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Light sterile neutrinos in the early universe

XI CPAN Meeting, Oviedo (ES), 21-23/10/2019

1 Light sterile neutrino

2 Light sterile neutrino and cosmology

3 A new interaction to solve the thermalization problem?

4 Conclusions



Three Neutrino Oscillations

Analogous to CKM mixing for quarks:

[Pontecorvo, 1968] [Maki, Nakagawa, Sakata, 1962]

$$u_{\alpha} = \sum_{k=1}^{3} U_{\alpha k} \nu_k \quad (\alpha = e, \mu, \tau)$$

 u_{α} flavour eigenstates, $U_{\alpha k}$ PMNS mixing matrix, ν_k mass eigenstates.

Current knowledge of the 3 active ν mixing: [de Salas et al. (2018)]

 $\Delta m_{ji}^2 = m_j^2 - m_i^2$, θ_{ij} mixing angles NO: Normal Ordering, $m_1 < m_2 < m_3$ IO: Inverted Ordering, $m_3 < m_1 < m_2$



Three Neutrino Oscillations

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Short Baseline (SBL) anomaly

[SG+, JPG 43 (2016) 033001]

Problem: anomalies in SBL experiments

errors in flux calculations? deviations from 3-v description?

A short review:

- LSND search for $\bar{\nu}_{\mu} \rightarrow \bar{\nu}_{e}$, with $L/E = 0.4 \div 1.5$ m/MeV. Observed a 3.8 σ excess of $\bar{\nu}_e$ events [Aguilar et al., 2001]
- Reactor re-evaluation of the expected anti-neutrino flux \Rightarrow disappearance of $\bar{\nu}_e$ events compared to predictions ($\sim 3\sigma$) with L < 100 m [Mention et al, 2011], [Azabajan et al, 2012]
- Gallium calibration of GALLEX and SAGE Gallium solar neutrino experiments give a 2.7 σ anomaly (disappearance of ν_e) [Giunti, Laveder, 2011]

MiniBooNE

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See next
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Possible explanation: Additional squared mass difference $\Delta m_{SBI}^2 \simeq 1 \text{ eV}^2$



[DANSS, PLB 787 (2018) 56]



first *model independent* indications in favor of SBL oscillations

DANSS alone gives a $\Delta \chi^2 \simeq 13$ in favor of a light sterile neutrino!

[MiniBooNE, PRL 121 (2018) 221801]



[MiniBooNE, PRL 121 (2018) 221801]



[MiniBooNE, PRL 121 (2018) 221801]



[MINOS+, PRL 122 (2019) 091803]



with NEOS&DANSS



3+1 Neutrino Model

new $\Delta m_{SBL}^2 \Rightarrow 4$ neutrinos! ν_4 with $m_4 \simeq 1$ eV, no weak interactions light sterile neutrino (LS ν) 3 (active) + 1 (sterile) mixing: $u_{\alpha} = \sum U_{\alpha k} \nu_k \quad (\alpha = e, \mu, \tau, s)$ k=1 $\nu_{\rm s}$ is mainly $\nu_{\rm 4}$: $m_s \simeq m_4 \simeq \sqrt{\Delta m_{41}^2} \simeq \sqrt{\Delta m_{\rm SBL}^2}$ assuming $m_4 \gg m_i$ (i = 1, 2, 3)

[SG+, work in progress]



3+1 Neutrino Model

new
$$\Delta m_{\text{SBL}}^2 \Rightarrow 4$$
 neutrinos!
 \downarrow
 ν_4 with $m_4 \simeq 1$ eV,
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light sterile neutrino (LS ν)
3 (active) + 1 (sterile) mixing:
 $\nu_{\alpha} = \sum_{k=1}^{3+1} U_{\alpha k} \nu_k \quad (\alpha = e, \mu, \tau, s)$
 ν_s is mainly ν_4 :
 $m_s \simeq m_4 \simeq \sqrt{\Delta m_{41}^2} \simeq \sqrt{\Delta m_{\text{SBL}}^2}$
assuming $m_4 \gg m_i$ ($i = 1, 2, 3$)

[SG+, work in progress]



can ν_4 thermalize in the early Universe through oscillations?

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 $m_{\rm P1}$ Planck mass – ρ_T total energy density – $m_{W,Z}$ mass of the W, Z bosons – $G_{\rm F}$ Fermi constant – [., .] commutator

[SG+, JCAP 07 (2019) 014] ν oscillations in the early universe comoving coordinates: a = 1/T $x \equiv m_e a$ $y \equiv p a$ $z \equiv T_{\gamma} a$ $w \equiv T_{\nu} a$ density matrix: $\varrho(x, y) = \begin{pmatrix}
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\varrho_{\mu e} & \varrho_{\mu\mu} \equiv f_{\nu_{\mu}} & \varrho_{\mu\tau} & \varrho_{\mu s} \\
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\end{pmatrix}$ $\frac{\mathrm{d}\varrho(y,x)}{\mathrm{d}x} = \sqrt{\frac{3m_{\mathrm{Pl}}^2}{8\pi\rho_{\mathrm{T}}}} \left\{ -i\frac{x^2}{m_e^3} \left[\frac{\mathbb{M}_{\mathrm{F}}}{2y} - \frac{8\sqrt{2}G_{\mathrm{F}}ym_e^6}{3x^6} \left(\frac{\mathbb{E}_{\ell}}{m_{\mathrm{ev}}^2} + \frac{\mathbb{E}_{\nu}}{m_{\mathrm{e}}^2} \right), \varrho \right] + \frac{m_e^3G_{\mathrm{F}}^2}{(2\pi)^3x^4v^2} \mathcal{I}(\varrho) \right\} \left| \right|$ $m_{\rm P1}$ Planck mass – ρ_T total energy density – $m_{W,Z}$ mass of the W,Z bosons – $G_{\rm F}$ Fermi constant – [., .] commutator $M_{\rm F} = U M U^{\dagger}$ $U = R^{34} R^{24} R^{14} R^{23} R^{13} R^{12} \quad \text{e.g. } R^{14} = \begin{pmatrix} \cos \theta_{14} & 0 & 0 & \sin \theta_{14} \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 1 & 0 \\ -\sin \theta_{14} & 0 & 0 & \cos \theta_{14} \end{pmatrix}$ $|U|^2 = \begin{pmatrix} \cdots \cdots \cdots \cos^2 \theta_{14} \sin^2 \theta_{14} \\ \cdots \cdots \cdots \cos^2 \theta_{14} \cos^2 \theta_{24} \sin^2 \theta_{34} \\ \cdots \cdots \cos^2 \theta_{14} \cos^2 \theta_{24} \cos^2 \theta_{34} \end{pmatrix}$ $\mathbb{M} = \operatorname{diag}(m_1^2, \ldots, m_N^2)$



take into account matter effects in oscillations

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ho_{e},
ho_{\mu}, 0, 0)$ $\mathbb{E}_{
u} = S_{a}\left(\int dyy^{3}arrho\right)S_{a}$ $\mathcal{I}(\rho)$ collision integrals

take into account neutrino-electron scattering and pair annihilation 2D integrals over the momentum, take most of the computation time

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ight) S_a$ $\mathbb{M}_{\mathbf{F}} = U\mathbb{M}U^{\dagger}$ $\mathcal{I}(\rho)$ collision integrals $\frac{\mathrm{d}z}{\mathrm{d}x} = \frac{\sum_{\ell=e,\mu} \left[\frac{r_{\ell}^2}{r} J(r_{\ell})\right] + G_1(r) - \frac{1}{2\pi^2 z^3} \int_0^{\infty} dy \, y^3 \sum_{\alpha=e}^{s} \frac{\mathrm{d}\varrho_{\alpha\alpha}}{\mathrm{d}x}}{\sum_{\alpha=e} \left[\frac{r_{\ell}^2}{r} J(r_{\ell}) + Y(r_{\ell})\right] + G_2(r) + \frac{2\pi^2}{15}}$ from continuity equation $\dot{\rho} = -3H(\rho + P)$ $\ell = e.\mu$ r = x/z, $r_{\ell} = m_{\ell}/m_e r$ J(r), Y(r) from non-relativistic transition of e^{\pm} , μ^{\pm} $G_1(r)$ and $G_2(r)$ from electromagnetic corrections

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neutrino temperature w: same equation as z, but electrons always relativistic

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"Light sterile neutrinos in the early universe" Ovied

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$N_{\rm eff}$ and the new mixing parameters

Only vary one angle and fix two to zero: do they have the same effect?







 $N_{\rm eff}$ and the new mixing parameters



 $I_{\rm N_{eff}}$ and the new mixing parameters







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BBN constraints: $N_{
m eff} = 2.90 \pm 0.22 \; ({
m BBN} + Y_p)$ [Peimbert et al., 2016]

Summary: $\Delta N_{\text{eff}} = 1$ from LS ν incompatible with CMB and BBN!

Planck18=Planck 2018 TT,TE,EE + lowE + lensing ~~ All the constraints are at $2\sigma~$ CL ~

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Incomplete Thermalization

Active-sterile oscillations in the early Universe: mixing parameters from SBL data $\implies \Delta N_{\rm eff} \simeq 1$ [SG+, 2019]

Many probes constrain $\Delta \textit{N}_{\rm eff} < 1.$ Do we need

- a mechanism to suppress oscillations and full thermalization of ν_s ?
- to compensate $\Delta N_{
 m eff} = 1$ with additional mechanisms in Cosmology?
- Some ideas (an incomplete list!):
 - Iarge lepton asymmetry [Foot et al., 1995; Mirizzi et al., 2012; many more]
 - new neutrino interactions [Bento et al., 2001; Dasgupta et al., 2014; Hannestad et al., 2014; Saviano et al., 2014; Archidiacono et al. 2016; many more]
 - entropy production after neutrino decoupling [Ho et al., 2013]
 - very low reheating temperature [Gelmini et al., 2004; Smirnov et al., 2006]
 - time varying dark energy components [Giusarma et al., 2012]
 - Iarger expansion rate at the time of ν_s production [Rehagen et al., 2014]
 - freedom in the Primordial Power Spectrum (PPS) of scalar perturbations from inflation compensate damping due to $N_{\rm eff} \neq 3.046$ [SG et al., 2015]

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Adding a new interaction



Constraints on the pseudoscalar interaction?



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 Problems with ΔN_{eff} = 1? solved (incomplete thermalization due to suppression of active-sterile oscillations in primordial plasma);

- mass bounds avoided
 - \Rightarrow large m_s allowed and (mild) preference for $m_s \simeq 4$ eV;
- high values of H_0 predicted by cosmology
 - \Rightarrow more compatible with local measurements.

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APP-DIS tension! Systematics or $LS\nu$ or new physics?

oscillations in the early universe $\Rightarrow N_{\rm eff} \simeq 4.05$ Planck constrains $N_{\rm eff} \lesssim 3.3!$





2





[Danilov@EPS-HEP, 2019]

old data



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