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Neutrino decoupling in standard and non-standard scenarios

Based on JCAP 04 (2021) 073, JCAP 07 (2019) 014, JCAP 03 (2023) 046

TAsP meeting, Turin, 18-19/01/2024

1 Cosmic Neutrino Background

2 Standard three neutrino scenario

3 Non-standard 1: light sterile neutrino

4 Non-standard 2: non-unitarity

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5 **Conclusions**

History of the universe

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History of the universe

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History of the universe

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Relic neutrinos in cosmology: $N_{\rm eff}$

$$\widehat{\rho_r = \left[1 + \frac{7}{8} \left(\frac{4}{11}\right)^{4/3} N_{\text{eff}}\right] \rho_{\gamma}}$$

 ρ_{γ} photon energy density, 7/8 for fermions, $(4/11)^{4/3}$ due to photon reheating after neutrino decoupling prediction: measurement:

instantaneous decoupling: $N_{\rm eff}~=~1$ for each u family

> 3 because of entropy transfer to photons when electrons become non-relativistic

recommended value (3
$$\nu$$
):
 $N_{\rm eff} = 3.04$
[Bennett+, 2020] [Akita+, 2020]
[Froustey+, 2020] [Cielo+, 2023]



TAsP, 18/01/2024

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before BBN: neutrinos coupled to plasma ($\nu_{\alpha}\bar{\nu}_{\alpha} \leftrightarrow e^+e^-$, $\nu e \leftrightarrow \nu e$)



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 ν decouple mostly before $e^+e^- \to \gamma\gamma$ annihilation!

before BBN: neutrinos coupled to plasma ($\nu_{\alpha}\bar{\nu}_{\alpha} \leftrightarrow e^+e^-$, $\nu e \leftrightarrow \nu e$)



[Bennett, SG+, JCAP 2021] [Sigl, Raffelt, 1993] ν oscillations in the early universe comoving coordinates: a = 1/T $x \equiv m_e a$ $y \equiv p a$ $z \equiv T_{\gamma} a$ $w \equiv T_{\nu} a$ $\begin{array}{ccc} \text{density matrix:} & \varrho(x,y) = \begin{pmatrix} \varrho_{ee} \equiv f_{\nu_e} & \varrho_{e\mu} & \varrho_{e\tau} \\ \varrho_{\mu e} & \varrho_{\mu\mu} \equiv f_{\nu_{\mu}} & \varrho_{\mu\tau} \\ \varrho_{\tau e} & \varrho_{\tau\mu} & \varrho_{\tau\tau} \equiv f_{\nu_{\tau}} \end{pmatrix}$ off-diagonals to take into account coherency in the neutrino system ϱ evolution from $x H \frac{\mathrm{d}\varrho(y,x)}{\mathrm{d}x} = -ia[\mathcal{H}_{\mathrm{eff}},\varrho] + b\mathcal{I}$ *H* Hubble factor \rightarrow expansion (depends on universe content) effective Hamiltonian $\mathcal{H}_{eff} = \frac{\mathbb{M}_{F}}{2y} - \frac{2\sqrt{2}G_{F}ym_{e}^{6}}{x^{6}} \left(\frac{\mathbb{E}_{\ell} + \mathbb{P}_{\ell}}{m_{in}^{2}} + \frac{4}{3} \frac{\mathbb{E}_{\nu}}{m_{\gamma}^{2}}\right)$ vacuum oscillations + → matter effects \mathcal{I} collision integrals take into account $\nu - e$ scattering and pair annihilation, $\nu - \nu$ interactions

2D integrals over momentum, take most of the computation time

solve together with z evolution, from $x \frac{d\rho(x)}{dx} = \rho - 3P$

 $\rho,\,P$ total energy density and pressure, also take into account FTQED corrections

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"Neutrino decoupling in standard and non-standard scenarios"

TAsP, 18/01/2024

[Bennett, SG+, JCAP 2021] ν oscillations in the early universe [Sigl, Raffelt, 1993] comoving coordinates: a = 1/T $x \equiv m_e a$ $y \equiv p a$ $z \equiv T_{\gamma} a$ $w \equiv T_{\nu} a$ $\begin{array}{ll} \text{density matrix:} & \varrho(x,y) = \left(\begin{array}{cc} \varrho_{ee} \equiv f_{\nu_e} & \varrho_{e\mu} & \varrho_{e\tau} \\ \varrho_{\mu e} & \varrho_{\mu\mu} \equiv f_{\nu_{\mu}} & \varrho_{\mu\tau} \\ \varrho_{\tau e} & \varrho_{\tau\mu} & \varrho_{\tau\tau} \equiv f_{\nu_{-}} \end{array} \right) \end{array}$ off-diagonals to take into account coherency in the neutrino system ρ evolution from $xH\frac{\mathrm{d}\rho(y,x)}{\mathrm{d}x} = -ia[\mathcal{H}_{\mathrm{eff}},\rho] + b\mathcal{I}$ FORTran-Evolved PrimordIAl Neutrino Oscillations (FortEPiaNO) https://bitbucket.org/ahep cosmo/fortepiano public vacuum oscillations + → matter effects \mathcal{I} collision integrals take into account $\nu - e$ scattering and pair annihilation, $\nu - \nu$ interactions 2D integrals over momentum, take most of the computation time solve together with z evolution, from $x \frac{d\rho(x)}{dx} = \rho - 3P$ ρ , P total energy density and pressure, also take into account FTQED corrections S. Gariazzo "Neutrino decoupling in standard and non-standard scenarios" TAsP, 18/01/2024 4/15

Neutrino momentum distribution and $N_{\rm eff}$ [Bennett, SG+, JCAP 2021]

Distortion of the momentum distribution (f_{FD} : Fermi-Dirac at equilibrium)



Neutrino momentum distribution and $N_{\rm eff}$ [Bennett, SG+, JCAP 2021]



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Neutrino momentum distribution and $N_{\rm eff}$ [Bennett, SG+, JCAP 2021]

$$N_{\text{eff}}^{\text{any time}} = \frac{8}{7} \left(\frac{T_{\gamma}}{T_{\nu}}\right)^4 \frac{\rho_{\nu}}{\rho_{\gamma}} = \frac{8}{7} \left(\frac{T_{\gamma}}{T_{\nu}}\right)^4 \frac{1}{\rho_{\gamma}} \sum_i g_i \int \frac{d^3 p}{(2\pi)^3} E(p) f_{\nu,i}(p)$$



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"Neutrino decoupling in standard and non-standard scenarios"

[Bennett, SG+, JCAP 2021]

Effect of neutrino oscillations



[Bennett, SG+, JCAP 2021]

Effect of neutrino oscillations



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Full 3ν mixing results:



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[SG+, JPG 43 (2016) 033001]

Do three-neutrino oscillations explain all experimental results?

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[SG+, JPG 43 (2016) 033001]







 $\frac{\nu \text{ oscillations in the early universe}}{\text{comoving coordinates: } a = 1/T \quad x \equiv m_e \text{ a} \quad y \equiv p \text{ a}} \begin{bmatrix} \text{Bennett, SG+, JCAP 2021} \\ [Sigl, Raffelt, 1993] \\ z \equiv T_{\gamma} \text{ a} \quad w \equiv T_{\nu} \text{ a} \end{bmatrix}$ $\frac{e_{ee} \equiv f_{\nu_e}}{e_{\mu e}} e_{\mu\mu} \equiv f_{\nu\mu} e_{\mu\tau} e_{\mu\tau$

nsity matrix:
$$\rho(x, y) = \begin{pmatrix} \rho_{\mu}e & \rho_{\mu} & \rho_{\mu} & \rho_{\mu}r \\ \rho_{\tau e} & \rho_{\tau \mu} & \rho_{\tau \tau} & \rho_{\tau s} \\ \rho_{s e} & \rho_{s \mu} & \rho_{s \tau} & \rho_{s s} \\ \rho_{s e} & \rho_{s \mu} & \rho_{s \tau} & \rho_{s s} \\ \rho_{s e} & \rho_{s \mu} & \rho_{s \tau} & \rho_{s s} \\ \rho_{s e} & \rho_{s \mu} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s e} & \rho_{s \mu} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s e} & \rho_{s \mu} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s e} & \rho_{s \mu} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s e} & \rho_{s \mu} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s e} & \rho_{s \mu} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s e} & \rho_{s \mu} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s e} & \rho_{s \mu} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s e} & \rho_{s \mu} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s e} & \rho_{s \mu} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s e} & \rho_{s \mu} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s e} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s e} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s e} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s e} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s e} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s e} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s e} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s e} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s e} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s e} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s e} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s e} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} & \rho_{s \tau} \\ \rho_{s \tau} & \rho_{s \tau} & \rho_{s$$

$$\varrho$$
 evolution from $xH\frac{\mathrm{d}\varrho(y,x)}{\mathrm{d}x} = -ia[\mathcal{H}_{\mathrm{eff}},\varrho] + b\mathcal{I}$

H Hubble factor \rightarrow expansion (depends on universe content)

effective Hamiltonian
$$\mathcal{H}_{\text{eff}} = \frac{\mathbb{M}_{\text{F}}}{2y} - \frac{2\sqrt{2}G_{\text{F}}ym_{e}^{6}}{x^{6}} \left(\frac{\mathbb{E}_{\ell} + \mathbb{P}_{\ell}}{m_{W}^{2}} + \frac{4}{3}\frac{\mathbb{E}_{\nu}}{m_{Z}^{2}}\right)$$

vacuum oscillations \longleftarrow matter effects

$\mathcal I$ collision integrals

take into account $\nu-e$ scattering and pair annihilation, $\nu-\nu$ interactions

2D integrals over momentum, take most of the computation time

solve together with z evolution, from
$$x \frac{d\rho(x)}{dx} = \rho - 3P$$

 $\rho,\,P$ total energy density and pressure, also take into account FTQED corrections

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[SG+, JCAP 07 (2019) 014]



$N_{\rm eff}$ and the new mixing parameters

[SG+, JCAP 07 (2019) 014]



Comparing constraints

Cosmological constraints are stronger than most other probes

But much more model dependent (as all the cosmological constraints)!



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Warning: tension between reactor experiments and CMB bounds!

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[JCAP 03 (2023)]

Non-unitarity of the 3×3 mixing matrix

Consider we have N_{ν} neutrino states

Unitary
$$N_{\nu} \times N_{\nu}$$
 mixing matrix: $V = \begin{pmatrix} V_{e1} & V_{e2} & V_{e3} \\ V_{\mu 1} & V_{\mu 2} & V_{\mu 3} \\ V_{\tau 1} & V_{\tau 2} & V_{\tau 3} \\ \vdots & \ddots \end{pmatrix}$

the 3×3 sector (N)

describing mixing among lightest neutrinos is non-unitary

$$N = \begin{pmatrix} \alpha_{11} & 0 & 0 \\ \alpha_{21} & \alpha_{22} & 0 \\ \alpha_{31} & \alpha_{32} & \alpha_{33} \end{pmatrix} U$$

 α_{ii} real, α_{ij} $(i \neq j)$ complex \Rightarrow CP violation

 $U = R^{23}R^{13}R^{12}$ is the standard unitary mixing matrix

[JCAP 03 (2023)]

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 mixing matrix: $V = \begin{pmatrix} V_{e1} & V_{e2} & V_{e3} & \dots & V_{\mu 1} & V_{\mu 2} & V_{\mu 3} \\ V_{\pi 1} & V_{\pi 2} & V_{\pi 3} & \dots & V_{\pi 1} & V_{\pi 2} & V_{\pi 3} \\ \vdots & \vdots & \ddots & \ddots \end{pmatrix}$

the 3×3 sector (*N*) describing mixing among lightest neutrinos is non-unitary

Neutrino interactions depend only on kinematically accessible states Oscillations depend on all states

Oscillations with states n > 3 much heavier than $n \le 3$ are averaged out at experiments

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Non-unitarity and neutrino decoupling

Neutrino density matrix evolution in mass basis:

$$\frac{\mathrm{d}\varrho(y)}{\mathrm{d}x}\Big|_{\mathrm{M}} = \sqrt{\frac{3m_{\mathrm{Pl}}^2}{8\pi\rho}} \left\{ -i\frac{x^2}{m_e^3} \left[\frac{\mathbb{M}_{\mathrm{M}}}{2y} - \frac{2\sqrt{2}G_F y m_e^6}{x^6} \mathcal{E}_{\mathrm{M}}, \varrho \right] + \frac{m_e^3}{x^4} \mathcal{I}(\varrho) \right\}$$

Unitary case

interactions: $(Y_L)_{ab} \equiv \tilde{g}_L \mathbb{I} + (U^{\dagger})_{ea} U_{eb}$ $(Y_R)_{ab} \equiv g_R \mathbb{I}$ Non-unitary case

interactions:

$$\begin{array}{lcl} (Y_L)_{ab} &\equiv & \tilde{g}_L(V^{\dagger}V)_{ab} + (V^{\dagger})_{ea}V_{eb} \\ (Y_R)_{ab} &\equiv & g_R(V^{\dagger}V)_{ab} \end{array}$$

 matter effects: $\mathcal{E}_{\rm NU} \equiv \frac{\rho_e + P_e}{m_W^2} (Y_L - Y_R)$

Fermi constant: $G_F^{\mu} = G_F$ $G_F^{\mu} = G_F \sqrt{\alpha_{11}^2 (\alpha_{22}^2 + |\alpha_{21}|^2)}$ $G_F^{\mu} = 1.1663787(6) \times 10^{-5} \text{ GeV}^{-2} \text{ [CODATA]}$ $\mathcal{I}(\varrho) \propto G_F^2$

[JCAP 03 (2023)]

Non-unitarity parameters and $N_{\rm eff}$

[JCAP 03 (2023)]



Non-unitarity parameters and $N_{\rm eff}$

[JCAP 03 (2023)]



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Conclusions

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Neutrinos in the early universe - probe lowest energies







Active neutrinos: precision calculations



Non-standard scenarios: complementary bounds



Conclusions

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Non-standard scenarios: complementary bounds



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