, 0) \qquad \mathbb{E}_\nu = S_a \left(\int dy y^3 \varrho\right) S_a$$

$\mathcal{I}(\varrho)$ collision integrals

from continuity equation $\dot{\rho} = -3H(\rho + P)$

$$\frac{\mathrm{d}z}{\mathrm{d}x} = \frac{\sum_{\ell=e,\mu}\left[\frac{r_\ell^2}{r}J(r_\ell)\right] + G_1(r) - \frac{1}{2\pi^2 z^3}\int_0^\infty dy\, y^3 \sum_{\alpha=e}^{s}\frac{\mathrm{d}\varrho_{\alpha\alpha}}{\mathrm{d}x}}{\sum_{\ell=e,\mu}\left[r_\ell^2 J(r_\ell) + Y(r_\ell)\right] + G_2(r) + \frac{2\pi^2}{15}}$$

neutrino temperature w: same equation as z, but electrons always relativistic

initial conditions: $\varrho_{\alpha\alpha}$ = Fermi-Dirac at $x_{\mathrm{in}} \simeq 0.001$, with $w = z \simeq 1$

comoving coordinates: $a = 1/T \quad x \equiv m_e\, a \quad y \equiv p\, a \quad z \equiv T_\gamma\, a \quad w \equiv T_\nu\, a$

density matrix: $$\varrho(x,y) = \begin{pmatrix} \varrho_{ee} \equiv f_{\nu_e} & \varrho_{e\mu} & \varrho_{e\tau} & \varrho_{es} \\ \varrho_{\mu e} & \varrho_{\mu\mu} \equiv f_{\nu_\mu} & \varrho_{\mu\tau} & \varrho_{\mu s} \\ \varrho_{\tau e} & \varrho_{\tau\mu} & \varrho_{\tau\tau} \equiv f_{\nu_\tau} & \varrho_{\tau s} \\ \varrho_{se} & \varrho_{s\mu} & \varrho_{s\tau} & \varrho_{ss} \equiv f_{\nu_s} \end{pmatrix}$$

$$\frac{\mathrm{d}\varrho(y,x)}{\mathrm{d}x} = \sqrt{\frac{3 m_{\mathrm{Pl}}^2}{8\pi\rho_{\mathrm{T}}}} \left\{ -i\frac{x^2}{m_e^3}\left[\frac{\mathbb{M}_{\mathrm{F}}}{2y} - \frac{8\sqrt{2} G_{\mathrm{F}}\, y\, m_e^6}{3x^6}\left(\frac{\mathbb{E}_\ell}{m_W^2} + \frac{\mathbb{E}_\nu}{m_Z^2}\right), \varrho\right] + \frac{m_e^3 G_{\mathrm{F}}^2}{(2\pi)^3 x^4 y^2}\mathcal{I}(\varrho) \right\}$$

m_{Pl} Planck mass ... total energy density ... mass of the W, Z bosons ... Fermi constant ... [,] commutator

FORTran-Evolved PrimordIAl Neutrino Oscillations (FortEPiaNO)

https://bitbucket.org/ahep_cosmo/fortepiano

will be public soon

from continuity equation $\dot\rho = -3H(\rho + P)$

$$\frac{\mathrm{d}z}{\mathrm{d}x} = \frac{\sum_{\ell=e,\mu}\left[\; r \;\ldots\;\right]\ldots) - \frac{1}{2\pi^2 z^3}\int_0^\infty dy\, y^3 \sum_{\alpha=e}^{s}\frac{\mathrm{d}\varrho_{\alpha\alpha}}{\mathrm{d}x}}{\sum_{\ell=e,\mu}\left[r_\ell^2 J(r_\ell) + Y(r_\ell)\right] + G_2(r) + \frac{2\pi^2}{15}}$$

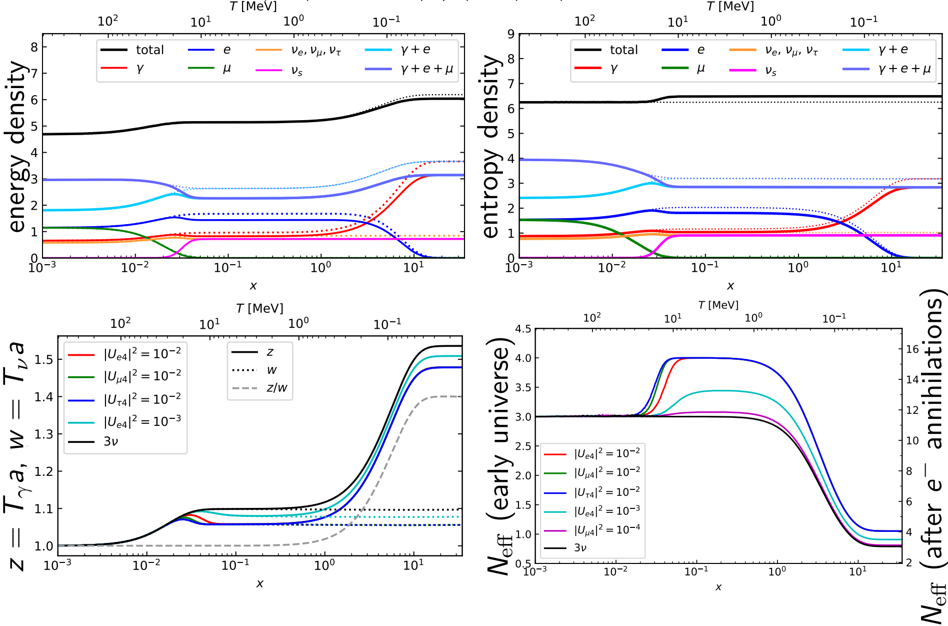
neutrino temperature w: same equation as z, but electrons always relativistic

initial conditions: $\varrho_{\alpha\alpha}$ = Fermi-Dirac at $x_{\mathrm{in}} \simeq 0.001$, with $w = z \simeq 1$

Energy, entropy, temperatures, $N_{\rm eff}$

[SG+, JCAP 07 (2019) 014]

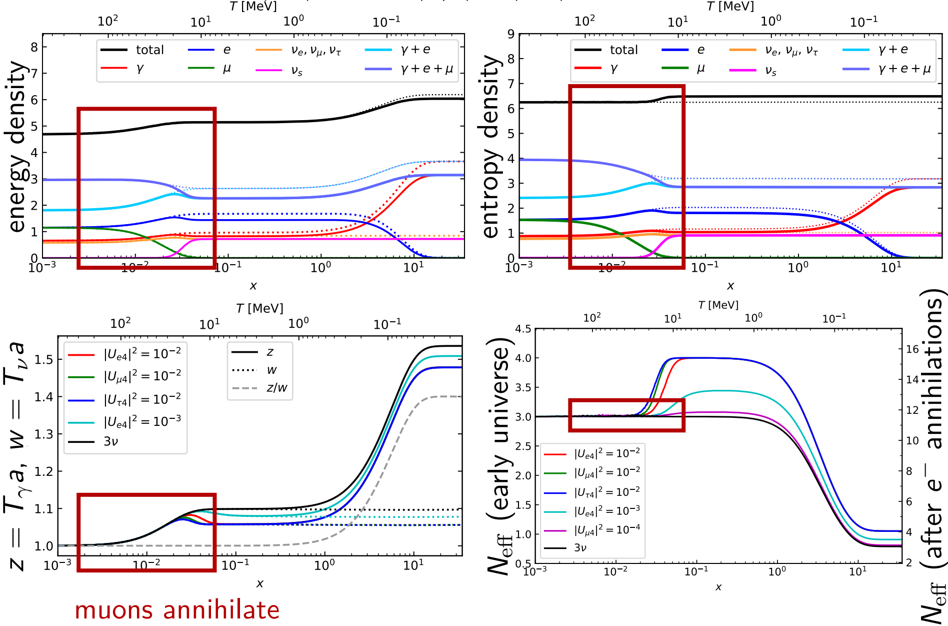
dashed: 3ν, solid: $|U_{e4}|^2 = 10^{-2}$, $|U_{\mu 4}|^2 = |U_{\tau 4}|^2 = 0$. $\Delta m^2_{41} = 1.29$ eV^2 always



Energy, entropy, temperatures, $N_{\rm eff}$

[SG+, JCAP 07 (2019) 014]

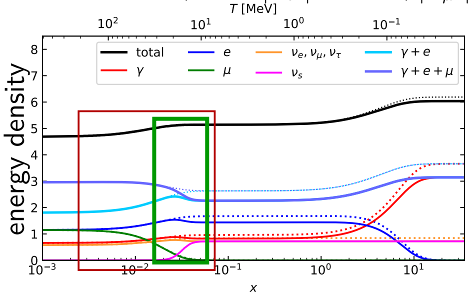
dashed: 3ν, solid: $|U_{e4}|^2 = 10^{-2}$, $|U_{\mu4}|^2 = |U_{\tau4}|^2 = 0$. $\Delta m^2_{41} = 1.29$ eV^2 always

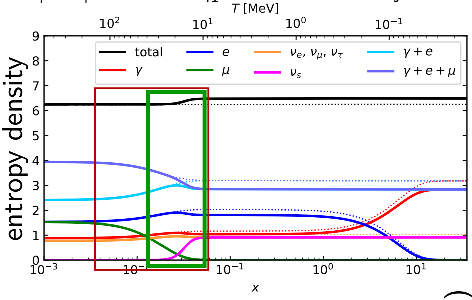


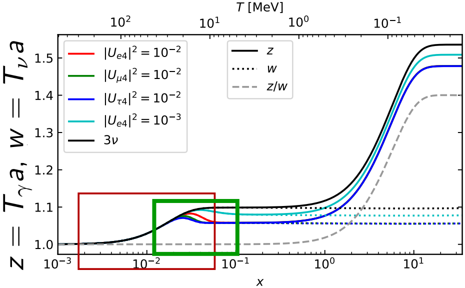
Energy, entropy, temperatures, $N_{\rm eff}$

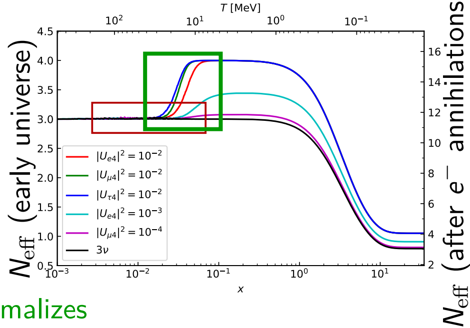
[SG+, JCAP 07 (2019) 014]

dashed: 3ν, solid: $|U_{e4}|^2 = 10^{-2}$, $|U_{\mu 4}|^2 = |U_{\tau 4}|^2 = 0$. $\Delta m^2_{41} = 1.29$ eV^2 always









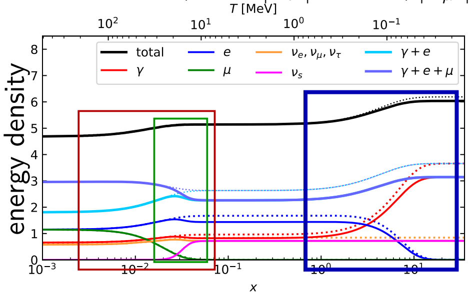
muons annihilate

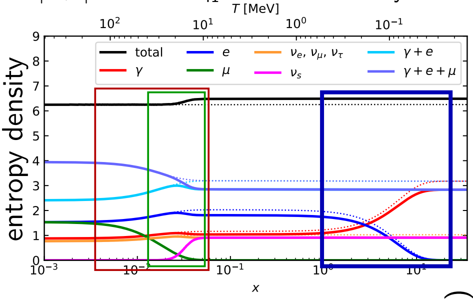
ν_s thermalizes

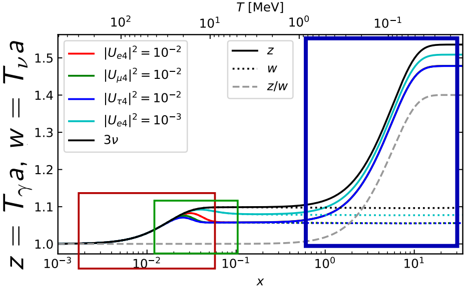
Energy, entropy, temperatures, $N_{\rm eff}$

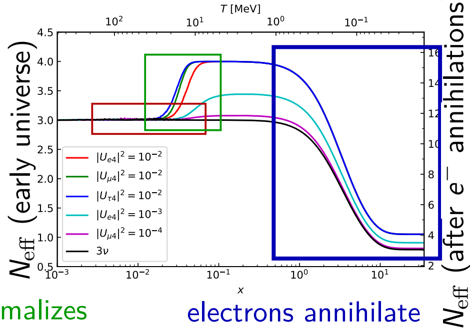
[SG+, JCAP 07 (2019) 014]

dashed: 3ν, solid: $|U_{e4}|^2 = 10^{-2}$, $|U_{\mu4}|^2 = |U_{\tau4}|^2 = 0$. $\Delta m^2_{41} = 1.29$ eV^2 always







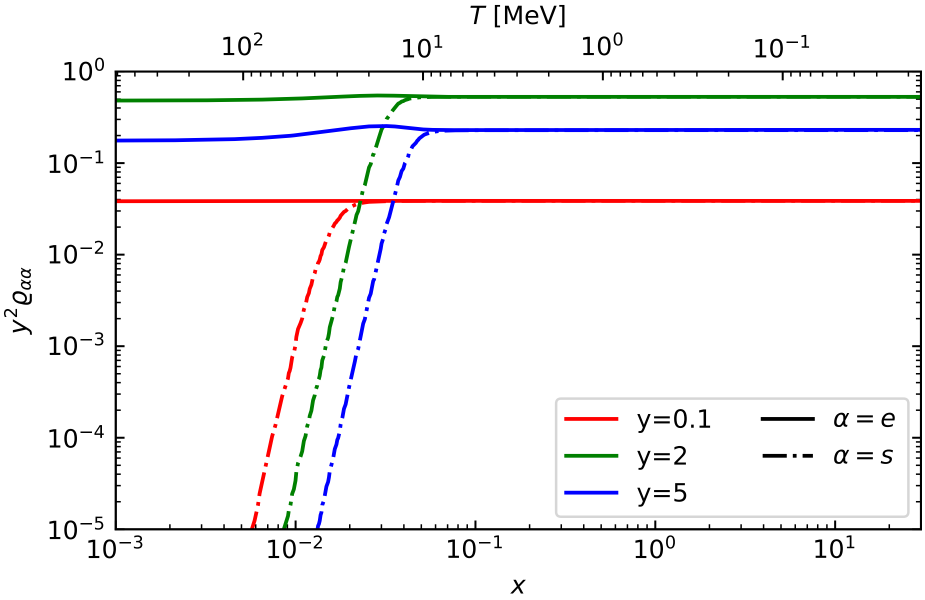


muons annihilate ν_s thermalizes electrons annihilate

Momentum distributions

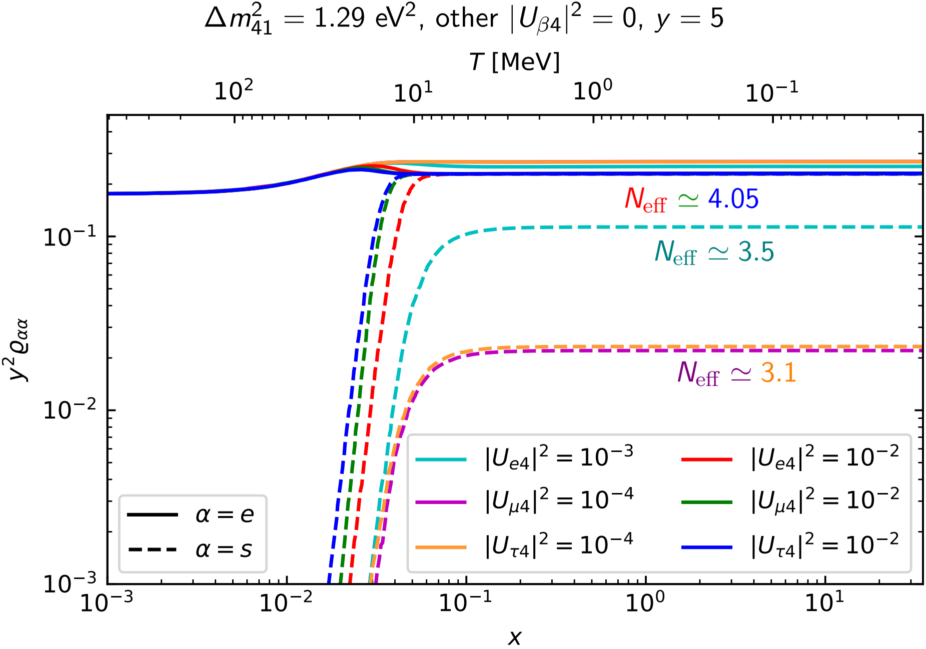
[SG+, JCAP 07 (2019) 014]

$\Delta m^2_{41} = 1.29$ eV^2, $|U_{e4}|^2 = 10^{-2}$, $|U_{\mu4}|^2 = |U_{\tau4}|^2 = 0$, $N_{\rm eff} \simeq 4.05$



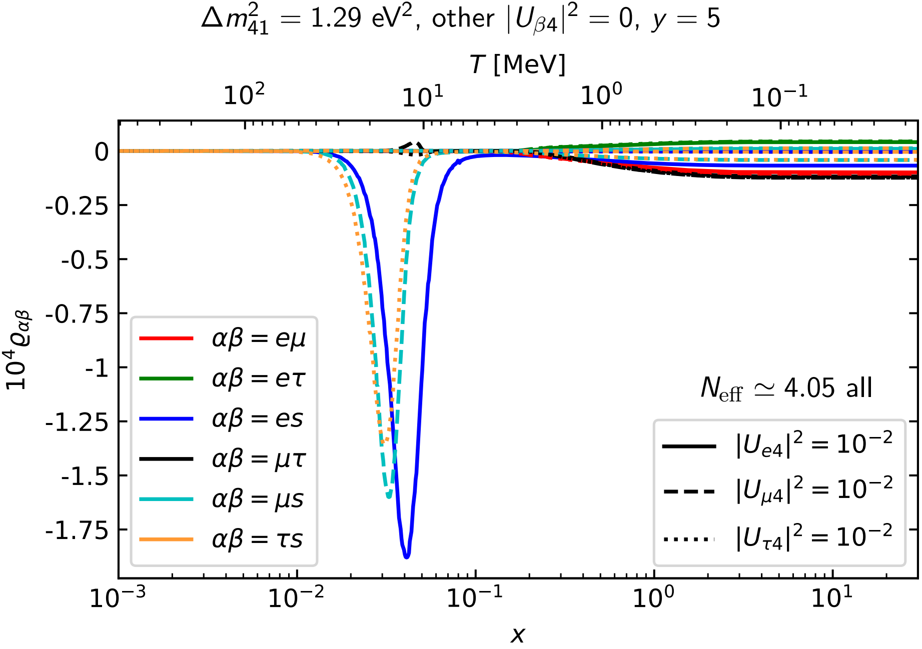
Momentum distributions

[SG+, JCAP 07 (2019) 014]



Momentum distributions

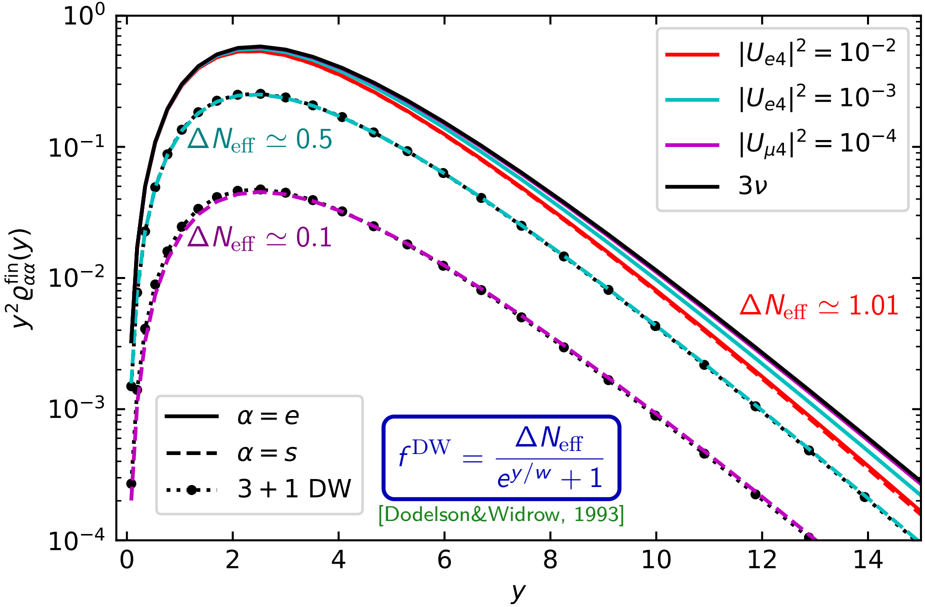
[SG+, JCAP 07 (2019) 014]



Momentum distributions

[SG+, JCAP 07 (2019) 014]

$\Delta m^2_{41} = 1.29$ eV^2, other $|U_{\beta 4}|^2 = 0$, $\Delta N_{\rm eff} = N_{\rm eff} - N_{\rm eff}^{\rm active}$



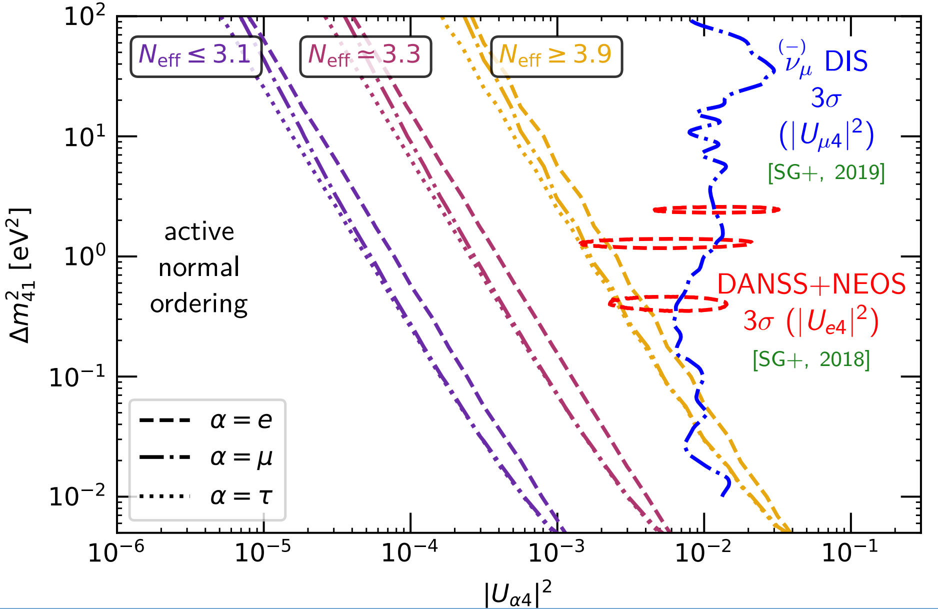
$$f^{\rm DW} = \frac{\Delta N_{\rm eff}}{e^{y/w}+1}$$

[Dodelson&Widrow, 1993]

$N_{\rm eff}$ and the new mixing parameters

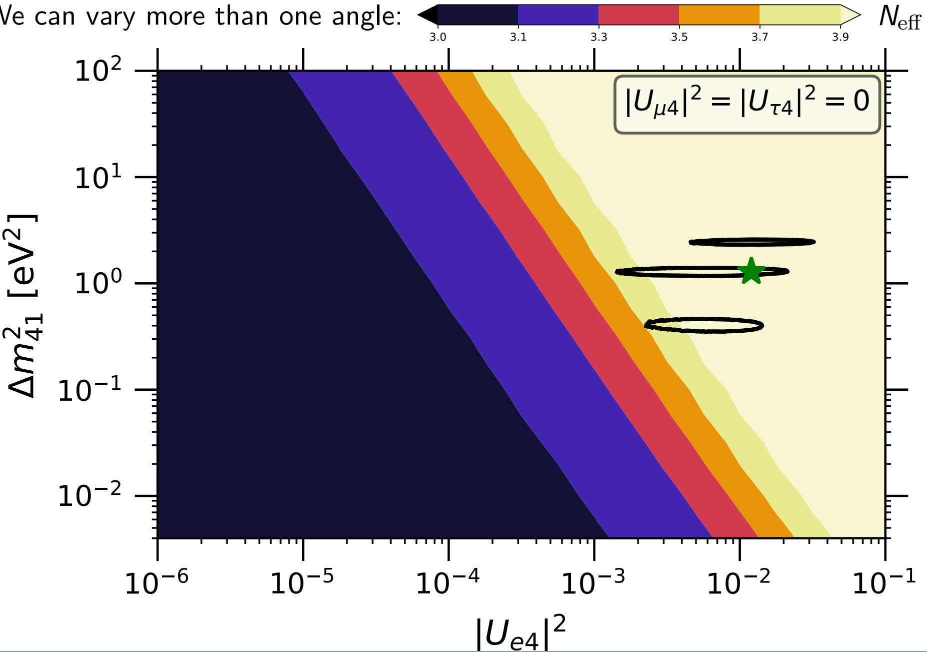
[SG+, JCAP 07 (2019) 014]

Only vary one angle and fix two to zero: do they have the same effect?



$N_{\rm eff}$ and the new mixing parameters

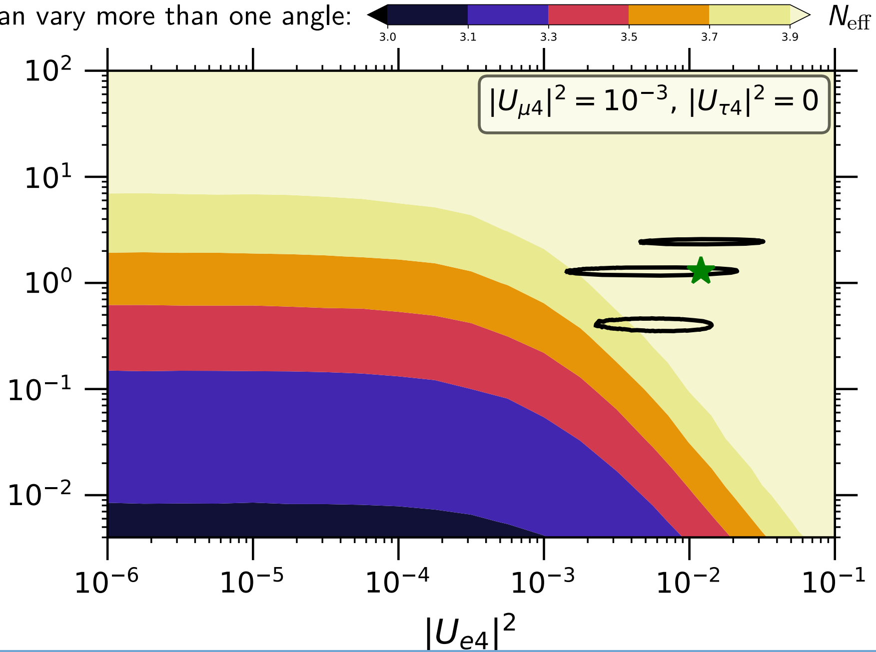
We can vary more than one angle:



$N_{\rm eff}$ and the new mixing parameters

[SG+, JCAP 07 (2019) 014]

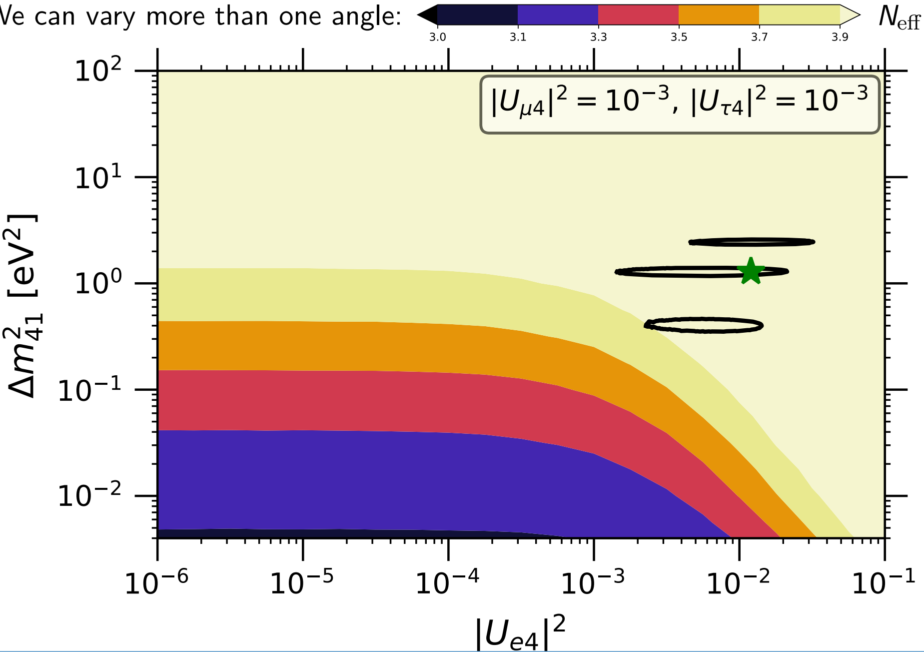
We can vary more than one angle:



$N_{\rm eff}$ and the new mixing parameters

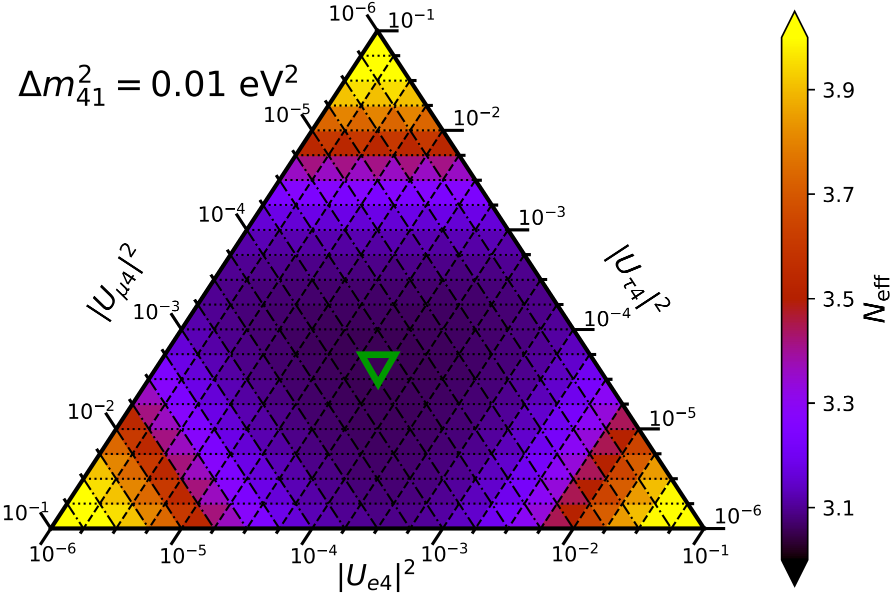
[SG+, JCAP 07 (2019) 014]

We can vary more than one angle:



$N_{\rm eff}$ and the new mixing parameters

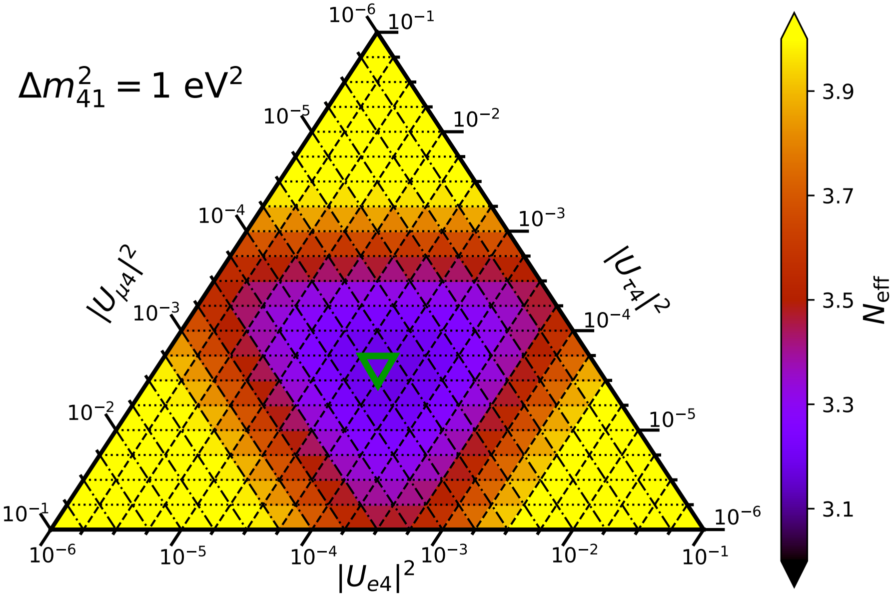
Sort of ternary plot (sum of $|U_{\alpha 4}|^2$ does not add up to 1!):



$N_{\rm eff}$ and the new mixing parameters

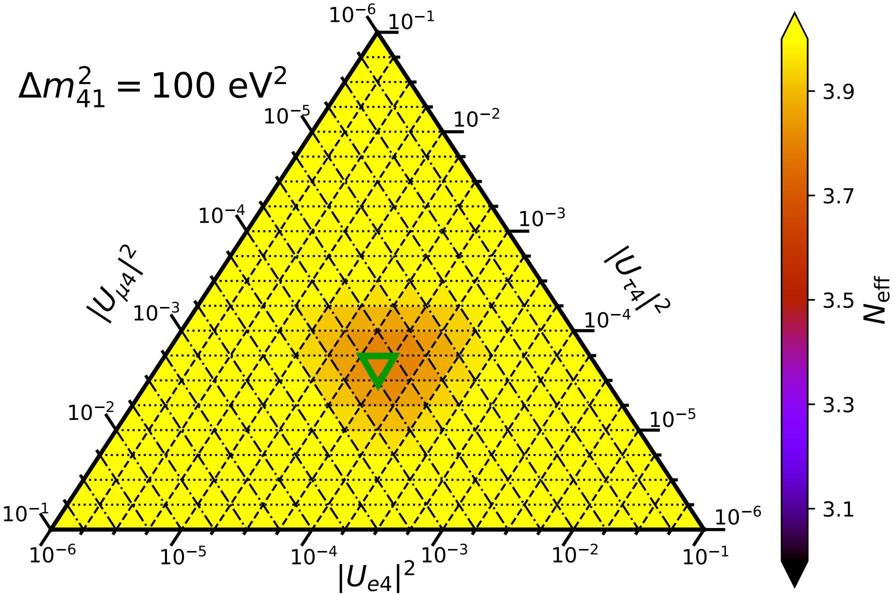
[SG+, JCAP 07 (2019) 014]

Sort of ternary plot (sum of $|U_{\alpha 4}|^2$ does not add up to 1!):



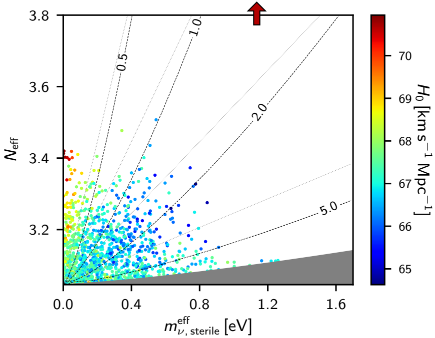
$N_{\rm eff}$ and the new mixing parameters

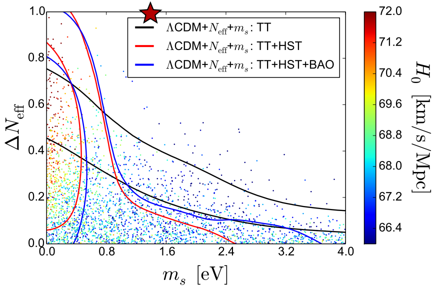
Sort of ternary plot (sum of $|U_{\alpha 4}|^2$ does not add up to 1!):



LSν constraints from cosmology

CMB+local: [Planck Collaboration, 2018]   [Archidiacono et al., JCAP 08 (2016) 067]





$$\begin{cases} N_{\rm eff} < 3.29 \\ m_s^{\rm eff} < 0.65 \text{ eV} \end{cases}$$ (Planck18+BAO) [$m_s < 10$ eV]

dataset	free $\Delta N_{\rm eff}$ [$m_s < 10$ eV]	$\Delta N_{\rm eff} = 1$
(TT)	$N_{\rm eff} < 3.5$	$m_s < 0.66$ eV
(+H_0)	$N_{\rm eff} < 3.9$	$m_s < 0.55$ eV
(+BAO)	$N_{\rm eff} < 3.8$	$m_s < 0.53$ eV

BBN constraints: $N_{\rm eff} = 2.90 \pm 0.22$ (BBN+Y_p) [Peimbert et al., 2016]

Summary: $\Delta N_{\rm eff} = 1$ from LSν incompatible with CMB and BBN!

Planck18=Planck 2018 TT,TE,EE + lowE + lensing   All the constraints are at 2σ CL

Conclusions

1. Unclear model-independent results from $\overset{(-)}{\nu}_e$ DIS,
plus discrepancy with Gallium anomaly and RAA

2. nothing seen in $\overset{(-)}{\nu}_\mu$ DIS
strong upper bounds on $|U_{\mu 4}|^2$,
but also first constraints on $|U_{\tau 4}|^2$

3. strong APP-DIS tension
What are LSND and MiniBooNE observing?
Systematics or LSν or new physics?

Conclusions

1. Unclear model-independent results from $\overset{(-)}{\nu}_e$ DIS, plus discrepancy with Gallium anomaly and RAA

2. nothing seen in $\overset{(-)}{\nu}_\mu$ DIS
strong upper bounds on $|U_{\mu 4}|^2$,
but also first constraints on $|U_{\tau 4}|^2$

3. strong APP-DIS tension
What are LSND and MiniBooNE observing?
Systematics or LSν or new physics?

+1 oscillations in the early universe $\rightarrow N_{\rm eff} \simeq 4.05$
Planck constrains $N_{\rm eff} \lesssim 3.3$!
new physics needed? where???

Conclusions

1. Unclear model-independent results from $\overset{(-)}{\nu}_e$ DIS, plus discrepancy with Gallium anomaly and RAA

2. nothing seen in $\overset{(-)}{\nu}_\mu$ DIS
strong upper bounds on $|U_{\mu 4}|^2$,
but also first constraints on $|U_{\tau 4}|^2$

3. strong APP-DIS tension
What are LSND and MiniBooNE observing?
Systematics or LSν or new physics?

+1 oscillations in the early universe $\rightarrow$ $N_{\rm eff} \simeq 4.05$
Planck constrains $N_{\rm eff} \lesssim 3.3$!
new physics needed? where???

Thank you for the attention!