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Horizon 2020
European Union funding
for Research & Innovation

Light sterile neutrinos: the current picture

Conference on Flavour Physics and CP violation (FPCP), 8–12/06/2020

<https://www.particlesforjustice.org/>

- 1 *Neutrino Oscillations - Some theory*
- 2 *Electron (anti)neutrino disappearance*
- 3 *Muon (anti)neutrino disappearance*
- 4 *Electron (anti)neutrino appearance*
- 5 *Global fit*
- 6 *Recent updates*
- 7 *Light sterile neutrino and cosmology*
- 8 *Conclusions*

Three Neutrino Oscillations

$$\nu_\alpha = \sum_{k=1}^3 U_{\alpha k} \nu_k \quad (\alpha = e, \mu, \tau)$$

$U_{\alpha k}$ described by 3 mixing angles θ_{12} , θ_{13} , θ_{23} and one CP phase δ

Current knowledge of the 3 active ν mixing: [in preparation]

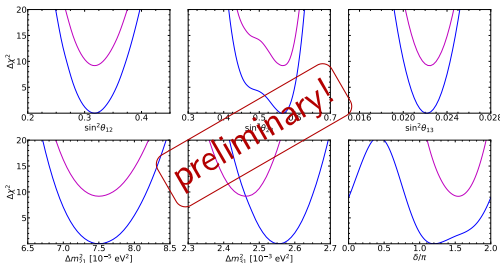
NO/NH: Normal Ordering/Hierarchy, $m_1 < m_2 < m_3$

IO/IH: Inverted O/H, $m_3 < m_1 < m_2$

$$\begin{aligned}\Delta m_{21}^2 &= (7.50^{+0.22}_{-0.20}) \cdot 10^{-5} \text{ eV}^2 \\ |\Delta m_{31}^2| &= (2.56^{+0.03}_{-0.04}) \cdot 10^{-3} \text{ eV}^2 \text{ (NO)} \\ &= (2.46 \pm 0.03) \cdot 10^{-3} \text{ eV}^2 \text{ (IO)}\end{aligned}$$

$$\begin{aligned}10 \sin^2(\theta_{12}) &= 3.18 \pm 0.16 \\ 10^2 \sin^2(\theta_{13}) &= 2.225^{+0.055}_{-0.078} \text{ (NO)} \\ &= 2.250^{+0.056}_{-0.076} \text{ (IO)} \\ 10 \sin^2(\theta_{23}) &= 5.66^{+0.16}_{-0.22} \text{ (NO)} \\ &= 5.66^{+0.18}_{-0.23} \text{ (IO)}\end{aligned}$$

$$\begin{aligned}\delta/\pi &= 1.20^{+0.23}_{-0.14} \text{ (NO)} \\ &= 1.54 \pm 0.13 \text{ (IO)}\end{aligned}$$



see next talk by I. Esteban

see also: <http://globalfit.astroparticles.es>

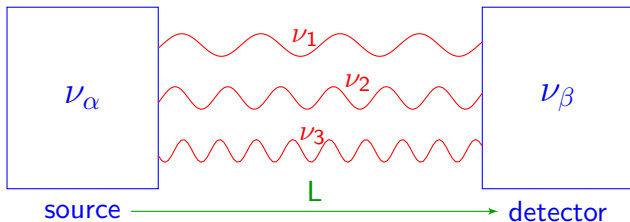
Two neutrino bases

flavor neutrinos ν_α

$$|\nu_\alpha\rangle = \sum_k U_{\alpha k} |\nu_k\rangle$$

massive neutrinos ν_k

$$|\nu(t=0)\rangle = |\nu_\alpha\rangle = U_{\alpha 1} |\nu_1\rangle + U_{\alpha 2} |\nu_2\rangle + U_{\alpha 3} |\nu_3\rangle$$



$$|\nu(t > 0)\rangle = |\nu_\beta\rangle = U_{\alpha 1} e^{-iE_1 t} |\nu_1\rangle + U_{\alpha 2} e^{-iE_2 t} |\nu_2\rangle + U_{\alpha 3} e^{-iE_3 t} |\nu_3\rangle \neq |\nu_\alpha\rangle$$

$$E_k^2 = p^2 + m_k^2 \longleftarrow \text{define} \longrightarrow t = L$$

$$P_{\nu_\alpha \rightarrow \nu_\beta}(L) = |\langle \nu_\alpha | \nu(L) \rangle|^2 = \sum_{k,j} U_{\beta k} U_{\alpha k}^* U_{\beta j}^* U_{\alpha j} \exp\left(-i \frac{\Delta m_{kj}^2 L}{2E}\right)$$

$$\Delta m_{ij}^2 = m_i^2 - m_j^2$$

A large family

In principle, previous discussion is valid for N neutrinos

only constraint: there are exactly three flavor neutrinos in the SM

[LEP, Phys. Rept. 427 (2006) 257,
arXiv:hep-ex/0509008]

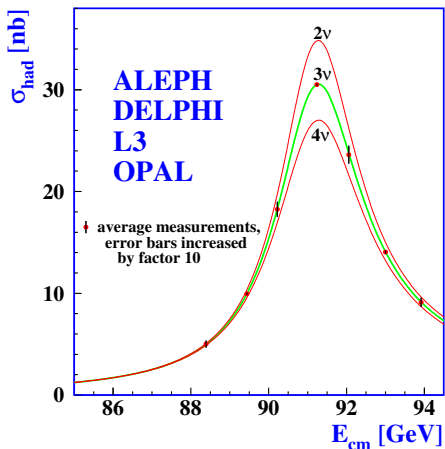
$$N_{\nu}^{(Z)} = 2.9840 \pm 0.0082$$

through the measurement
of the Z resonance

$$e^+e^- \rightarrow Z \rightarrow \sum_{a=e,\mu,\tau} \nu_a \bar{\nu}_a$$

neutrinos $\alpha > 3$ must be sterile

sterile neutrino = SM singlet: no couplings with other SM particles



A large family

In principle, previous discussion is valid for N neutrinos

$N \times N$ mixing matrix, N flavor neutrinos, N massive neutrinos

$$\begin{pmatrix} |\nu_e\rangle \\ |\nu_\mu\rangle \\ |\nu_\tau\rangle \\ |\nu_{s1}\rangle \\ \dots \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & U_{e4} & \vdots \\ U_{\mu1} & U_{\mu2} & U_{\mu3} & U_{\mu4} & \\ U_{\tau1} & U_{\tau2} & U_{\tau3} & U_{\tau4} & \\ U_{s11} & U_{s12} & U_{s13} & U_{s14} & \\ \dots & & & & \ddots \end{pmatrix} \begin{pmatrix} |\nu_1\rangle \\ |\nu_2\rangle \\ |\nu_3\rangle \\ |\nu_4\rangle \\ \dots \end{pmatrix}$$

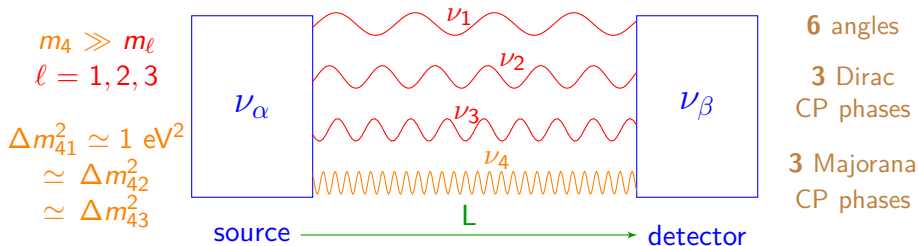
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Our case will be 3 (active)+1 (sterile), a perturbation of 3 neutrinos case



Short BaseLine (SBL)

$$P_{\nu_\alpha \rightarrow \nu_\beta}(L) = |\langle \nu_\alpha | \nu(L) \rangle|^2 = \sum_{k,j} U_{\beta k} U_{\alpha k}^* U_{\beta j}^* U_{\alpha j} \exp\left(-i \frac{\Delta m_{kj}^2 L}{2E}\right)$$

If $m_4 \gg m_\ell$, faster oscillations

ν_4 oscillations are averaged in most neutrino oscillation experiments

Effect of 4th neutrino only visible as global normalization

Short BaseLine (SBL) oscillations: $\frac{\Delta m_{41}^2 L}{E} \simeq 1$

At SBL, oscillations due to Δm_{21}^2 and $|\Delta m_{31}^2|$ do not develop

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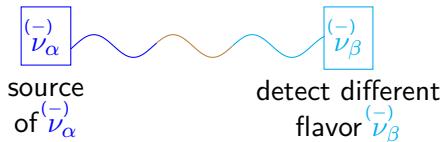
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APPEARANCE ($\alpha \neq \beta$)



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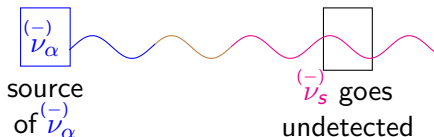
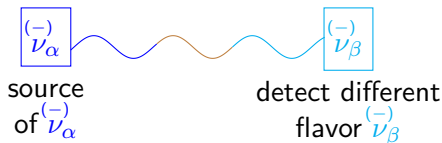
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DISappearance



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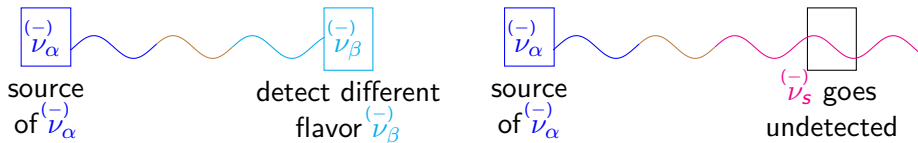
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DISappearance



CP violation cannot be observed in SBL experiments!

New mixings in the 3+1 scenario

4 × 4 mixing matrix:

$$\begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & U_{e4} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} & U_{\mu4} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} & U_{\tau4} \\ U_{s11} & U_{s12} & U_{s13} & U_{s14} \end{pmatrix}$$

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DISappearance

$$P_{\nu_{\alpha}^{(-)} \rightarrow \nu_{\alpha}^{(-)}}^{\text{SBL}} \simeq 1 - \sin^2 2\vartheta_{\alpha\alpha} \sin^2 \left(\frac{\Delta m_{41}^2 L}{4E} \right)$$

$$\sin^2 2\vartheta_{\alpha\alpha} = 4|U_{\alpha 4}|^2(1 - |U_{\alpha 4}|^2)$$

$\nu_e^{(-)} \rightarrow \nu_e^{(-)}$

reactor
gallium

$$|U_{e4}|^2 = \sin^2 \vartheta_{14}$$

$\nu_{\mu}^{(-)} \rightarrow \nu_{\mu}^{(-)}$

accelerator
atmospheric

$$|U_{\mu 4}|^2 = \cos^2 \vartheta_{14} \sin^2 \vartheta_{24}$$

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atmospheric

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APPEARance

$$P_{\nu_{\alpha} \rightarrow \nu_{\beta}}^{SBL(-)} \simeq \sin^2 2\vartheta_{\alpha\beta} \sin^2 \left(\frac{\Delta m_{41}^2 L}{4E} \right)$$

$$\sin^2 2\vartheta_{\alpha\beta} = 4|U_{\alpha 4}|^2 |U_{\beta 4}|^2$$

$$\nu_{\mu}^{(-)} \rightarrow \nu_e^{(-)}$$

LSND
MiniBooNE
KARMEN
OPERA
...

$$\sin^2 2\vartheta_{e\mu} = 4|U_{e4}|^2 |U_{\mu 4}|^2$$

quadratically suppressed!

for small $|U_{e4}|^2$, $|U_{\mu 4}|^2$

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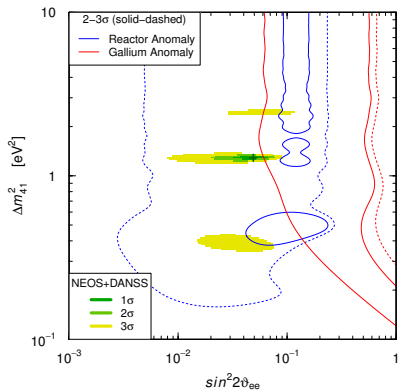
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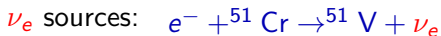


Gallium anomaly

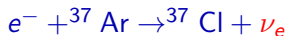
[SAGE, 2006][Laveder, 2007][Giunti&Laveder, 2011]

$L \simeq 1.9 \text{ m}$ $L \simeq 0.6 \text{ m}$

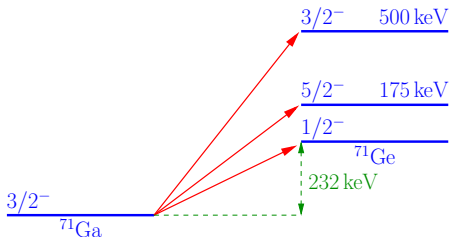
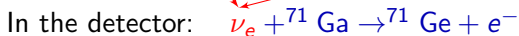
Gallium radioactive source experiments: **GALLEX** and **SAGE**



$E \simeq 0.75 \text{ MeV}$



$E \simeq 0.81 \text{ MeV}$



cross sections of
the transitions from

[Krofcheck et al., PRL 55 (1985) 1051]

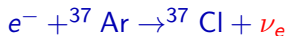
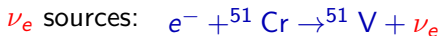
[Frekers et al., PLB 706 (2011) 134]

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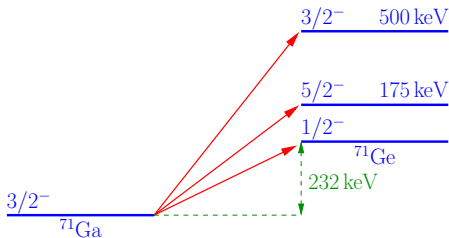
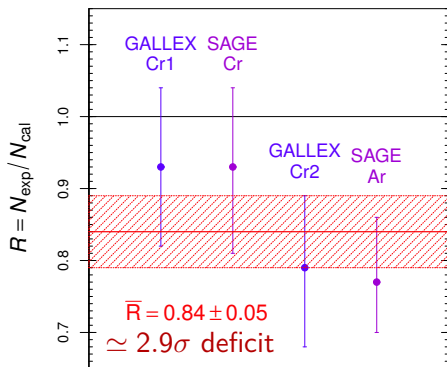


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Test detection of solar ν_e



cross sections of
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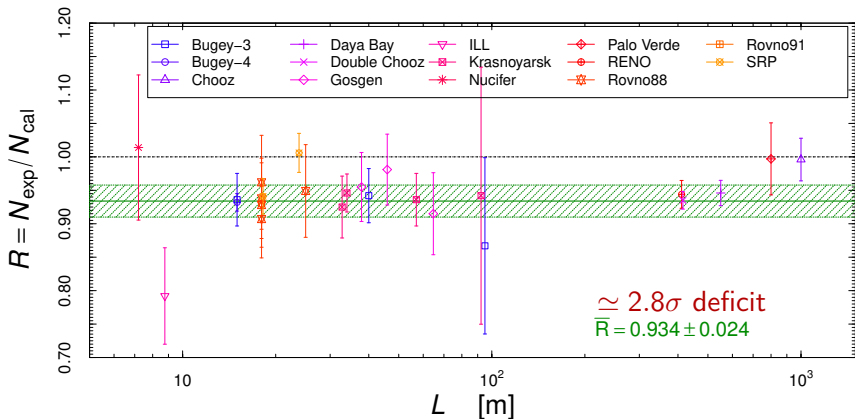
[Krofcheck et al., PRL 55 (1985) 1051]

[Frekers et al., PLB 706 (2011) 134]

2011: new reactor $\bar{\nu}_e$ fluxes by Huber and Mueller+ (HM)

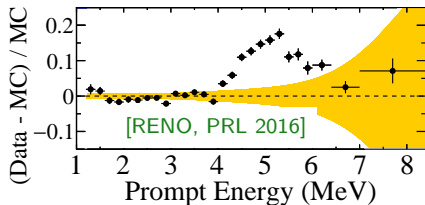
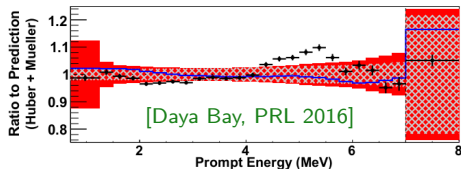
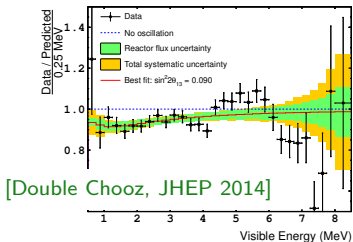
[Huber, PRC 84 (2011) 024617] [Mueller et al., PRC 83 (2011) 054615]

Previous reactor rates evaluated with new fluxes \Rightarrow deficit



Suppression at detector due to active-sterile oscillations?

Can we trust the HM fluxes?



known since 2014:
bump in the spectrum
around 5 MeV!

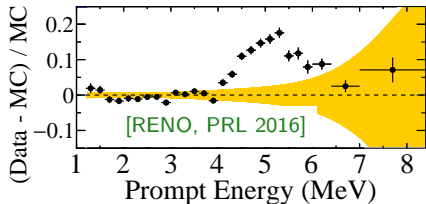
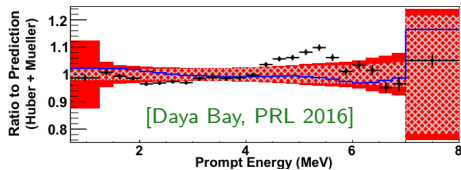
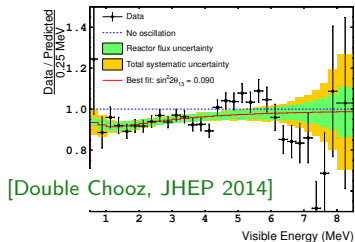
cannot be explained
by SBL oscillations

(averaged at the ob-
served distances)

many attempts of
possible explanations,
how to clarify the issue?

see also talk by
B. Roskovec
on Friday!

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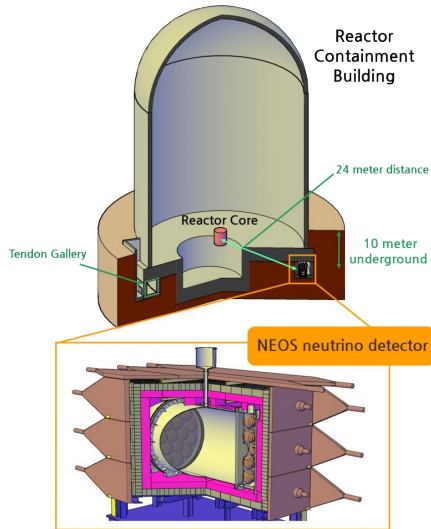
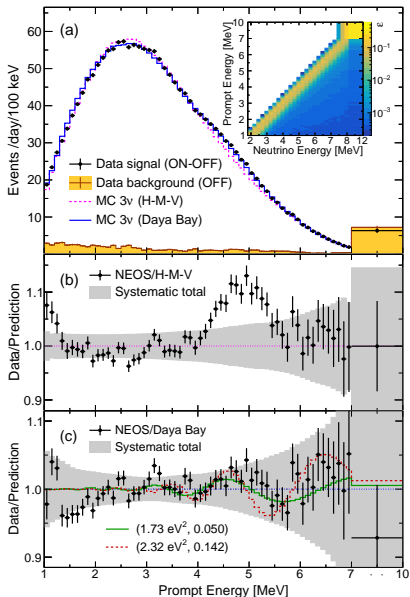
Model independent information!

(i.e. take ratio of spectra
at different distances)

$$\Phi_1 = \Phi_0(E)f(L_1, E) \quad \Phi_2 = \Phi_0(E)f(L_2, E)$$

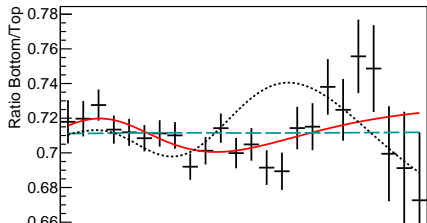
$$\Phi_1/\Phi_2 = f(L_1, E)/f(L_2, E)$$

Single detector experiment

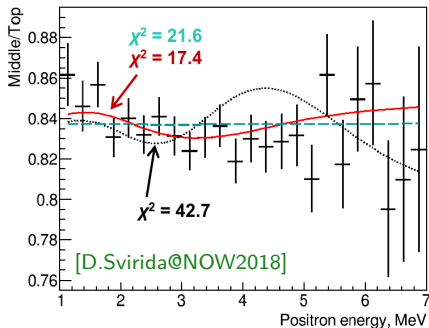


Ratio to DayaBay measurement to be model independent

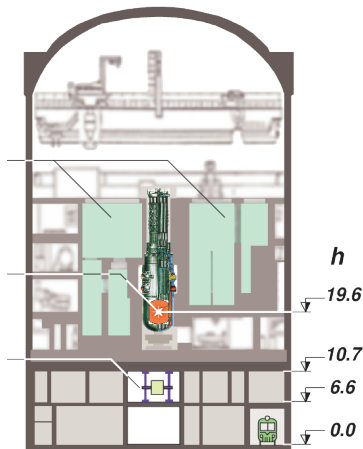
Single movable detector



~ 3 σ preference for 3+1 oscillations

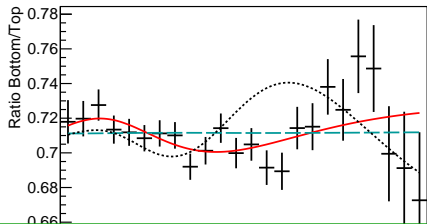


[D.Svirida@NOW2018]

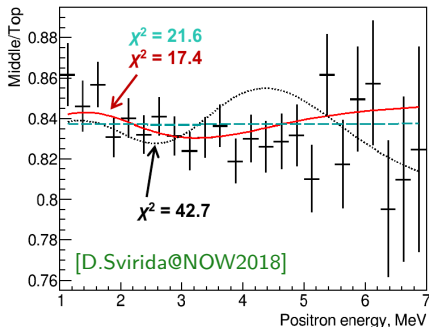


Detector can be at ~ 10.5, ~ 11.5
or ~ 12.5 m from reactor core

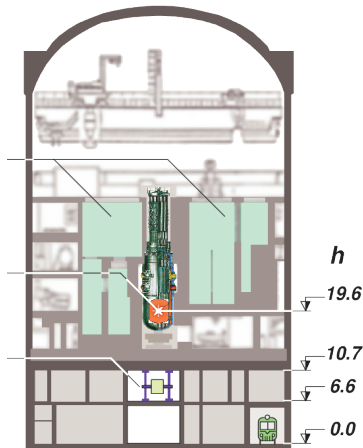
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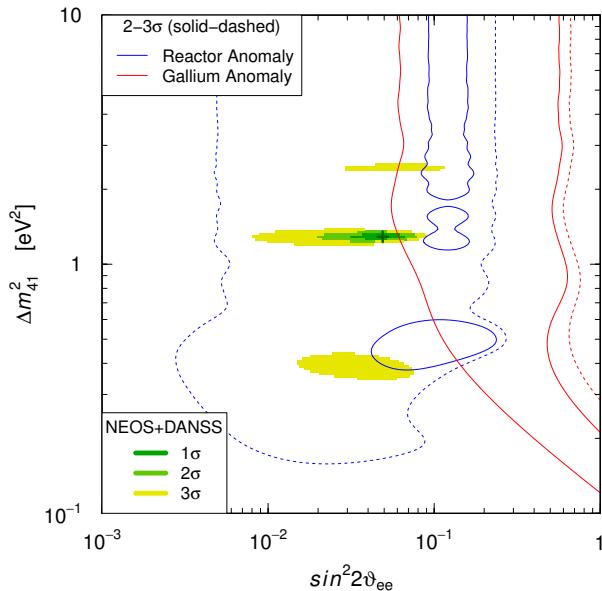
[D.Svirida@NOW2018]



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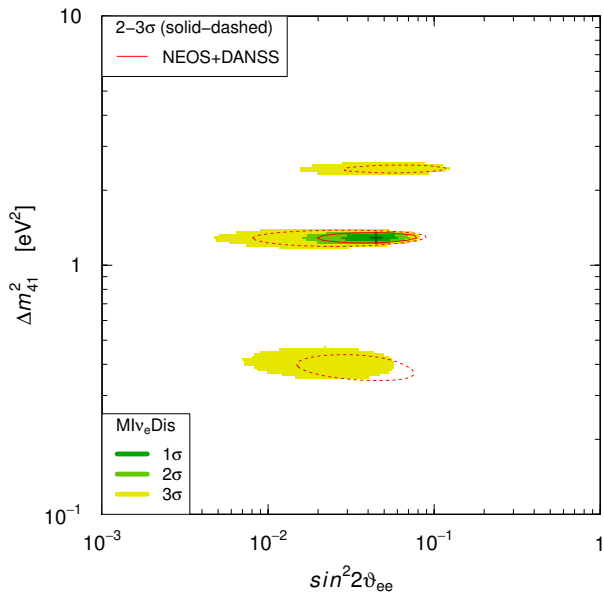
see later for 2019 update!

DANSS + NEOS + RAA + Gallium



DANSS + NEOS
do not agree with
Gallium and RAA

All data:



Fit dominated by
 DANSS + NEOS

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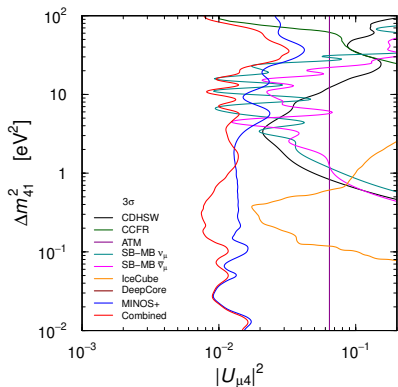
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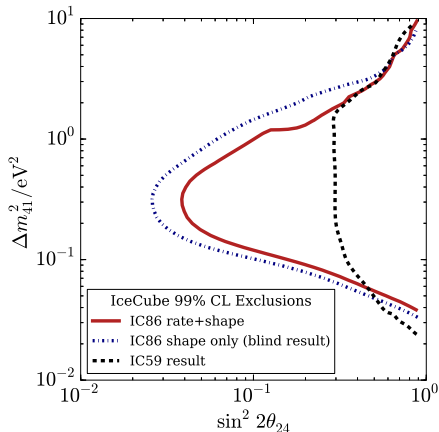
IceCube and DeepCore

IceCube

$\mathcal{O}(10 \text{ km}) \lesssim L \lesssim \mathcal{O}(10^4 \text{ km})$

$\sim 2 \times 10^4$ High energy μ events

$320 \text{ GeV} < E < 20 \text{ TeV}$



[PRL 117 (2016) 071801]

IceCube and DeepCore

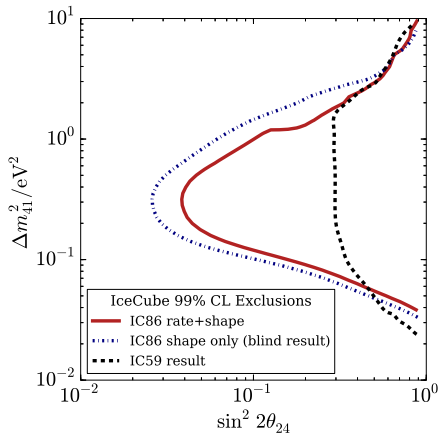
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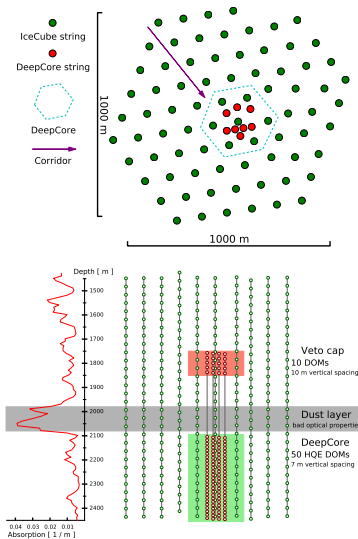
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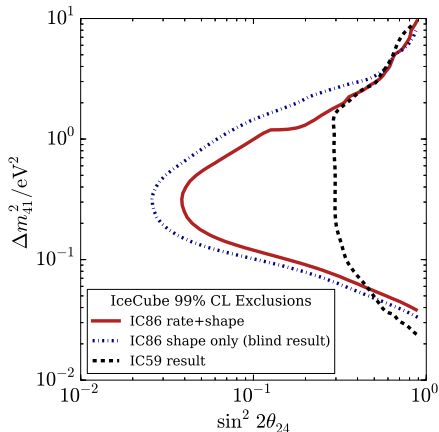
DeepCore

$\sim 2 \times 10^4$ High energy μ events

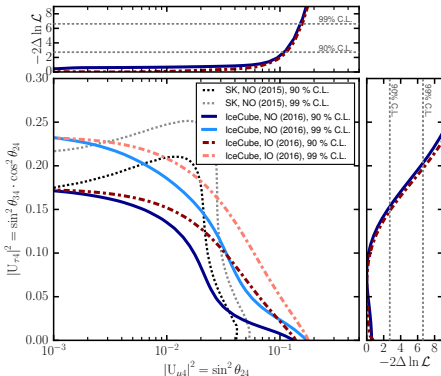
$320 \text{ GeV} < E < 20 \text{ TeV}$

$\sim 5 \times 10^3$ tracklike events

$6 \text{ GeV} \lesssim E \lesssim 60 \text{ GeV}$



[PRL 117 (2016) 071801]

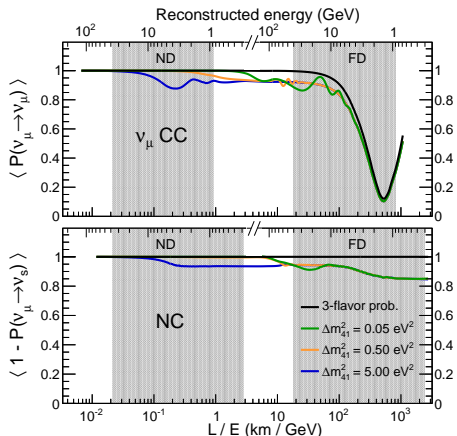


[PRD 95 (2017) 112002]

Both also constrain $|U_{\tau 4}|^2$

Near (ND, $\simeq 500$ m) and
far (FD, $\simeq 800$ km) detector

$1 \text{ GeV} \lesssim E \lesssim 40 \text{ GeV}$,
peak at 3 GeV



[PRL 117 (2016) 151803]:

far-to-near ratio

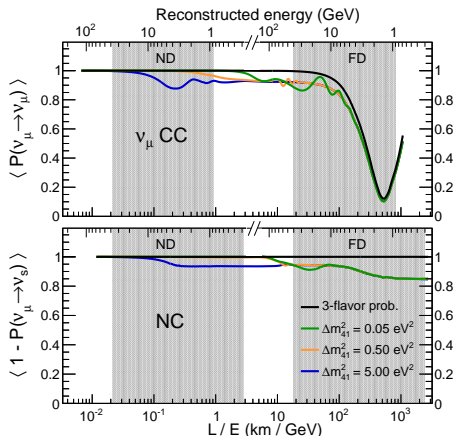
[PRL 122 (2019) 091803]:

full two-detectors fit

MINOS & MINOS+

Near (ND, $\simeq 500$ m) and
far (FD, $\simeq 800$ km) detector

$1 \text{ GeV} \lesssim E \lesssim 40 \text{ GeV}$,
peak at 3 GeV

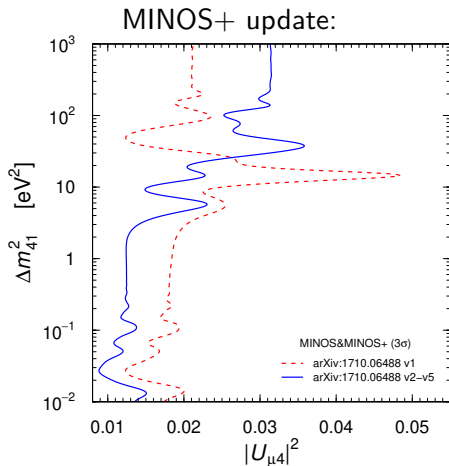


[PRL 117 (2016) 151803]:

far-to-near ratio

[PRL 122 (2019) 091803]:

full two-detectors fit

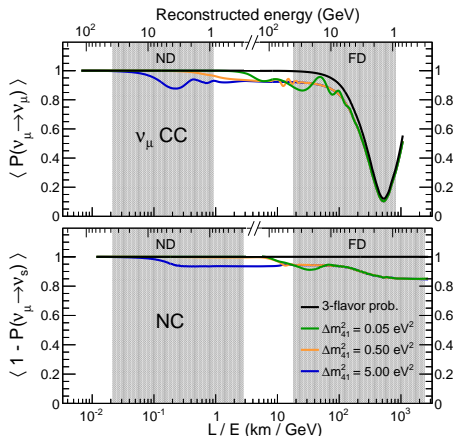


[SG+, in preparation]

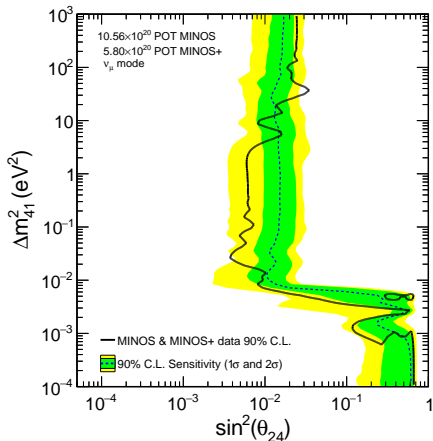
MINOS & MINOS+

Near (ND, ≈ 500 m) and
far (FD, ≈ 800 km) detector

$1 \text{ GeV} \lesssim E \lesssim 40 \text{ GeV}$,
peak at 3 GeV



Sensitivity and exclusion limit:



[PRL 117 (2016) 151803]:

far-to-near ratio

[PRL 122 (2019) 091803]:

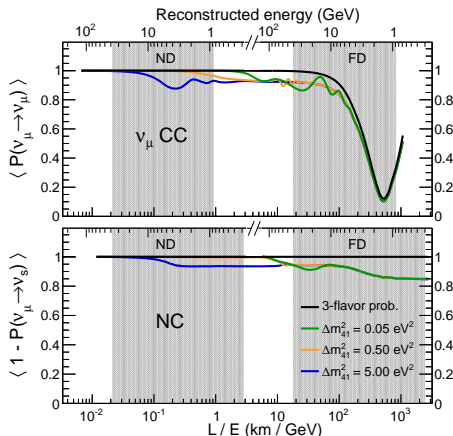
full two-detectors fit

[PRL 122 (2019) 091803]

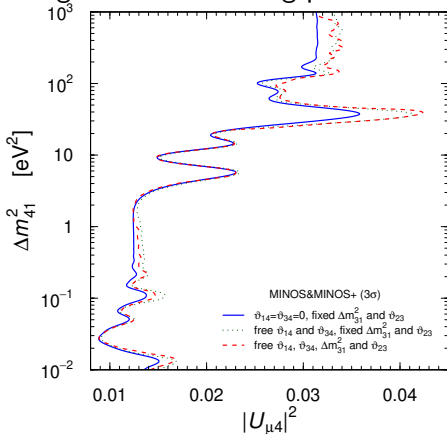
MINOS & MINOS+

Near (ND, $\simeq 500$ m) and
far (FD, $\simeq 800$ km) detector

$1 \text{ GeV} \lesssim E \lesssim 40 \text{ GeV}$,
peak at 3 GeV



Marginalize over mixing parameters:



[PRL 117 (2016) 151803]:

far-to-near ratio

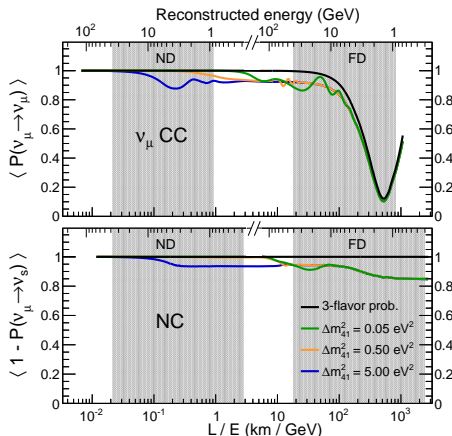
[PRL 122 (2019) 091803]:

full two-detectors fit

[SG+, in preparation]

MINOS & MINOS+

Near (ND, $\simeq 500$ m) and
far (FD, $\simeq 800$ km) detector



[PRL 117 (2016) 151803]:

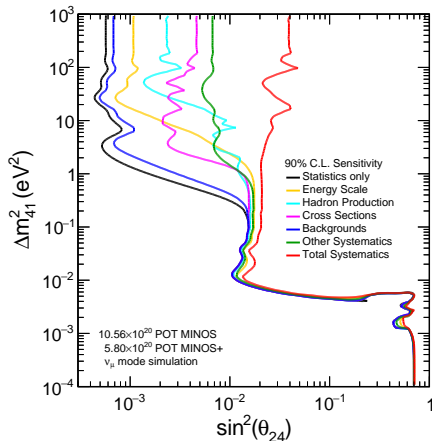
far-to-near ratio

[PRL 122 (2019) 091803]:

full two-detectors fit

$1 \text{ GeV} \lesssim E \lesssim 40 \text{ GeV}$,
peak at 3 GeV

Systematics:

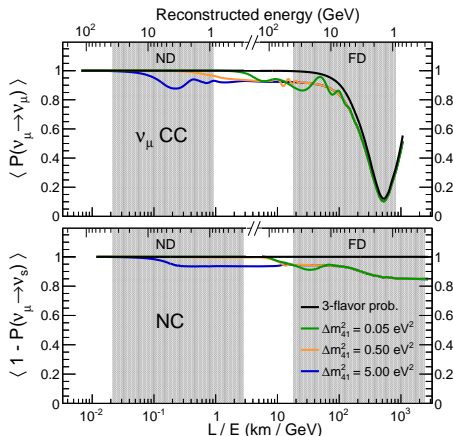


[PRL 122 (2019) 091803]

MINOS & MINOS+

Near (ND, $\simeq 500$ m) and
far (FD, $\simeq 800$ km) detector

$1 \text{ GeV} \lesssim E \lesssim 40 \text{ GeV}$,
peak at 3 GeV



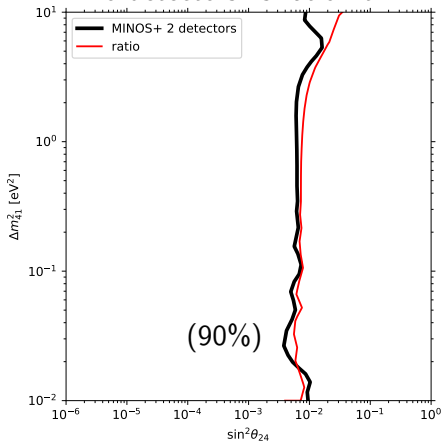
[PRL 117 (2016) 151803]:

far-to-near ratio

[PRL 122 (2019) 091803]:

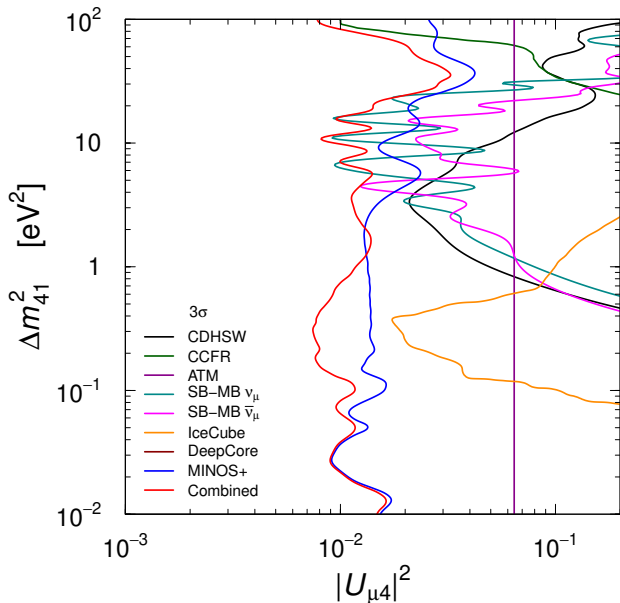
full two-detectors fit

Two detectors vs ratio fit:



[SG+, in preparation]

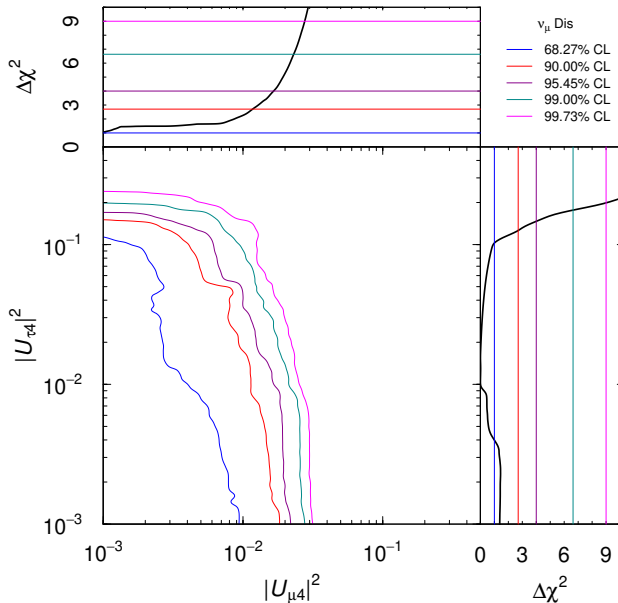
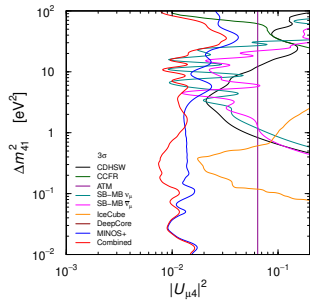
Global fit of $(\bar{\nu}_\mu^-)$ DIS



MINOS+
dominates
at small Δm_{41}^2

IceCube (1 yr)
important at
 $\Delta m_{41}^2 \simeq 0.2 \text{ eV}^2$

see later for
IceCube 8 yr!



1 *Neutrino Oscillations - Some theory*

2 *Electron (anti)neutrino disappearance*

3 *Muon (anti)neutrino disappearance*

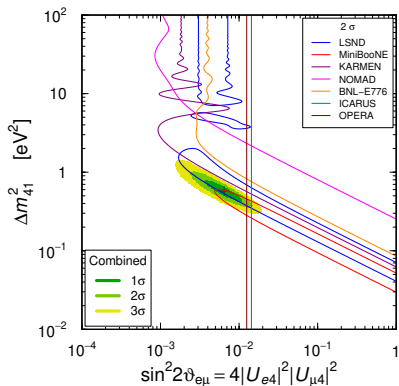
4 ***Electron (anti)neutrino appearance***

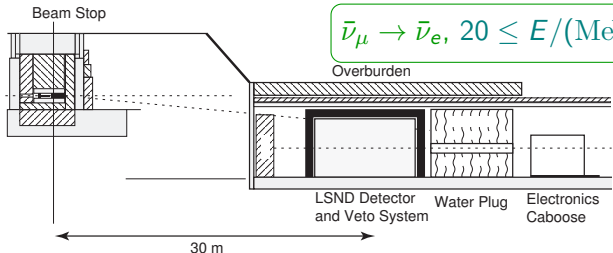
5 *Global fit*

6 *Recent updates*

7 *Light sterile neutrino and cosmology*

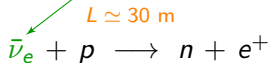
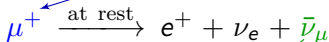
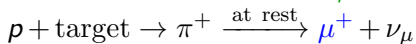
8 *Conclusions*





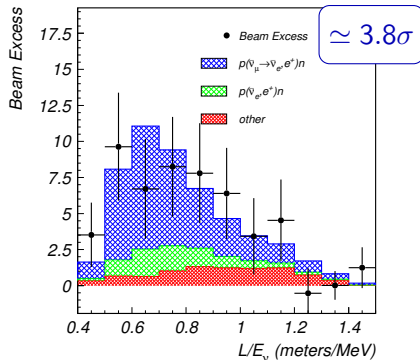
$$\bar{\nu}_\mu \rightarrow \bar{\nu}_e, 20 \leq E/(\text{MeV}) \leq 52.8$$

well known source of $\bar{\nu}_\mu$:



No signal seen in KARMEN ($L \simeq 18 \text{ m}$)

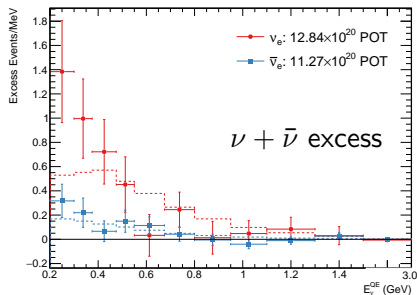
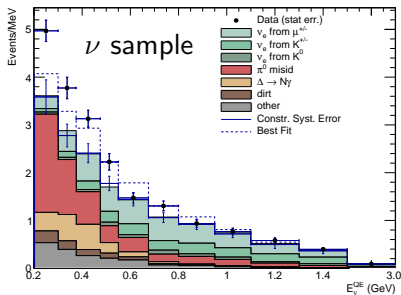
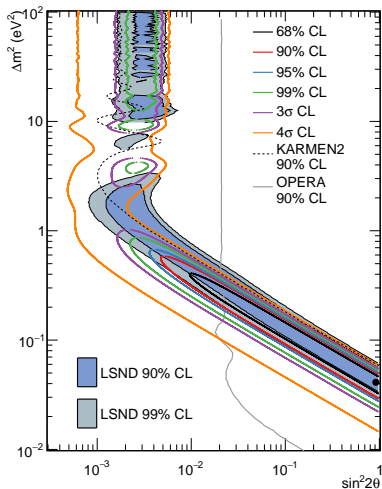
[PRD 65 (2002) 112001]



purpose: check LSND signal

$L \simeq 541$ m, $200 \text{ MeV} \leq E \lesssim 3 \text{ GeV}$

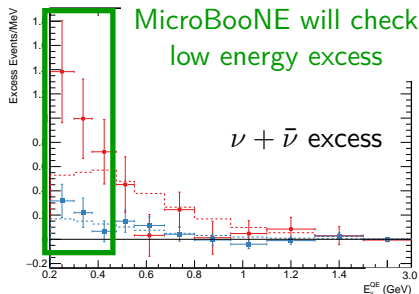
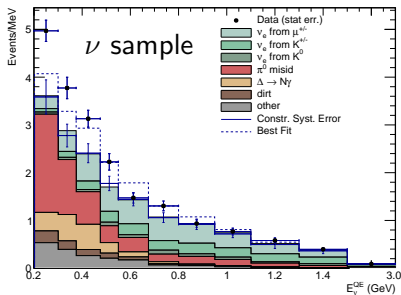
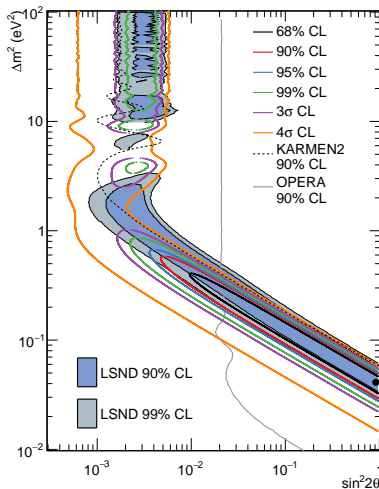
no money, no near detector



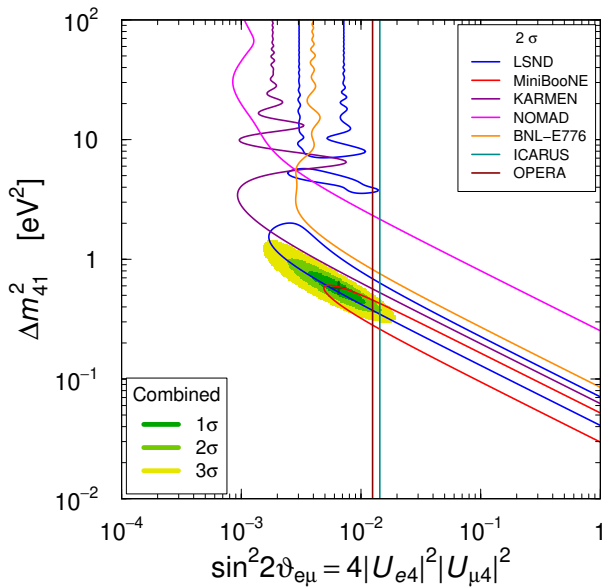
purpose: check LSND signal

$L \simeq 541$ m, $200 \text{ MeV} \leq E \lesssim 3 \text{ GeV}$

no money, no near detector



Global fit of $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ APP



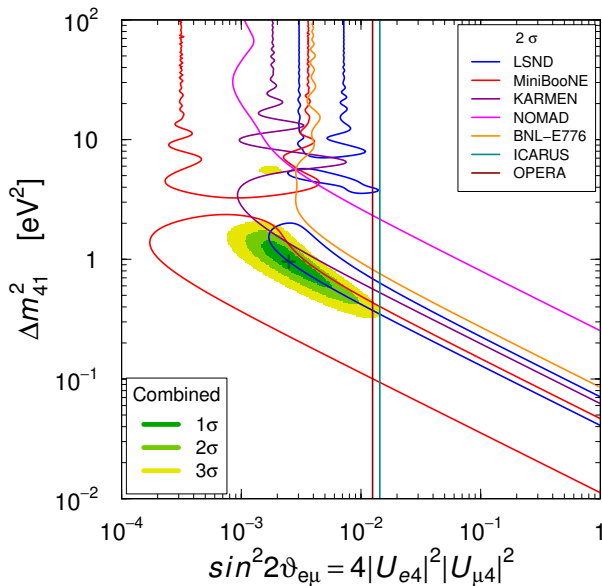
with full MiniBooNE data

ICARUS and OPERA
exclude

MiniBooNE best fit

LSND and MiniBooNE
only partially
in agreement

KARMEN cuts part
of LSND region



ICARUS and OPERA

exclude

MiniBooNE best fit

LSND and MiniBooNE

only partially
in agreement

KARMEN cuts part
of LSND region

without MiniBooNE low energy bins

1 *Neutrino Oscillations - Some theory*

2 *Electron (anti)neutrino disappearance*

3 *Muon (anti)neutrino disappearance*

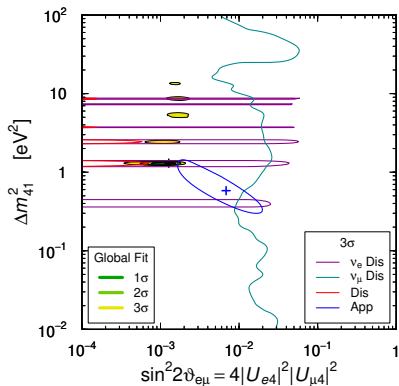
4 *Electron (anti)neutrino appearance*

5 **Global fit**

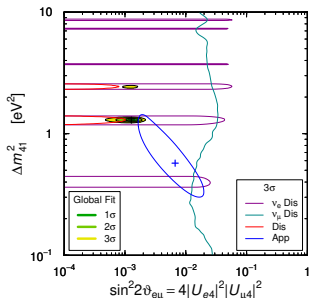
6 *Recent updates*

7 *Light sterile neutrino and cosmology*

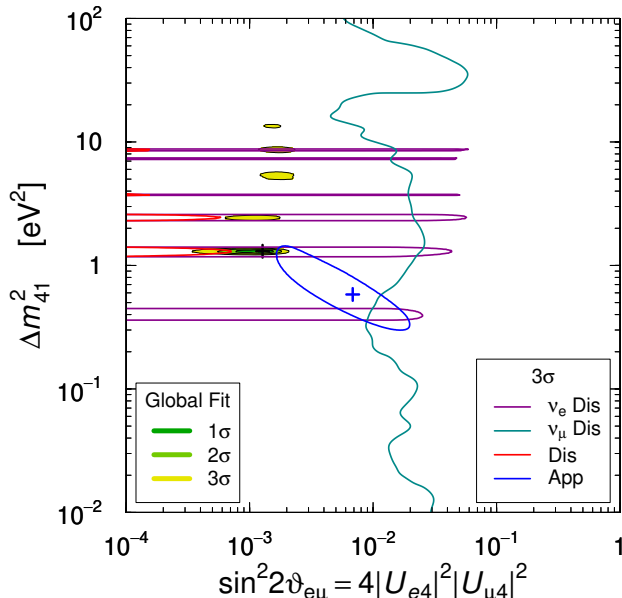
8 *Conclusions*



Status just after
Neutrino 2018:

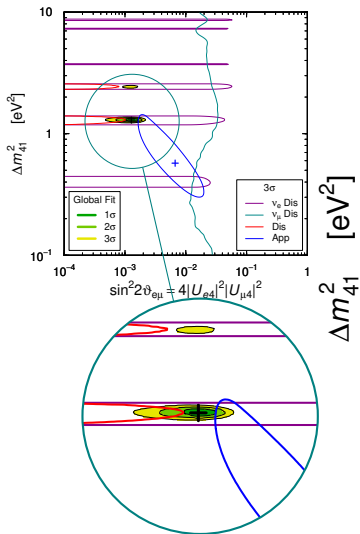


Status in early 2019

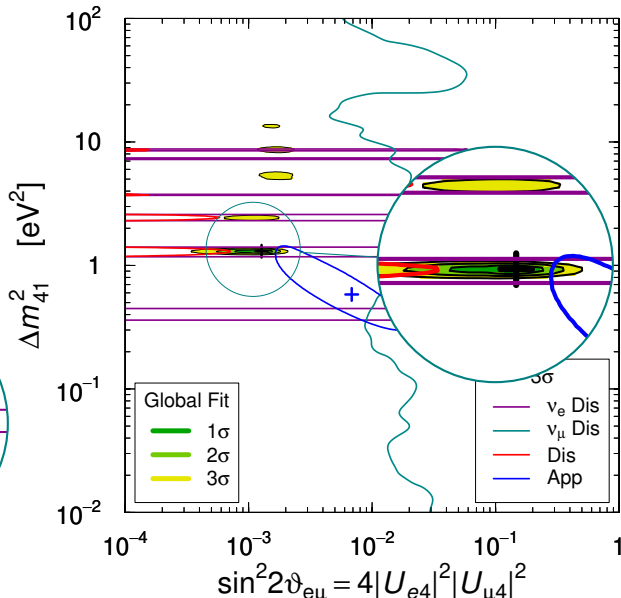


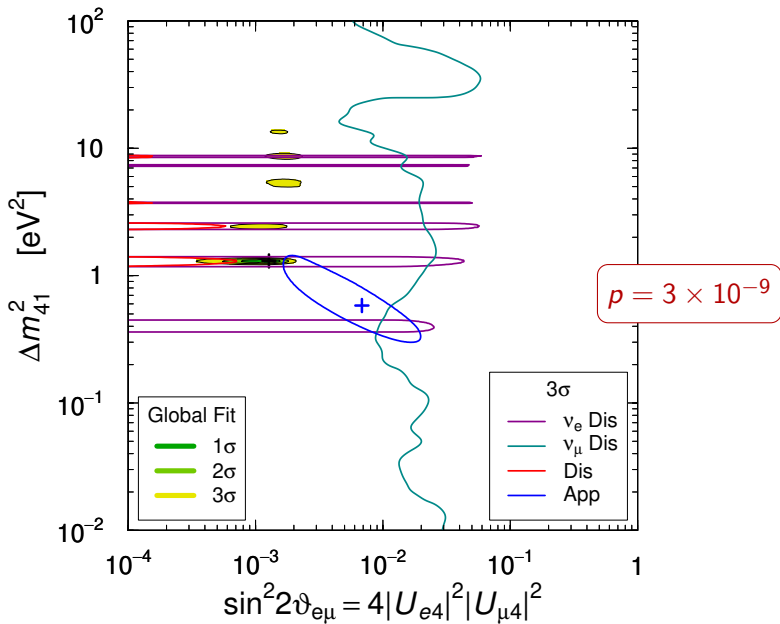
MINOS+ update,
new data
including MiniBooNE
(all bins)

Status just after
Neutrino 2018:



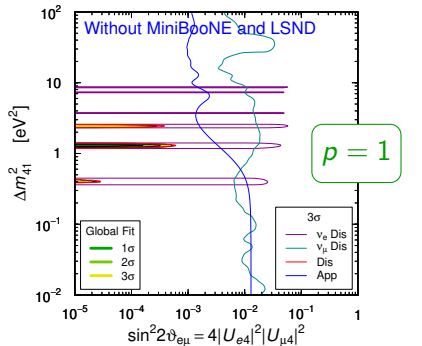
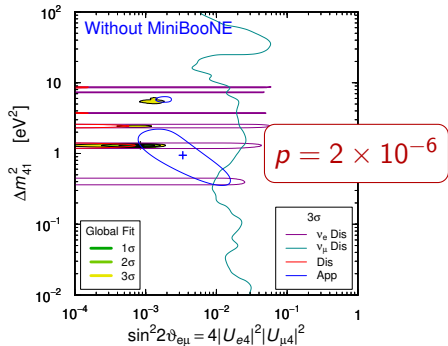
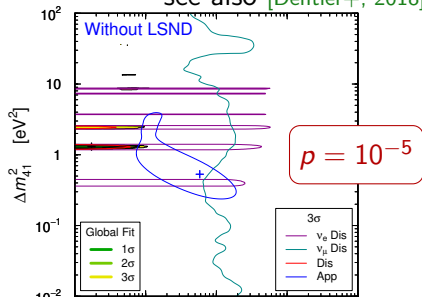
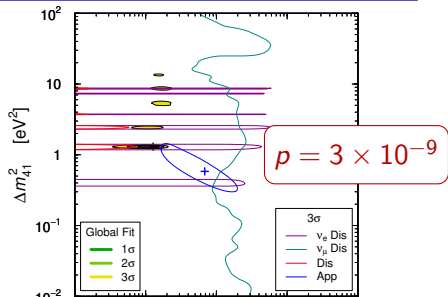
Status in early 2019





APP – DIS tension in 2019

[SG+, in preparation]
see also [Dentler+, 2018]



1 *Neutrino Oscillations - Some theory*

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3 *Muon (anti)neutrino disappearance*

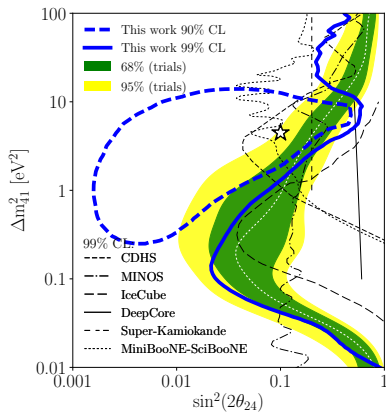
4 *Electron (anti)neutrino appearance*

5 *Global fit*

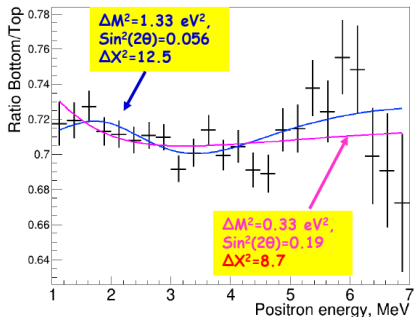
6 **Recent updates**

7 *Light sterile neutrino and cosmology*

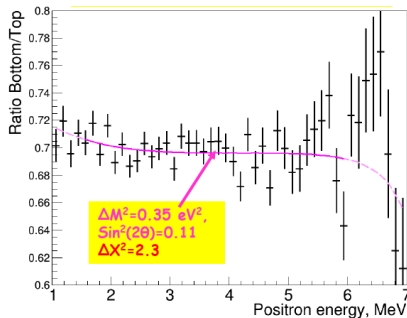
8 *Conclusions*



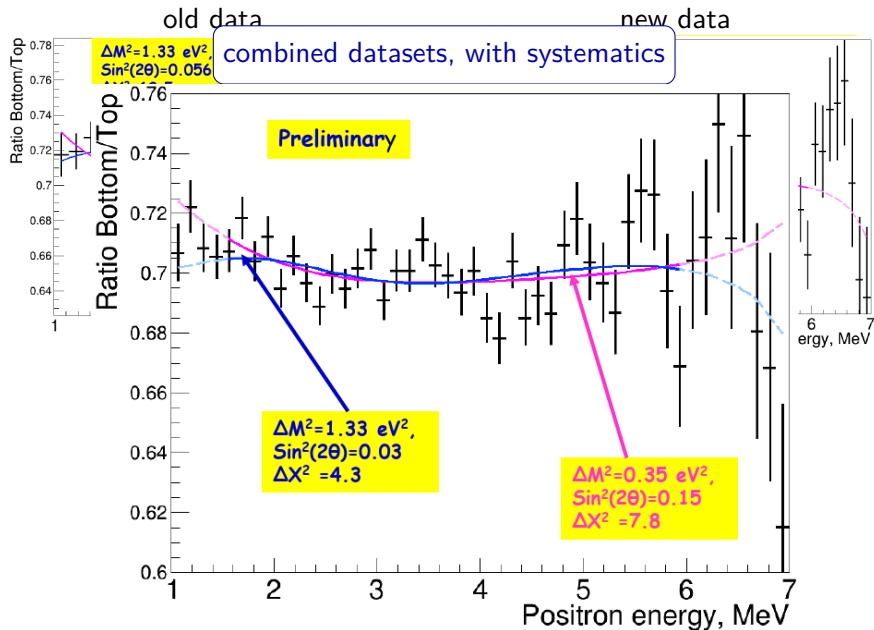
old data

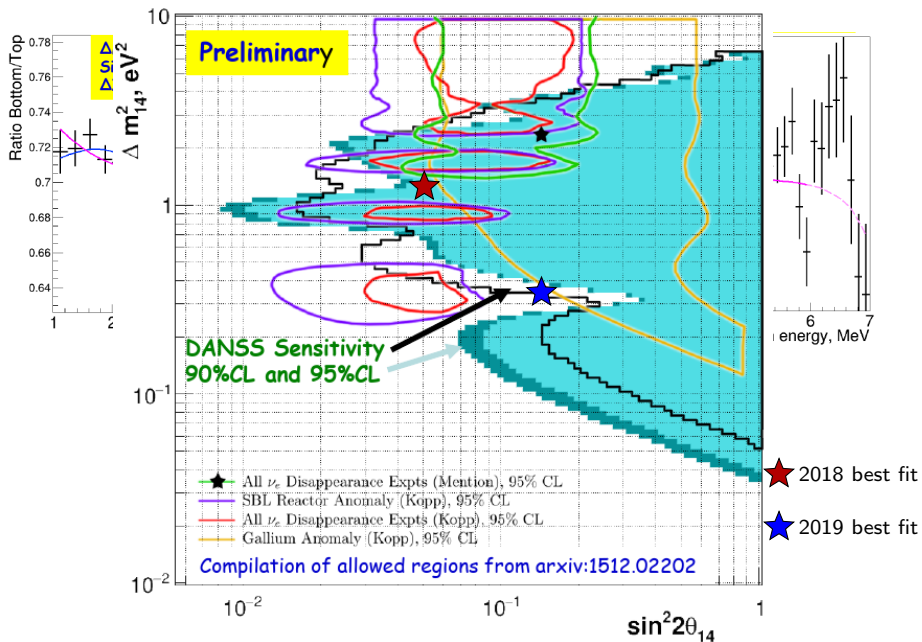


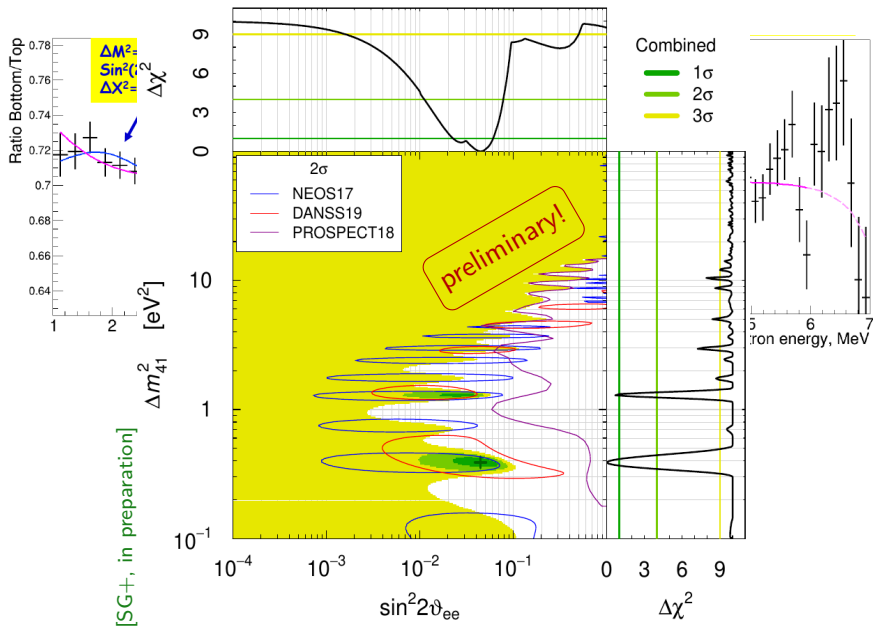
new data

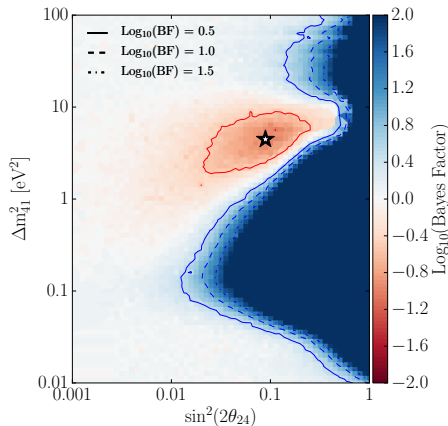
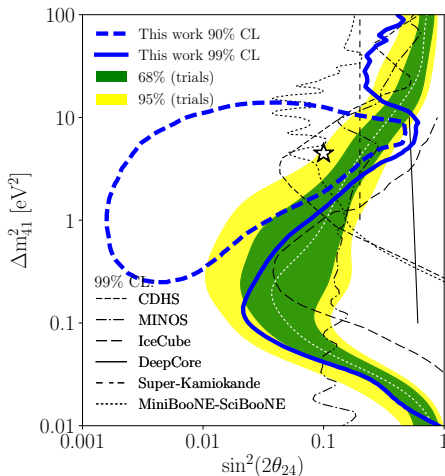


New analysis also
 considers systematics!









first indication in favor of sterile from ν_μ DIS!

although rather weak: $\log_{10} BF \simeq 1$ (weak preference)
or compatible with no oscillations at p -value of 8%

1 *Neutrino Oscillations - Some theory*

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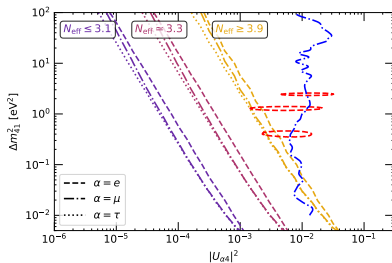
4 *Electron (anti)neutrino appearance*

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8 *Conclusions*



ν oscillations in the early universe

comoving coordinates: $a = 1/T$ $x \equiv m_e a$ $y \equiv p a$ $z \equiv T_\gamma a$ $w \equiv T_\nu a$

$$\text{density matrix: } \varrho(x, y) = \begin{pmatrix} \varrho_{ee} \equiv f_{\nu_e} & \varrho_{e\mu} & \varrho_{e\tau} & \varrho_{es} \\ \varrho_{\mu e} & \varrho_{\mu\mu} \equiv f_{\nu_\mu} & \varrho_{\mu\tau} & \varrho_{\mu s} \\ \varrho_{\tau e} & \varrho_{\tau\mu} & \varrho_{\tau\tau} \equiv f_{\nu_\tau} & \varrho_{\tau s} \\ \varrho_{se} & \varrho_{s\mu} & \varrho_{s\tau} & \varrho_{ss} \equiv f_{\nu_s} \end{pmatrix}$$

$$\frac{d\varrho(y, x)}{dx} = \sqrt{\frac{3m_{\text{Pl}}^2}{8\pi\rho_T}} \left\{ -i \frac{x^2}{m_e^3} \left[\frac{M_F}{2y} - \frac{8\sqrt{2}G_F y m_e^6}{3x^6} \left(\frac{\mathbb{E}_\ell}{m_W^2} + \frac{\mathbb{E}_\nu}{m_Z^2} \right), \varrho \right] + \frac{m_e^3 G_F^2}{(2\pi)^3 x^4 y^2} \mathcal{I}(\varrho) \right\}$$

m_{Pl} Planck mass – ρ_T total energy density – $m_{W,Z}$ mass of the W, Z bosons – G_F Fermi constant – $[\cdot, \cdot]$ commutator

ν oscillations in the early universe

comoving coordinates: $a = 1/T$ $x \equiv m_e a$ $y \equiv p a$ $z \equiv T_\gamma a$ $w \equiv T_\nu a$

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$$\frac{d\varrho(y, x)}{dx} = \sqrt{\frac{3m_{\text{Pl}}^2}{8\pi\rho_T}} \left\{ -i \frac{x^2}{m_e^3} \left[\frac{\mathbb{M}_F}{2y} - \frac{8\sqrt{2}G_F y m_e^6}{3x^6} \left(\frac{\mathbb{E}_\ell}{m_W^2} + \frac{\mathbb{E}_\nu}{m_Z^2} \right), \varrho \right] + \frac{m_e^3 G_F^2}{(2\pi)^3 x^4 y^2} \mathcal{I}(\varrho) \right\}$$

m_{Pl} Planck mass – ρ_T total energy density – $m_{W,Z}$ mass of the W, Z bosons – G_F Fermi constant – $[\cdot, \cdot]$ commutator

$$\mathbb{M}_F = U M U^\dagger$$

$$M = \text{diag}(m_1^2, \dots, m_N^2)$$

$$U = R^{34} R^{24} R^{14} R^{23} R^{13} R^{12} \quad \text{e.g. } R^{14} = \begin{pmatrix} \cos \theta_{14} & 0 & 0 & \sin \theta_{14} \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ -\sin \theta_{14} & 0 & 0 & \cos \theta_{14} \end{pmatrix}$$

$$|U|^2 = \begin{pmatrix} \dots & \dots & \dots & \sin^2 \theta_{14} \\ \dots & \dots & \dots & \cos^2 \theta_{14} \sin^2 \theta_{24} \\ \dots & \dots & \dots & \cos^2 \theta_{14} \cos^2 \theta_{24} \sin^2 \theta_{34} \\ \dots & \dots & \dots & \cos^2 \theta_{14} \cos^2 \theta_{24} \cos^2 \theta_{34} \end{pmatrix}$$

ν oscillations in the early universe

comoving coordinates: $a = 1/T$ $x \equiv m_e a$ $y \equiv p a$ $z \equiv T_\gamma a$ $w \equiv T_\nu a$

$$\text{density matrix: } \varrho(x, y) = \begin{pmatrix} \varrho_{ee} \equiv f_{\nu_e} & \varrho_{e\mu} & \varrho_{e\tau} & \varrho_{es} \\ \varrho_{\mu e} & \varrho_{\mu\mu} \equiv f_{\nu_\mu} & \varrho_{\mu\tau} & \varrho_{\mu s} \\ \varrho_{\tau e} & \varrho_{\tau\mu} & \varrho_{\tau\tau} \equiv f_{\nu_\tau} & \varrho_{\tau s} \\ \varrho_{se} & \varrho_{s\mu} & \varrho_{s\tau} & \varrho_{ss} \equiv f_{\nu_s} \end{pmatrix}$$

$$\frac{d\varrho(y, x)}{dx} = \sqrt{\frac{3m_{\text{Pl}}^2}{8\pi\rho_T}} \left\{ -i \frac{x^2}{m_e^3} \left[\frac{\mathbb{M}_F}{2y} - \frac{8\sqrt{2}G_F y m_e^6}{3x^6} \left(\frac{\mathbb{E}_\ell}{m_W^2} + \frac{\mathbb{E}_\nu}{m_Z^2} \right), \varrho \right] + \frac{m_e^3 G_F^2}{(2\pi)^3 x^4 y^2} \mathcal{I}(\varrho) \right\}$$

m_{Pl} Planck mass – ρ_T total energy density – $m_{W,Z}$ mass of the W, Z bosons – G_F Fermi constant – $[\cdot, \cdot]$ commutator

$$\mathbb{M}_F = U \mathbb{M} U^\dagger$$

$$\mathbb{E}_\ell = \text{diag}(\rho_e, \rho_\mu, 0, 0) \quad \mathbb{E}_\nu = S_a \left(\int dy y^3 \varrho \right) S_a \quad \text{with } S_a = \text{diag}(1, 1, 1, 0)$$

lepton densities

neutrino densities

(only for active neutrinos)

take into account matter effects in oscillations

ν oscillations in the early universe

comoving coordinates: $a = 1/T$ $x \equiv m_e a$ $y \equiv p a$ $z \equiv T_\gamma a$ $w \equiv T_\nu a$

$$\text{density matrix: } \varrho(x, y) = \begin{pmatrix} \varrho_{ee} \equiv f_{\nu_e} & \varrho_{e\mu} & \varrho_{e\tau} & \varrho_{es} \\ \varrho_{\mu e} & \varrho_{\mu\mu} \equiv f_{\nu_\mu} & \varrho_{\mu\tau} & \varrho_{\mu s} \\ \varrho_{\tau e} & \varrho_{\tau\mu} & \varrho_{\tau\tau} \equiv f_{\nu_\tau} & \varrho_{\tau s} \\ \varrho_{se} & \varrho_{s\mu} & \varrho_{s\tau} & \varrho_{ss} \equiv f_{\nu_s} \end{pmatrix}$$

$$\frac{d\varrho(y, x)}{dx} = \sqrt{\frac{3m_{\text{Pl}}^2}{8\pi\rho_T}} \left\{ -i \frac{x^2}{m_e^3} \left[\frac{\mathbb{M}_F}{2y} - \frac{8\sqrt{2}G_F y m_e^6}{3x^6} \left(\frac{\mathbb{E}_\ell}{m_W^2} + \frac{\mathbb{E}_\nu}{m_Z^2} \right), \varrho \right] + \frac{m_e^3 G_F^2}{(2\pi)^3 x^4 y^2} \mathcal{I}(\varrho) \right\}$$

m_{Pl} Planck mass – ρ_T total energy density – $m_{W,Z}$ mass of the W, Z bosons – G_F Fermi constant – $[\cdot, \cdot]$ commutator

$$\mathbb{M}_F = U M U^\dagger \quad \mathbb{E}_\ell = \text{diag}(\rho_e, \rho_\mu, 0, 0) \quad \mathbb{E}_\nu = S_a \left(\int dy y^3 \varrho \right) S_a$$

$\mathcal{I}(\varrho)$ collision integrals

take into account neutrino-electron scattering and pair annihilation

2D integrals over the momentum, take most of the computation time

ν oscillations in the early universe

comoving coordinates: $a = 1/T$ $x \equiv m_e a$ $y \equiv p a$ $z \equiv T_\gamma a$ $w \equiv T_\nu a$

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m_{Pl} Planck mass – ρ_T total energy density – $m_{W,Z}$ mass of the W, Z bosons – G_F Fermi constant – $[\cdot, \cdot]$ commutator

$$\mathbb{M}_F = U \mathbb{M} U^\dagger$$

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$\mathcal{I}(\varrho)$ collision integrals

from continuity
equation
 $\dot{\rho} = -3H(\rho + P)$

$$\frac{dz}{dx} = \frac{\sum_{\ell=e,\mu} \left[\frac{r_\ell^2}{r} J(r_\ell) \right] + G_1(r) - \frac{1}{2\pi^2 z^3} \int_0^\infty dy y^3 \sum_{\alpha=e} \frac{d\varrho_{\alpha\alpha}}{dx}}{\sum_{\ell=e,\mu} \left[r_\ell^2 J(r_\ell) + Y(r_\ell) \right] + G_2(r) + \frac{2\pi^2}{15}}$$

$r = x/z$, $r_\ell = m_\ell/m_e r$ $J(r)$, $Y(r)$ from non-relativistic transition of e^\pm , μ^\pm
 $G_1(r)$ and $G_2(r)$ from electromagnetic corrections

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FORTran-Evolved Primordial Neutrino Oscillations (FortEPiano)

https://bitbucket.org/ahp_cosmo/fortepiano

from continuity equation

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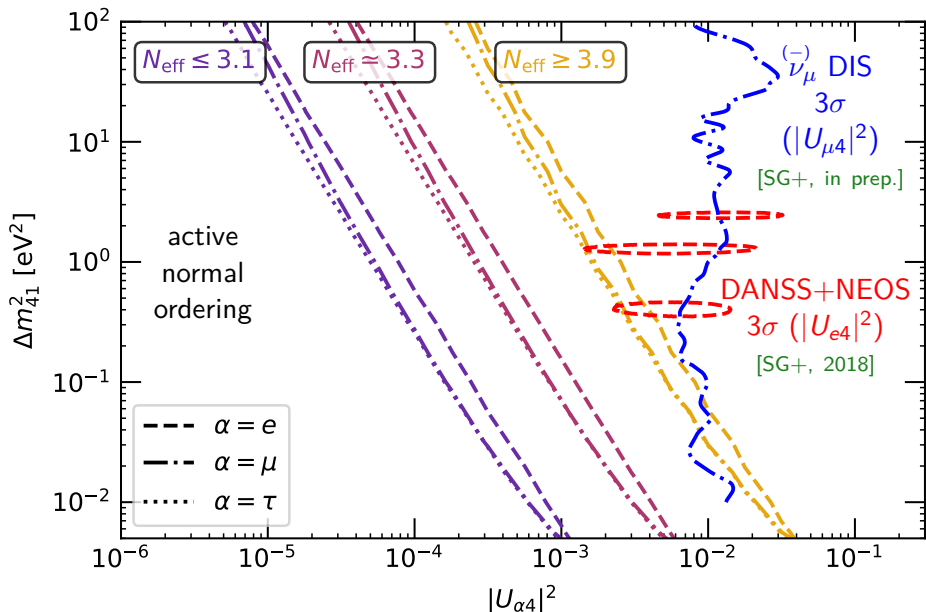
will be public soon

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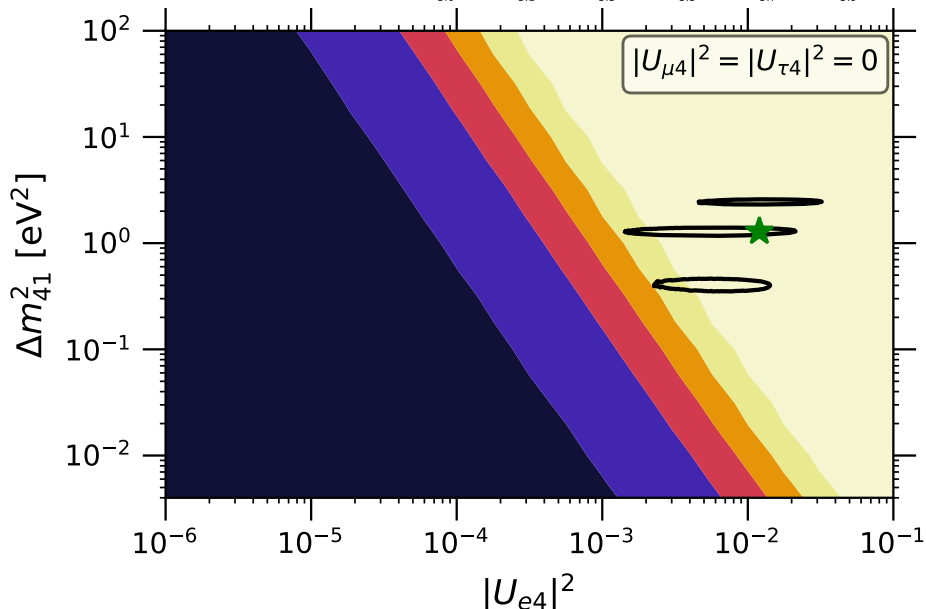
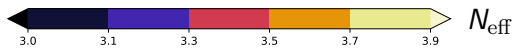
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N_{eff} and the new mixing parameters

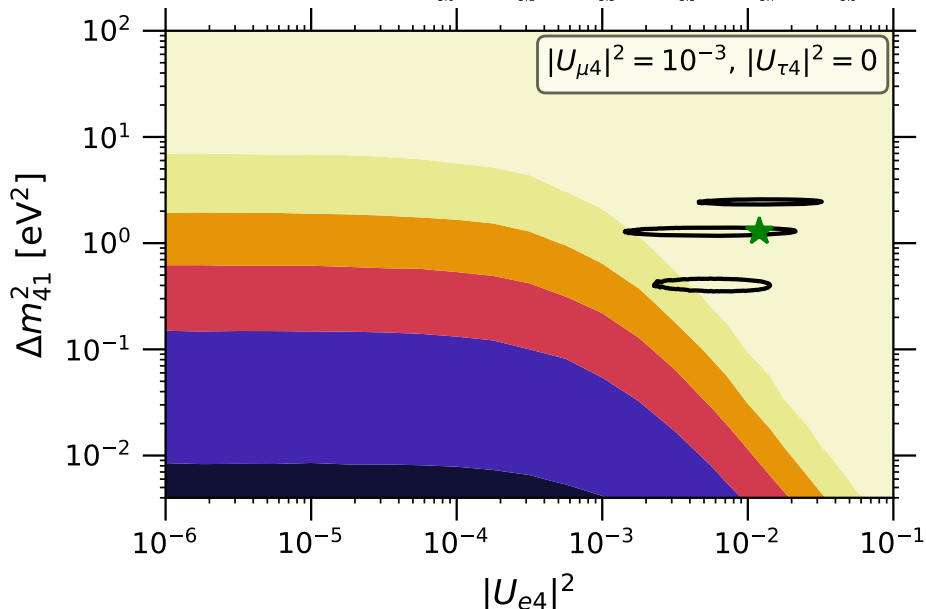
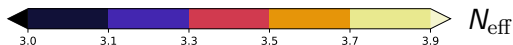
Only vary one angle and fix two to zero: do they have the same effect?



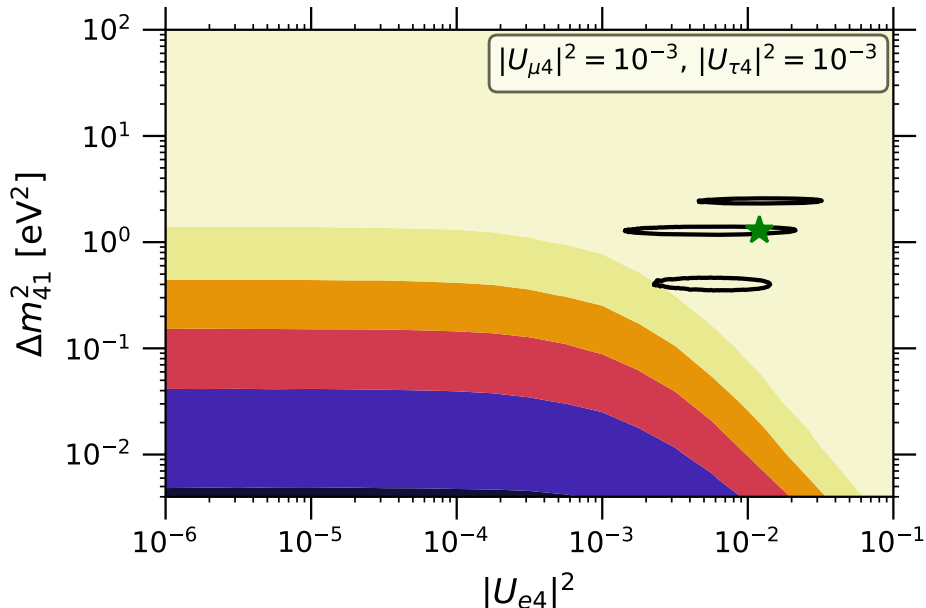
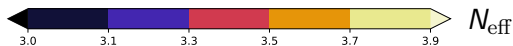
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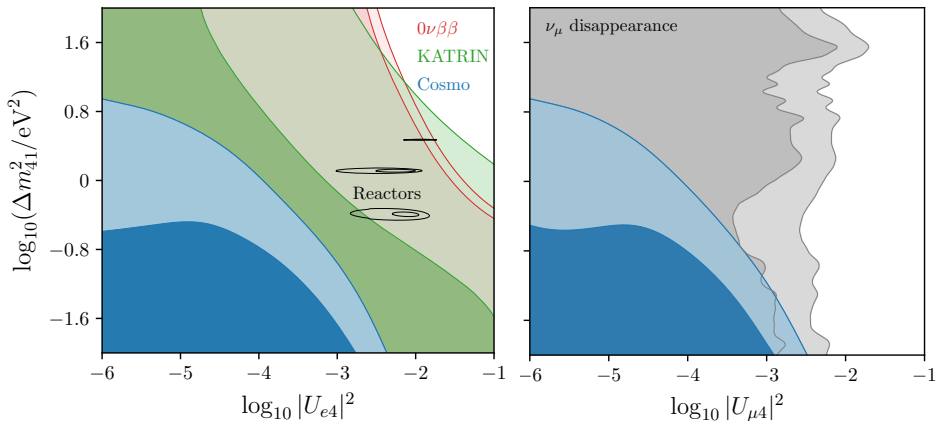
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Comparing constraints

Cosmological constraints are stronger than most other probes

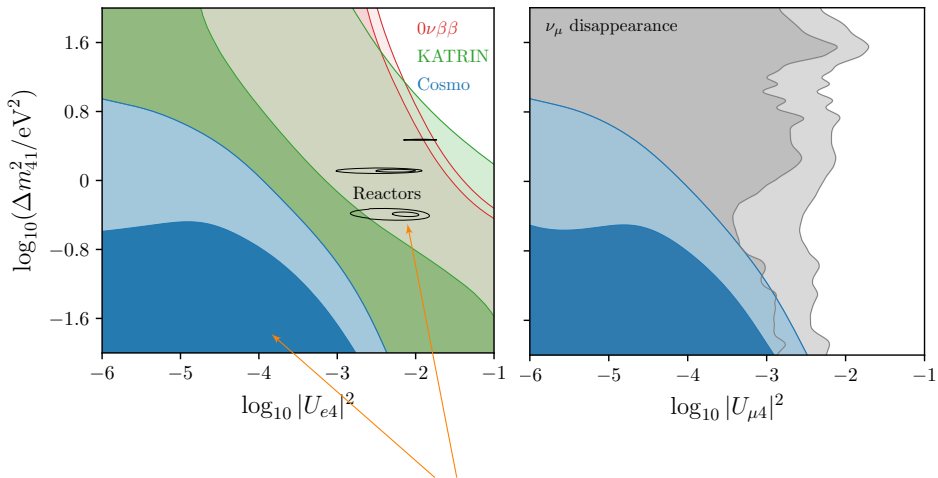
But much more model dependent (as all the cosmological constraints)!



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Warning: tension between reactor experiments and CMB bounds!

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- 2 *Electron (anti)neutrino disappearance*
- 3 *Muon (anti)neutrino disappearance*
- 4 *Electron (anti)neutrino appearance*
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- 6 *Recent updates*
- 7 *Light sterile neutrino and cosmology*
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Conclusions

1

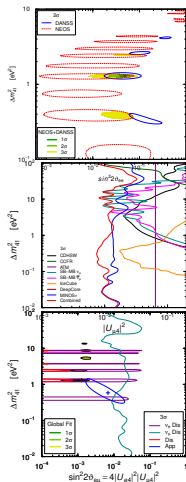
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with Gallium anomaly and RAA

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but also constraints on $|U_{\tau 4}|^2$

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strong APP-DIS tension
What are LSND/MiniBooNE observing?
Systematics or $LS\nu$ or new physics?



Conclusions

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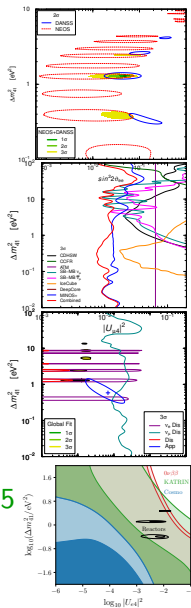
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Thank you for the attention!

