

La Fisica dei Neutrini: Stato e Prospettive

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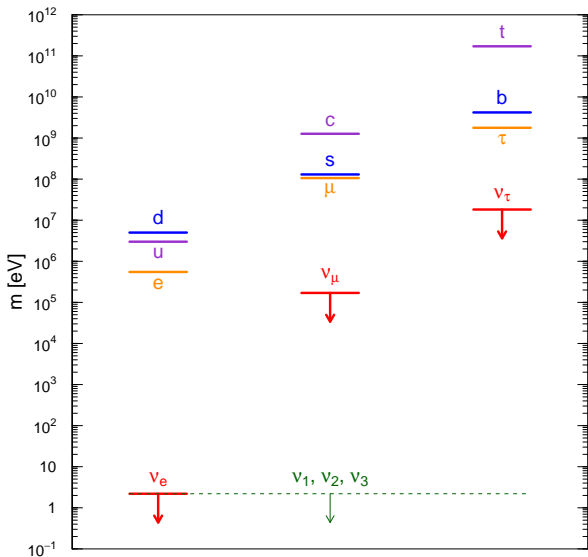
Neutrino Unbound: <http://www.nu.to.infn.it>

IFAE2011

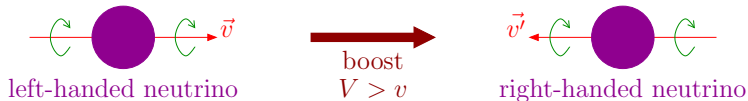
Incontri di Fisica delle Alte Energie

Perugia, 27–29 Aprile 2011

Fermion Mass Spectrum



Standard Model: Massless Neutrinos



Standard Model: $\nu_L \implies$ no Dirac mass term

$$\mathcal{L}^D = -m^D (\overline{\nu}_R \nu_L + \overline{\nu}_L \nu_R) \quad (\text{no } \nu_R)$$

Majorana Neutrino: $\nu = \nu^c \implies \nu_R = \nu_L^c \implies$ Majorana mass term

$$\mathcal{L}^M = -\frac{1}{2} m^M (\overline{\nu}_L^c \nu_L + \overline{\nu}_L \nu_L^c)$$

Standard Model: Majorana mass term **not** allowed by $SU(2)_L \times U(1)_Y$
(no Higgs triplet)

Extension of the SM: Massive Neutrinos

Standard Model can be extended with ν_R

Dirac neutrino mass term $\mathcal{L}^D = -m^D (\overline{\nu_R} \nu_L + \overline{\nu_L} \nu_R) \Rightarrow m^D \lesssim 100 \text{ GeV}$

surprise: Majorana neutrino mass for ν_R is allowed!

$$\mathcal{L}_R^M = -\frac{1}{2} m_R^M (\overline{\nu_R^c} \nu_R + \overline{\nu_R} \nu_R^c)$$

total neutrino mass term $\mathcal{L}^{D+M} = -\frac{1}{2} \begin{pmatrix} \overline{\nu_L^c} & \overline{\nu_R} \end{pmatrix} \begin{pmatrix} 0 & m^D \\ m^D & m_R^M \end{pmatrix} \begin{pmatrix} \nu_L \\ \nu_R^c \end{pmatrix} + \text{H.c.}$

m_R^M can be arbitrarily large (not protected by SM symmetries)

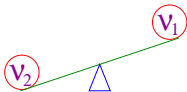
$m_R^M \sim$ scale of new physics beyond Standard Model $\Rightarrow m_R^M \gg m^D$

diagonalization of $\begin{pmatrix} 0 & m^D \\ m^D & m_R^M \end{pmatrix} \Rightarrow m_1 \simeq \frac{(m^D)^2}{m_R^M}, \quad m_2 \simeq m_R^M$

natural explanation of smallness
of light neutrino masses

massive neutrinos are Majorana!

3-GEN \Rightarrow effective low-energy 3- ν mixing



see-saw mechanism

[Minkowski, PLB 67 (1977) 42]

[Yanagida (1979); Gell-Mann, Ramond, Slansky (1979); Mohapatra, Senjanovic, PRL 44 (1980) 912]

Three-Neutrino Mixing

$$\mathcal{L}_{\text{mass}} \sim \begin{pmatrix} \bar{\nu}_e & \bar{\nu}_\mu & \bar{\nu}_\tau \end{pmatrix} \begin{pmatrix} m_{ee} & m_{e\mu} & m_{e\tau} \\ m_{\mu e} & m_{\mu\mu} & m_{\mu\tau} \\ m_{\tau e} & m_{\tau\mu} & m_{\tau\tau} \end{pmatrix} \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix}$$

diagonalization of mass matrix

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

$$\nu_\alpha = \sum_{k=1}^3 U_{\alpha k} \nu_k \quad (\alpha = e, \mu, \tau)$$

$$\mathcal{L}_{\text{mass}} \sim \begin{pmatrix} \bar{\nu}_1 & \bar{\nu}_2 & \bar{\nu}_3 \end{pmatrix} \begin{pmatrix} m_1 & 0 & 0 \\ 0 & m_2 & 0 \\ 0 & 0 & m_3 \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix} = \sum_{k=1}^3 m_k \bar{\nu}_k \nu_k$$

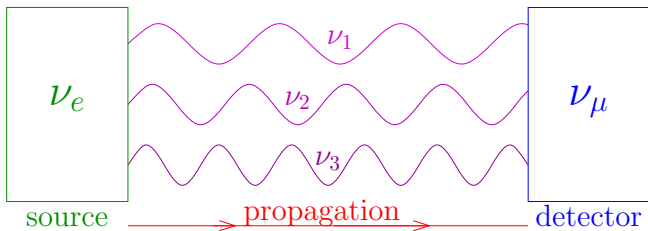
Neutrino Oscillations

- ▶ 1957: Bruno Pontecorvo proposed Neutrino Oscillations in analogy with $K^0 \Leftrightarrow \bar{K}^0$ oscillations (Gell-Mann and Pais, 1955)
- ▶ Flavor Neutrinos: ν_e, ν_μ, ν_τ produced in Weak Interactions
- ▶ Massive Neutrinos: ν_1, ν_2, ν_3 propagate from Source to Detector
- ▶ A Flavor Neutrino is a **superposition** of Massive Neutrinos

$$\begin{aligned} |\nu_e\rangle &= U_{e1} |\nu_1\rangle + U_{e2} |\nu_2\rangle + U_{e3} |\nu_3\rangle \\ |\nu_\mu\rangle &= U_{\mu1} |\nu_1\rangle + U_{\mu2} |\nu_2\rangle + U_{\mu3} |\nu_3\rangle \\ |\nu_\tau\rangle &= U_{\tau1} |\nu_1\rangle + U_{\tau2} |\nu_2\rangle + U_{\tau3} |\nu_3\rangle \end{aligned}$$

- ▶ U is the 3×3 Neutrino Mixing Matrix

$$|\nu(t=0)\rangle = |\nu_e\rangle = U_{e1} |\nu_1\rangle + U_{e2} |\nu_2\rangle + U_{e3} |\nu_3\rangle$$



$$|\nu(t > 0)\rangle = U_{e1} e^{-iE_1 t} |\nu_1\rangle + U_{e2} e^{-iE_2 t} |\nu_2\rangle + U_{e3} e^{-iE_3 t} |\nu_3\rangle \neq |\nu_e\rangle$$

at the detector there is a **probability** > 0 to see the neutrino as a ν_μ

Neutrino Oscillations are Flavor Transitions

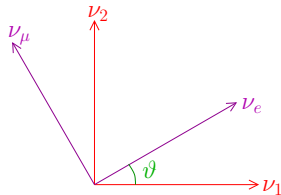
$$\nu_e \rightarrow \nu_\mu \quad \nu_e \rightarrow \nu_\tau \quad \nu_\mu \rightarrow \nu_e \quad \nu_\mu \rightarrow \nu_\tau$$

$$\bar{\nu}_e \rightarrow \bar{\nu}_\mu \quad \bar{\nu}_e \rightarrow \bar{\nu}_\tau \quad \bar{\nu}_\mu \rightarrow \bar{\nu}_e \quad \bar{\nu}_\mu \rightarrow \bar{\nu}_\tau$$

$$\Delta L_e, \Delta L_\mu, \Delta L_\tau = \pm 1 \quad \Delta L = 0$$

Two-Neutrino Mixing and Oscillations

$$|\nu_\alpha\rangle = \sum_{k=1}^2 U_{\alpha k} |\nu_k\rangle \quad (\alpha = e, \mu)$$



$$U = \begin{pmatrix} \cos\vartheta & \sin\vartheta \\ -\sin\vartheta & \cos\vartheta \end{pmatrix}$$

$$\begin{aligned} |\nu_e\rangle &= \cos\vartheta |\nu_1\rangle + \sin\vartheta |\nu_2\rangle \\ |\nu_\mu\rangle &= -\sin\vartheta |\nu_1\rangle + \cos\vartheta |\nu_2\rangle \end{aligned}$$

$$\Delta m^2 \equiv \Delta m_{21}^2 \equiv m_2^2 - m_1^2$$

Transition Probability: $P_{\nu_e \rightarrow \nu_\mu} = P_{\nu_\mu \rightarrow \nu_e} = \sin^2 2\vartheta \sin^2 \left(\frac{\Delta m^2 L}{4E} \right)$

Survival Probabilities: $P_{\nu_e \rightarrow \nu_e} = P_{\nu_\mu \rightarrow \nu_\mu} = 1 - P_{\nu_e \rightarrow \nu_\mu}$

Solar and Atmospheric Neutrino Oscillations

Solar
 $\nu_e \rightarrow \nu_\mu, \nu_\tau$

Reactor
 $\bar{\nu}_e$ disappearance

Homestake
Kamiokande
GALLEX/GNO & SAGE
Super-Kamiokande
SNO
BOREXino
(KamLAND)

\rightarrow

$\Delta m_{\text{SOL}}^2 \simeq 7.6 \times 10^{-5} \text{ eV}^2$
 $\sin^2 \vartheta_{\text{SOL}} \simeq 0.30$

Atmospheric
 $\nu_\mu \rightarrow \nu_\tau$

Accelerator
 ν_μ disappearance

Kamiokande
IMB
Super-Kamiokande
MACRO
Soudan-2
(K2K & MINOS)

\rightarrow

$\Delta m_{\text{ATM}}^2 \simeq 2.4 \times 10^{-3} \text{ eV}^2$
 $\sin^2 \vartheta_{\text{ATM}} \simeq 0.50$

Two scales of Δm^2 : $\Delta m_{\text{ATM}}^2 \simeq 30 \Delta m_{\text{SOL}}^2$

Large mixings: $\vartheta_{\text{ATM}} \simeq 45^\circ$, $\vartheta_{\text{SOL}} \simeq 33^\circ$

Three-Neutrino Mixing

$$\nu_{\alpha L} = \sum_{k=1}^3 U_{\alpha k} \nu_{kL} \quad (\alpha = e, \mu, \tau)$$

three flavor fields: ν_e, ν_μ, ν_τ

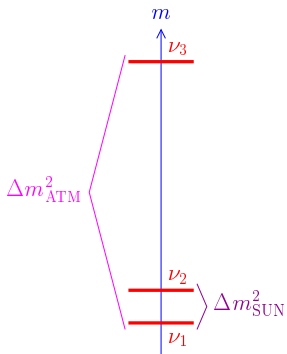
three massive fields: ν_1, ν_2, ν_3

$$\Delta m_{21}^2 + \Delta m_{32}^2 + \Delta m_{13}^2 = m_2^2 - m_1^2 + m_3^2 - m_2^2 + m_1^2 - m_3^2 = 0$$

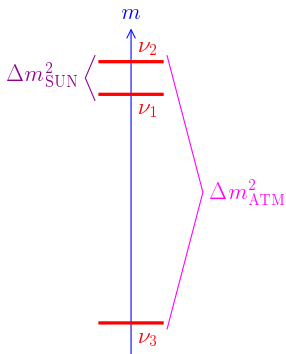
$$\Delta m_{\text{SOL}}^2 = \Delta m_{21}^2 \simeq 7.6 \times 10^{-5} \text{ eV}^2$$

$$\Delta m_{\text{ATM}}^2 \simeq |\Delta m_{31}^2| \simeq |\Delta m_{32}^2| \simeq 2.4 \times 10^{-3} \text{ eV}^2$$

Allowed Three-Neutrino Schemes



"normal"



"inverted"

different signs of $\Delta m_{31}^2 \simeq \Delta m_{32}^2$

absolute scale is not determined by neutrino oscillation data

Mixing Matrix

$$\Delta m_{21}^2 \ll |\Delta m_{31}^2|$$

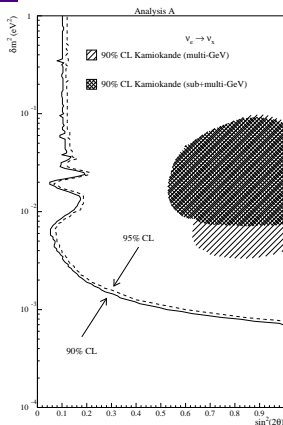
$$U = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix}$$

SOL →
↑
 ATM & LBL

$$\text{CHOOZ: } \begin{cases} \Delta m_{\text{CHOOZ}}^2 = \Delta m_{31}^2 = \Delta m_{\text{ATM}}^2 \\ \sin^2 2\vartheta_{\text{CHOOZ}} = 4|U_{e3}|^2(1 - |U_{e3}|^2) \end{cases}$$

$$|U_{e3}|^2 \lesssim 5 \times 10^{-2}$$

SOLAR AND ATMOSPHERIC ν OSCILLATIONS
ARE PRACTICALLY DECOUPLED!



[CHOOZ, PLB 466 (1999) 415]

[Palo Verde, PRD 64 (2001) 112001]

$$|U_{e1}|^2 \simeq \cos^2 \vartheta_{\text{SOL}} \quad |U_{e2}|^2 \simeq \sin^2 \vartheta_{\text{SOL}}$$

$$|U_{\mu 3}|^2 \simeq \sin^2 \vartheta_{\text{ATM}} \quad |U_{\tau 3}|^2 \simeq \cos^2 \vartheta_{\text{ATM}}$$

$$U = \underbrace{\begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix}}_{\vartheta_{23} \simeq \vartheta_{\text{ATM}}} \underbrace{\begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta_{13}} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta_{13}} & 0 & c_{13} \end{pmatrix}}_{\vartheta_{12} \simeq \vartheta_{\text{SOL}}} \underbrace{\begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}}_{\beta\beta_{0\nu}} \underbrace{\begin{pmatrix} 1 & 0 & 0 \\ 0 & e^{i\lambda_2} & 0 \\ 0 & 0 & e^{i\lambda_3} \end{pmatrix}}$$

$$\Delta m_{21}^2 = \left(7.65_{-0.20}^{+0.23}\right) \times 10^{-5} \text{ eV}^2$$

$$|\Delta m_{31}^2| = \left(2.40_{-0.11}^{+0.12}\right) \times 10^{-3} \text{ eV}^2$$

$$\sin^2 \vartheta_{12} = 0.304_{-0.016}^{+0.022}$$

$$\sin^2 \vartheta_{23} = 0.50_{-0.06}^{+0.07}$$

$$\sin^2 \vartheta_{13} < 0.035 \quad (90\% \text{ C.L.})$$

[Schwetz, Tortola, Valle, arXiv:0808.2016v3, 11 Feb 2010]

$$U \simeq \begin{pmatrix} \sqrt{2/3} & 1/\sqrt{3} & 0 \\ -1/\sqrt{6} & 1/\sqrt{3} & 1/\sqrt{2} \\ 1/\sqrt{6} & -1/\sqrt{3} & 1/\sqrt{2} \end{pmatrix}$$

Tri-Bimaximal Mixing

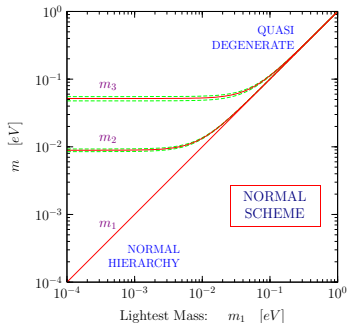
[Harrison, Perkins, Scott, PLB 530 (2002) 167]

Current Research

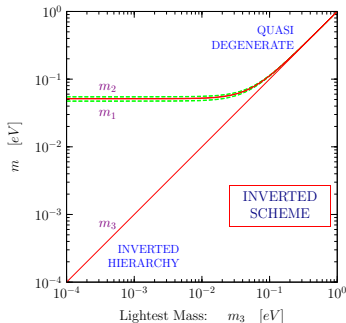
measure $\vartheta_{13} \neq 0 \implies$ CP violation, matter effects, mass hierarchy

Absolute Scale of Neutrino Masses

normal scheme



inverted scheme



$$m_2^2 = m_1^2 + \Delta m_{21}^2 = m_1^2 + \Delta m_{\text{SOL}}^2$$

$$m_3^2 = m_1^2 + \Delta m_{31}^2 = m_1^2 + \Delta m_{\text{ATM}}^2$$

$$m_1^2 = m_3^2 - \Delta m_{31}^2 = m_3^2 + \Delta m_{\text{ATM}}^2$$

$$m_2^2 = m_1^2 + \Delta m_{21}^2 \simeq m_3^2 + \Delta m_{\text{ATM}}^2$$

Quasi-Degenerate for $m_1 \simeq m_2 \simeq m_3 \simeq m_\nu \gg \sqrt{\Delta m_{\text{ATM}}^2} \simeq 5 \times 10^{-2} \text{ eV}$

▶ Tritium Beta-Decay

$$m_\nu < 2.2 \text{ eV} \quad (95\% \text{ C.L.}) \quad \text{Mainz \& Troitsk} \quad [\text{hep-ex}/0210050]$$

$$\text{KATRIN sensitivity: } m_\nu \simeq 0.2 \text{ eV} \quad [\text{hep-ex}/0109033, \text{hep-ex}/0309007]$$

▶ Neutrinoless Double-Beta Decay

$$|m_{\beta\beta}| \lesssim 0.3 - 0.7 \text{ eV} \quad (90\% \text{ C.L.}) \quad \text{CUORICINO} \quad [\text{arXiv}:1012.3266]$$

▶ Cosmology

$$m_\nu \lesssim 0.07 - 0.2 \text{ eV} \quad (95\% \text{ C.L.}) \quad [\text{hep-ph}/0805.2517, \text{arXiv}:1006.3795]$$

Anomalies Beyond 3- ν Mixing

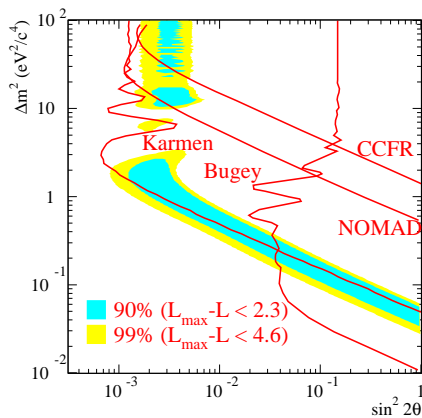
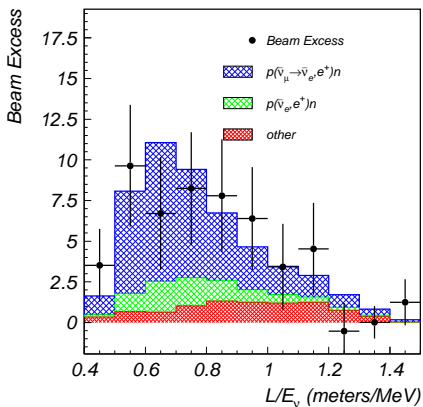
LSND

[LSND, PRL 75 (1995) 2650; PRC 54 (1996) 2685; PRL 77 (1996) 3082; PRD 64 (2001) 112007]

$$\bar{\nu}_\mu \rightarrow \bar{\nu}_e$$

$$L \simeq 30 \text{ m}$$

$$20 \text{ MeV} \leq E \leq 200 \text{ MeV}$$



$$\Delta m_{\text{LSND}}^2 \gtrsim 0.2 \text{ eV}^2 \quad (\gg \Delta m_{\text{ATM}}^2 \gg \Delta m_{\text{SOL}}^2)$$

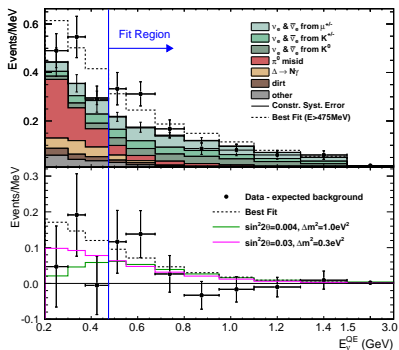
MiniBooNE Antineutrinos

[PRL 103 (2009) 111801; PRL 105 (2010) 181801]

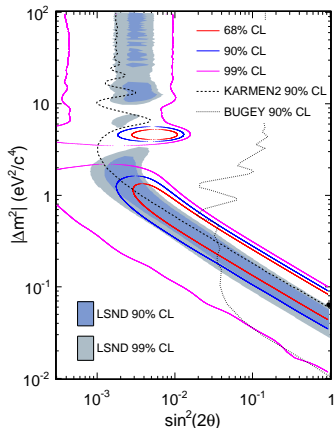
$$\bar{\nu}_\mu \rightarrow \bar{\nu}_e$$

$$L \simeq 541 \text{ m}$$

$$475 \text{ MeV} \leq E \lesssim 3 \text{ GeV}$$



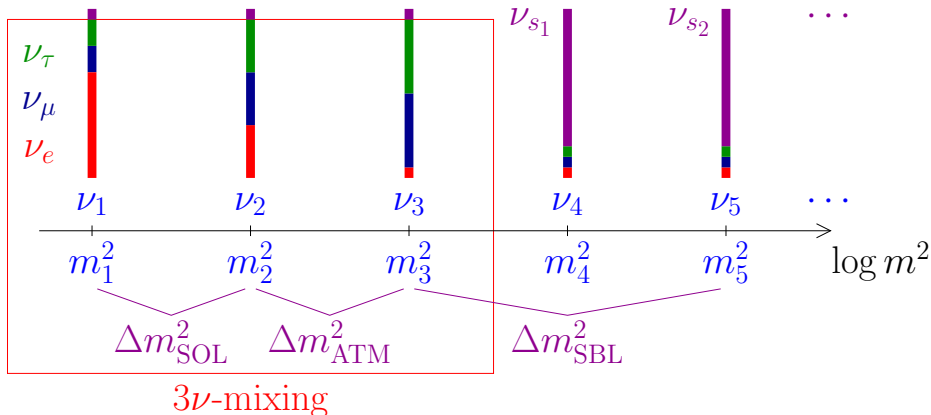
[MiniBooNE, PRL 105 (2010) 181801, arXiv:1007.1150]



Agreement with LSND $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ signal!

Similar L/E but different L and $E \implies$ Oscillations!

Beyond Three-Neutrino Mixing



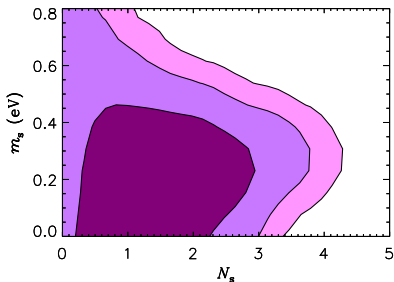
Sterile Neutrinos

- ▶ Sterile means no standard model interactions
- ▶ Active neutrinos (ν_e, ν_μ, ν_τ) can oscillate into sterile neutrinos (ν_s)
- ▶ Observables:
 - ▶ Disappearance of active neutrinos
 - ▶ Indirect evidence through combined fit of data (current indication)
- ▶ Powerful window on new physics beyond the Standard Model

Cosmology

- ▶ N_s = number of thermalized sterile neutrinos (not necessarily integer)

- ▶ CMB and LSS in Λ CDM:



[Hamann, Hannestad, Raffelt, Tamborra, Wong, PRL 105 (2010) 181301, arXiv:1006.5276]

$$N_s = 1.61 \pm 0.92$$

$$m_{\nu_s} < 0.70 \text{ eV} \quad (95\% \text{ C.L.})$$

[Giusarma, Corsi, Archidiacono, de Putter, Melchiorri, Mena, Pandolfi, arXiv:1102.4774]

- ▶ BBN:

$$N_s = 0.22 \pm 0.59$$

[Cyburt, Fields, Olive, Skillman, AP 23 (2005) 313, astro-ph/0408033]

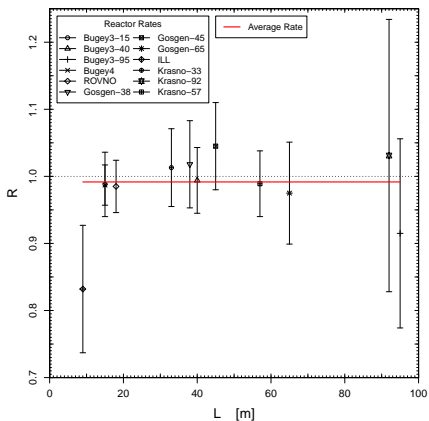
$$N_s = 0.64^{+0.40}_{-0.35}$$

[Izotov, Thuan, ApJL 710 (2010) L67, arXiv:1001.4440]

Reactor Antineutrino Anomaly

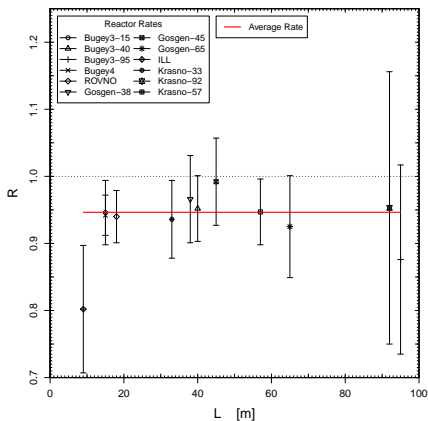
- ▶ Improved Predictions of Reactor Antineutrino Spectra, [arXiv:1101.2663](#)
- ▶ The Reactor Antineutrino Anomaly, [arXiv:1101.2755](#)

Standard Reactor $\bar{\nu}_e$ Fluxes



$$\bar{R} = 0.992 \pm 0.024$$

New Reactor $\bar{\nu}_e$ Fluxes



$$\bar{R} = 0.946 \pm 0.024$$

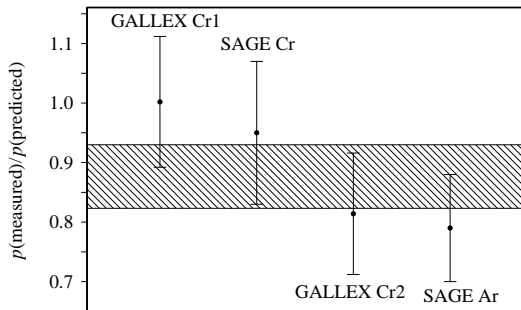
Gallium Anomaly

Gallium Radioactive Source Experiments

Tests of the solar neutrino detectors GALLEX (Cr1, Cr2) and SAGE (Cr, Ar)

Detection Process: $\nu_e + {}^{71}\text{Ga} \rightarrow {}^{71}\text{Ge} + e^-$

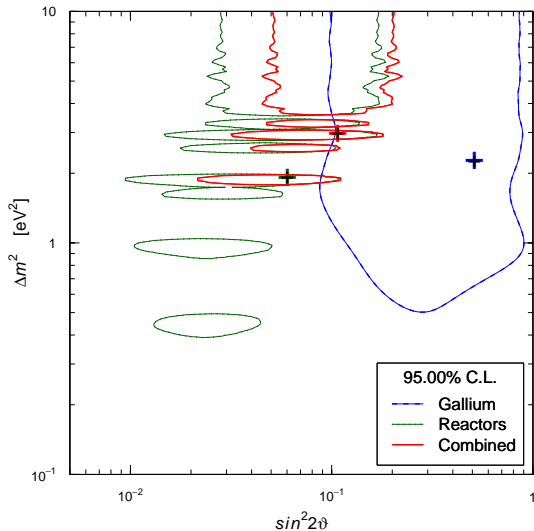
ν_e Sources: $e^- + {}^{51}\text{Cr} \rightarrow {}^{51}\text{V} + \nu_e$ $e^- + {}^{37}\text{Ar} \rightarrow {}^{37}\text{Cl} + \nu_e$



$$R_{\text{Ga}} = 0.86 \pm 0.05$$

[SAGE, PRC 73 (2006) 045805, nucl-ex/0512041]

Gallium Anomaly + Reactor Anomaly



$$\chi_{\min}^2 = 59.6$$

$$\text{NdF} = 71$$

$$\text{GoF} = 83\%$$

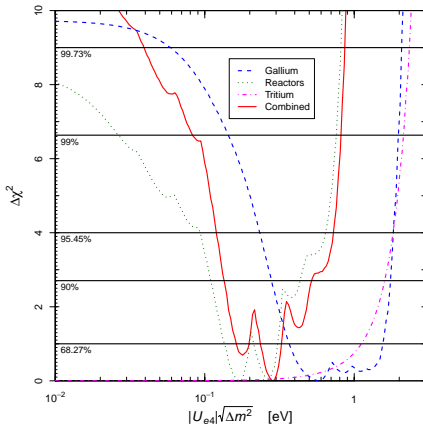
$$\sin^2 2\theta = 0.11$$

$$\Delta m^2 = 2.95 \text{ eV}^2$$

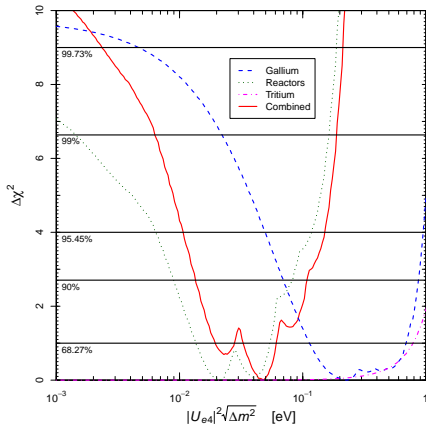
$$\text{PGoF} = 4.6\%$$

Implications of Gallium and Reactor Anomalies

β Decay



$(\beta\beta)_{0\nu}$ Decay



$$m_{\beta}^2 = \sum_k |U_{ek}|^2 m_k^2$$

$$m_{\beta\beta} = \left| \sum_k U_{ek}^2 m_k \right|$$

[Giunti, Laveder, In Preparation]

Conclusions - Three-Neutrino Mixing

$\nu_e \rightarrow \nu_\mu, \nu_\tau$ with $\Delta m_{\text{SOL}}^2 \simeq 8.3 \times 10^{-5} \text{ eV}^2$ (SOL, KamLAND)

$\nu_\mu \rightarrow \nu_\tau$ with $\Delta m_{\text{ATM}}^2 \simeq 2.4 \times 10^{-3} \text{ eV}^2$ (ATM, K2K, MINOS)



Bilarge 3ν -Mixing with $|U_{e3}|^2 \ll 1$ (CHOOZ)

β & $\beta\beta_{0\nu}$ Decay and Cosmology $\implies m_\nu \lesssim 1 \text{ eV}$

To Do

Theory: Why lepton mixing \neq quark mixing?

(Due to Majorana nature of ν 's?)

Why only $|U_{e3}|^2 \ll 1$?

Exp.: Measure $|U_{e3}| > 0 \implies$ CP viol., matter effects, mass hierarchy.

Find absolute mass scale.

Conclusions - Anomalies

- ▶ Existence of sterile neutrinos is a very interesting possibility: powerful window on new physics beyond the Standard Model
- ▶ Λ CDM cosmology and BBN hint at $N_s > 0$
- ▶ Suggestive LSND and MiniBooNE agreement on SBL $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$
- ▶ Interesting agreement of
 - ▶ Gallium Anomaly (SBL ν_e disappearance)
 - ▶ Reactor Anomaly (SBL $\bar{\nu}_e$ disappearance)

Testable Predictions:

- ▶ $m_\beta \sim 0.12 - 0.71 \text{ eV} \quad (2\sigma)$
- ▶ $m_{\beta\beta} \sim 0.011 - 0.15 \text{ eV} \quad (2\sigma)$
- ▶ Phenomenological Work in Progress: combined explanation of LSND and MiniBooNE + Gallium and Reactor Anomalies
- ▶ Experimental Work in Progress: study of new short-baseline neutrino oscillation experiments to clarify anomalies