

Closing Remarks of NuMass 2022

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NuMass 2022

Determination of the Absolute Electron (Anti)-Neutrino Mass

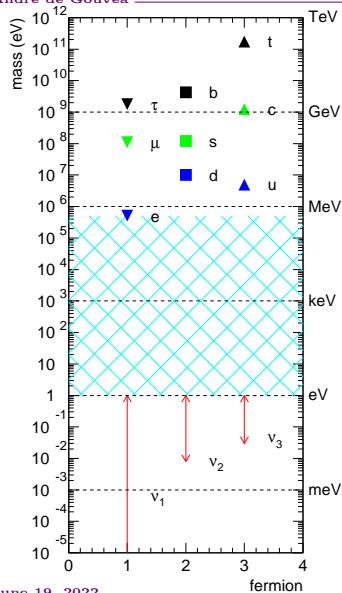
6–10 June 2022





Caveats

- ▶ This is a summary of some topics presented and discussed in the workshop.
- ▶ The selection of the topics is based on my limited understanding.
- ▶ I apologize to the many presenters of wonderful technical achievements that I am not able to report.



NEUTRINOS HAVE MASS

[albeit very tiny ones...]

So What?



NEW PHYSICS

The standard 3ν framework: parameters

Mixing matrix: **CKM** → **PMNS** (Pontecorvo-Maki-Nakagawa-Sakata)

$$U_{\alpha i} = \left[\begin{array}{ccc} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{array} \right] \left[\begin{array}{ccc} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{array} \right] \left[\begin{array}{ccc} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{array} \right] \left[\begin{array}{ccc} 1 & 0 & 0 \\ 0 & e^{i\alpha/2} & 0 \\ 0 & 0 & e^{i\beta/2} \end{array} \right]$$

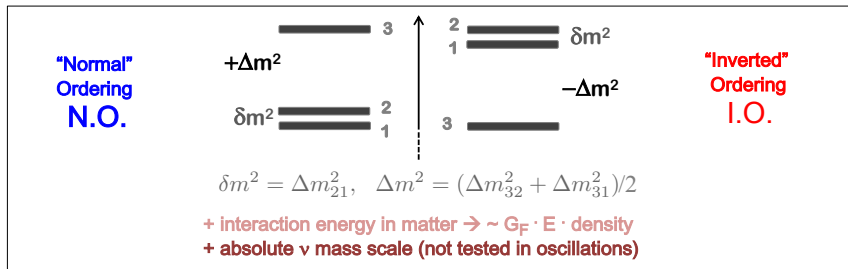
2-3 rotation
1-3 rotation
1-2 rotation
Extra CPV phases

$U(\nu) \rightarrow U^*(\bar{\nu})$
[if Majorana]

not tested in oscillat.

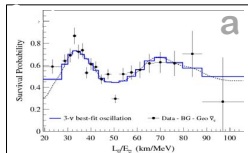
Mass [squared] spectrum

($E \sim p + m^2/2E + \text{“interaction energy”}$)

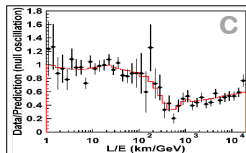


Beautiful ν oscillation data have established this 3 ν framework...

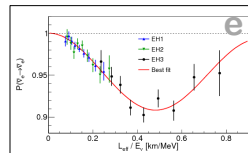
$e \rightarrow e$ (KamLAND, KL)



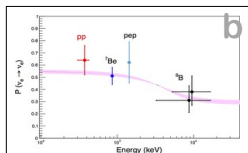
$\mu \rightarrow \mu$ (Atmospheric)



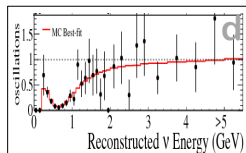
$e \rightarrow e$ (SBL React.)



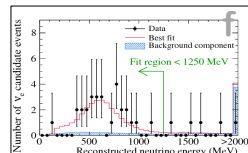
$e \rightarrow e$ (Solar)



$\mu \rightarrow \mu$ (LBL Accel)



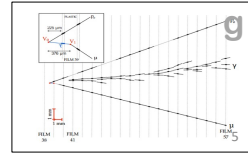
$\mu \rightarrow e$ (LBL Accel)



LBL = Long baseline (few x 100 km); SBL = short baseline (~1 km)

(a) KamLAND reactor [plot]; (b) Borexino [plot], Homestake, Super-K, SAGE, GALLEX/GNO, SNO; (c) Super-K atmosph. [plot], DeepCore, MACRO, MINOS etc.; (d) T2K [plot], NOvA, MINOS, K2K LBL accel.; (e) Daya Bay [plot], RENO, Double Chooz SBL reactor; (f) T2K [plot], MINOS, NOvA LBL accel.; (g) OPERA [plot] LBL accel., Super-K and IC-CD atmospheric.

$\mu \rightarrow \tau$ (OPERA, SK, DC)



Sketchy 3ν picture (1 significant digit)

5 knowns:

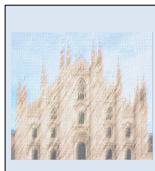
$$\delta m^2 \sim 8 \times 10^{-5} \text{ eV}^2$$

$$\Delta m^2 \sim 2 \times 10^{-3} \text{ eV}^2$$

$$\sin^2 \theta_{12} \sim 0.3$$

$$\sin^2 \theta_{23} \sim 0.5$$

$$\sin^2 \theta_{13} \sim 0.02$$



5 unknowns:

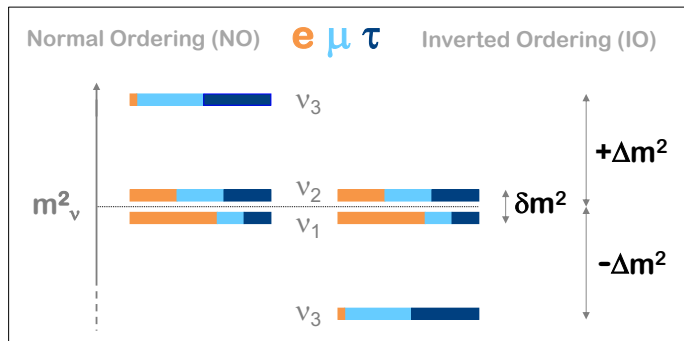
δ CPV Dirac phase

$\text{sign}(\Delta m^2) \rightarrow \text{NO/IO}$

θ_{23} octant degeneracy

absolute mass scale

Dirac/Majorana nature



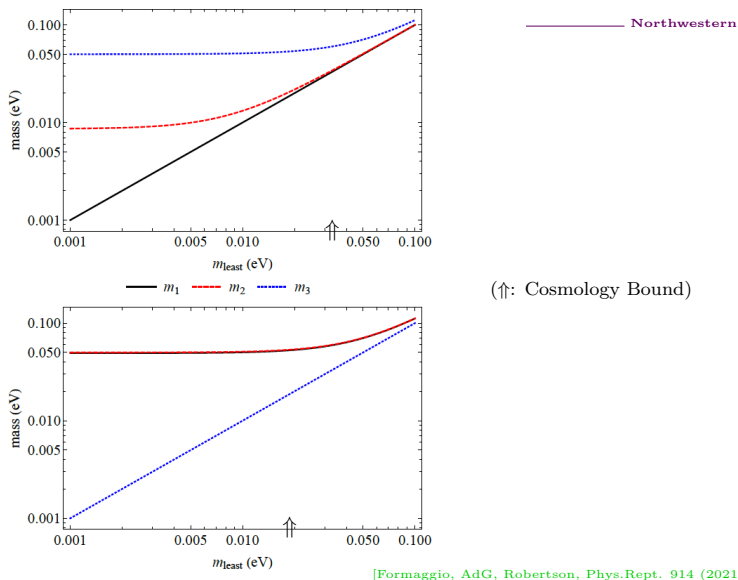


FIG. 4: Current best-fit values of the neutrino masses m_1, m_2, m_3 as a function of the lightest neutrino

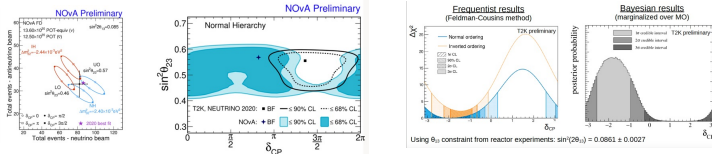
June 19, 2022

mass, for the normal mass-ordering (top) and the inverted mass ordering (bottom).

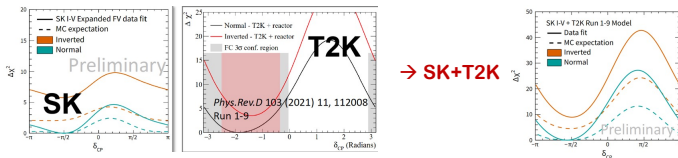
ν Mass BSM

NO/IO and δ_{CP} overview

T2K and NOvA: still in tension; but no joint analysis yet. IC-DC: IO and δ not discussed yet.



But SK atm. alone, and SK atm. +T2K \rightarrow Increased preference for NO and for $\sin\delta < 0$



\rightarrow SK+T2K

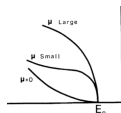
- Comment #1: Separately revised cross sections have not shed light on T2K vs NOvA tension
- Comment #2: ... but joint T2K+NOvA analysis with common interaction model still lacking
- Comment #3: SK and T2K synergy strengthens current hints on NO/IO and δ_{CP}
- Comment #4: ... but SK speaker admits that "Results from both experiments exceed sensitivity"

Absolute neutrino mass: main focus of NuMass talks!

Including the last 3ν unknowns & their observables (m_β , $m_{\beta\beta}$, Σ)

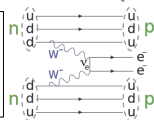
β decay, sensitive to the “effective electron neutrino mass”:

$$m_\beta = \left[c_{13}^2 c_{12}^2 m_1^2 + c_{13}^2 s_{12}^2 m_2^2 + s_{13}^2 m_3^2 \right]^{\frac{1}{2}}$$



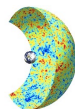
$0\nu\beta\beta$ decay: only if Majorana. “Effective Majorana mass”:

$$m_{\beta\beta} = \left| c_{13}^2 c_{12}^2 m_1 + c_{13}^2 s_{12}^2 m_2 e^{i\phi_2} + s_{13}^2 m_3 e^{i\phi_3} \right|$$



Cosmology: Dominantly sensitive to sum of neutrino masses:

$$\Sigma = m_1 + m_2 + m_3$$

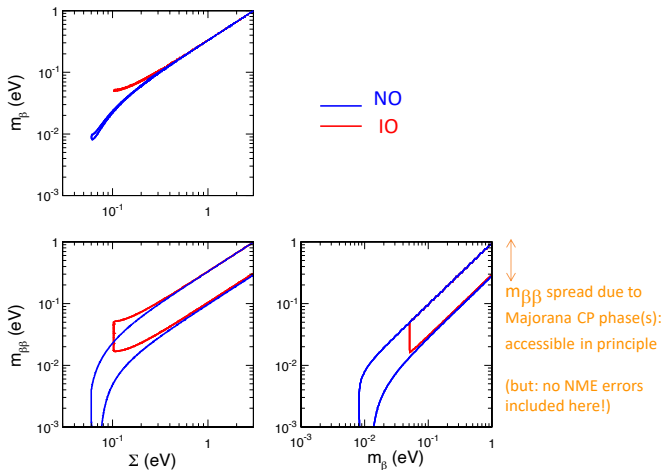


Note 1: These observables may provide handles to distinguish NO/IO.

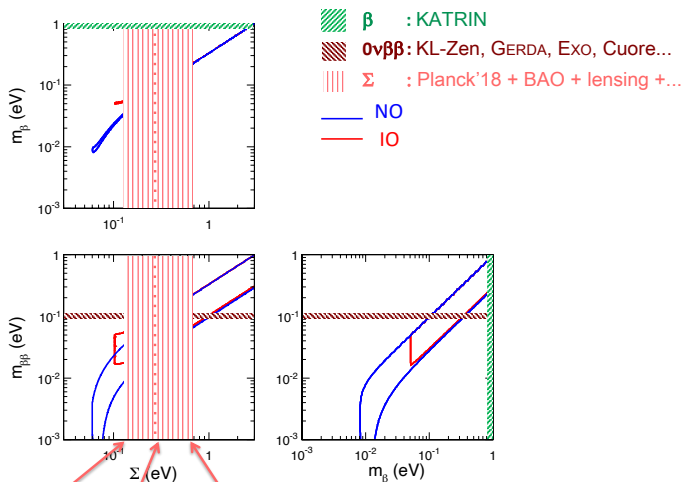
Note 2: Majorana case gives a new source of CPV (unconstrained)

Note 2: The three observables are correlated by oscillation data →

Impact of oscillations on non-oscillation parameter space (2σ)

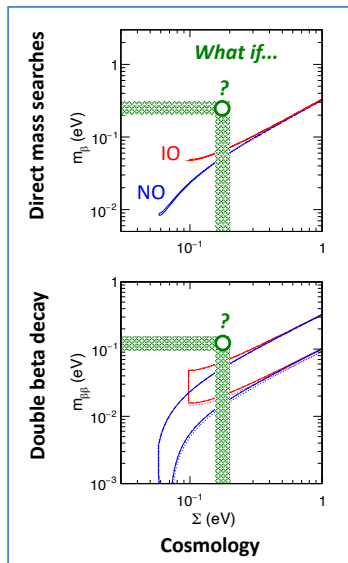


Cosmology: variety of upper bounds, with IO “under pressure”



“Aggressive” “Default” “Conservative” cosmological limits

Future data might also bring us beyond 3ν and re-shape the field...



Lack of convergence among data (barring expt mistakes) **might point towards new possibilities:**

- *Cosmology beyond Λ CDM*
 - *Alternative DBD mechanisms*
 - *New interactions (NSI)*
 - *New neutrino states ...*
- See talks by: de Gouvea, Tyagi

Main contender in current ν physics: **Light sterile ν at $O(1 \text{ eV})$ scale** but... confusing/unconfirmed hints
 → See talk by Ternes

In any case: generic expectations for new possible ν mass state(s)

Indirect probes of C_vB

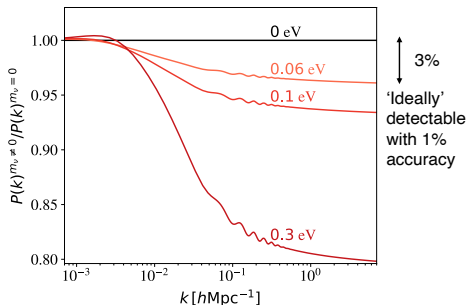
- Impact on power spectrum of matter density fluctuations

$$\delta_{\text{cdm}}^{m_\nu=0} \propto a$$

$$\delta_{\text{cdm}}^{m_\nu \neq 0} \propto a^{1-\frac{3}{5}\frac{\Omega_\nu}{\Omega_m}}$$

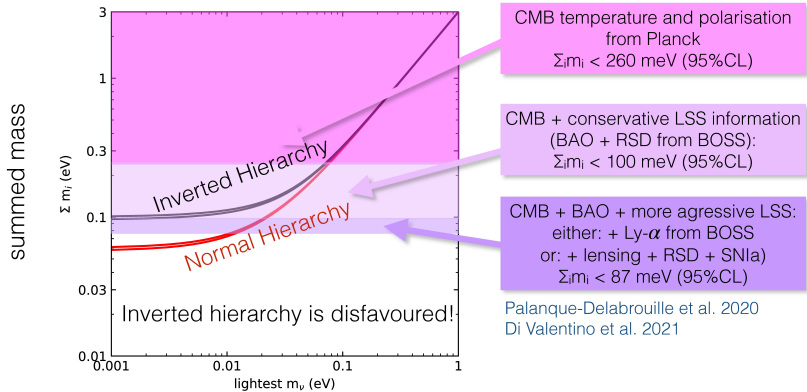
Planck TT,TE,EE + low E + lensing + BAO
 $\Sigma m_\nu < 0.12$ eV (95%cl)

Planck TT,TE,EE + low E + lensing + BAO
 +Lyman- α
 $\Sigma m_\nu < 0.09$ (95%cl) [Palanque-Delabrouille et al. 2020]



Bounds on Σm_ν

95%CL upper bounds on $\Sigma_i m_i$ for 7 parameters



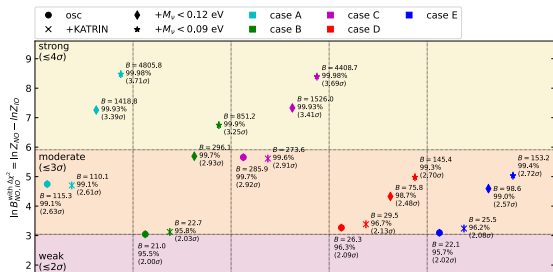
- ▶ The bound $\sum_i m_i < 0.09 \text{ eV}$ (95% C.L.) was obtained assuming $\sum_i m_i > 0$.
- ▶ An analysis adopting the oscillation prior $\sum_i m_i \gtrsim 0.05 \text{ eV}$ would return a looser bound.
- ▶ Thus, inverted hierarchy cannot be considered as disfavored at the 2σ level.

[Palanque-Delabrouille, Yeche, Schoneberg, Lesgourgues, Walther, Chabanier, Armengaud, arXiv:1911.09073]

The debate over the hierarchy

Degenerate hierarchy (DH) approximation: $m_1 = m_2 = m_3$

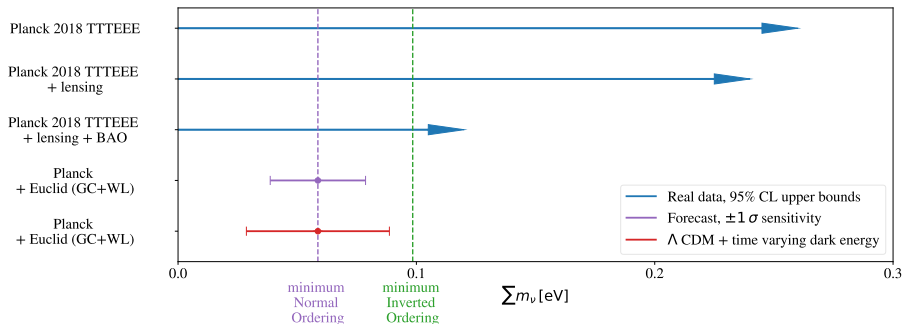
«Moderate evidence, mostly driven by neutrino oscillation data» [Gariazzo et al. 2022](#) (see also [Hergt et al. 2021](#))



case A is based on Jimenez et al.: a Gaussian prior on the logarithm of the three neutrino mass eigenstates

“The significance of the preference in favor of NO changes significantly when we consider different parameterizations.”

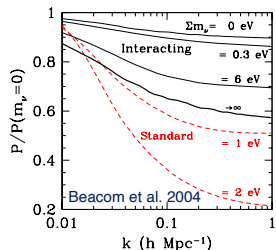
Neutrino mass constraints: the future



The nightmare scenario

Cosmological data are more and more pointing towards $\Sigma m_\nu < 0.06$ eV.

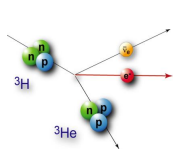
- Extended particle physics models (beyond SM)
 - Neutrino self-interactions and annihilation [“Neutrinoless Universe” Beacom et al. 2004, Esteban et al. 2021, Blinov et al. 2020, Kreisch et al. 2020, Archidiacono et al. 2020]



The interesting (exciting) scenario!

Direct Neutrino Mass Measurement with Beta Decay

Use **Kinematics only**, look at the end-point shape

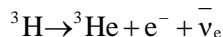


$$\frac{dN}{dE_e} = C \cdot F(E, Z) \cdot P_e \cdot (E_e + m_e c^2) \cdot (E_o - E_e) \sqrt{(E_o - E_e)^2 - m_{\nu_e}^2}$$

(some details/corrections not included)

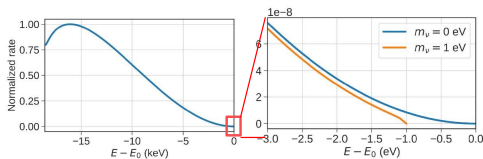
$$\sum_i |U_{ei}|^2 \cdot m_i^2 \sim m_i^2$$

in degenerated region



Tritium as beta-source

- **low end-point** (18.6 keV)
 - relatively large deformation
 - electro-statically reachable
- **short life** (12.3 y):
 - small source amount
 - less scattering in source
- super-allowed transition
 - matrix element reliably calculable
- simplest molecular:
 - molecular states calculable



only 2×10^{-13} of all beta in last 1 eV

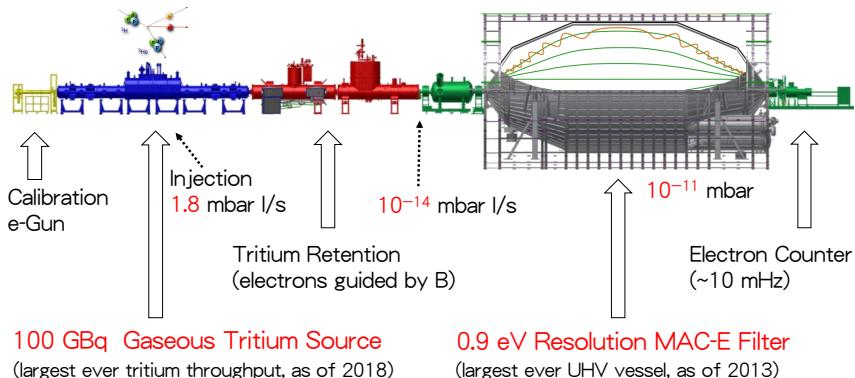
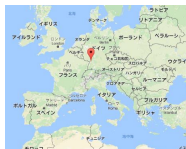
Needs:

- strong stable source
- high precision spectroscopy

KATRIN Experiment (Design)

KARlsruhe TRItium Neutrino Experiment

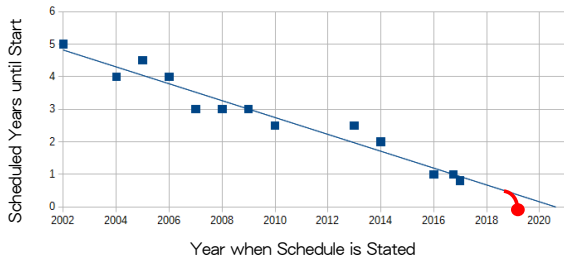
- located at Karlsruhe Institute of Technology, Karlsruhe, Germany
- design sensitivity: $m(\nu_e) < 0.2 \text{ eV}$ (90%CL, 3 years)



All numbers are from KATRIN Design Report (2004)

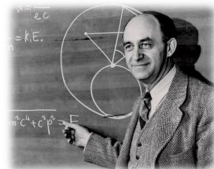
KATRIN Construction History

- First Collaboration Meeting in 2001
- Design Report in 2004, planned to start in 4 years
- Strat data taking in 2019



KATRIN has demonstrated end-points are deformed!!!

and that there can be one correct red point
in a plot with many wrong blue points



Versuch einer Theorie der β -Strahlen. I¹⁾.
Von E. Fermi in Rom.
Mit 3 Abbildungen. (Kingsgesungen am 16. Januar 1934.)

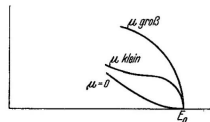
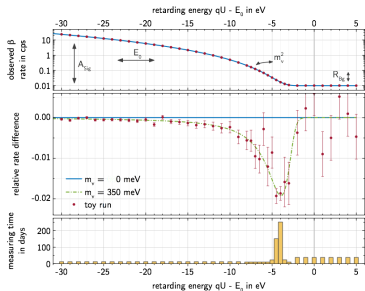


Fig. 1.

Neutrino Mass Analysis and Results

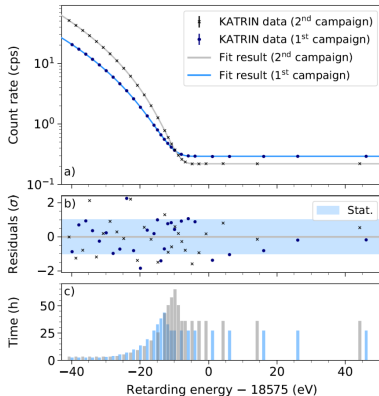
Spectrum Model and Measurement Strategy (MC made before measurements)



$$N(qU) = \underbrace{A}_{\text{activity}} \cdot \int_{qU}^{E_0} \underbrace{\frac{d\Gamma}{dE}(E; m_\nu^2, E_0)}_{\text{beta decay spectrum}} \cdot \underbrace{f(qU, E)}_{\text{apparatus response}} dE + \underbrace{B}_{\text{backgrounds}}$$

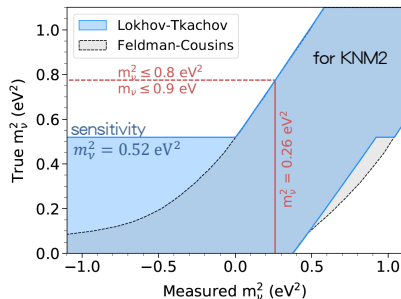
plus a number of nuisance parameters

Measured Spectra and Fitting Results



$$m_\nu^2 = 0.26 \pm 0.34 \text{ eV}^2$$

Mass Limit Setting



Feldman-Cousins (FC)

- De facto standard
- Provided only for supplementary

Lokhov-Tkachov (LT)

- No “tighter limit” from negative m^2
- For negative m^2 , stop at the sensitivity
- Coverage is still correct (no flip-flopping)

A. Lokhov is
in this meeting!

KNM 1

$$m_\nu^2 = -1.0^{+0.9}_{-1.1} \text{ eV}^2$$

$$m_\nu < 1.1 \text{ eV (LT 90\%CL)}$$

$$(m_\nu < 0.8 \text{ eV (FC 90\%CL)})$$

KNM 2

$$m_\nu^2 = 0.26 \pm 0.34 \text{ eV}^2$$

$$m_\nu < 0.9 \text{ eV (90\%CL)}$$

KNM 1 & 2 combined

$$m_\nu^2 = 0.1 \pm 0.3 \text{ eV}^2$$

$$m_\nu < 0.8 \text{ eV (90\%CL)}$$

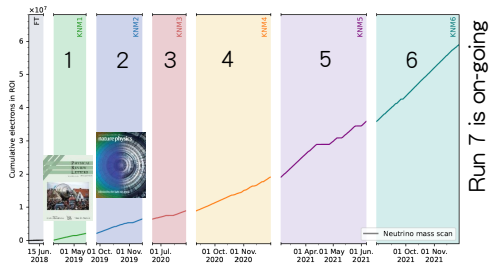
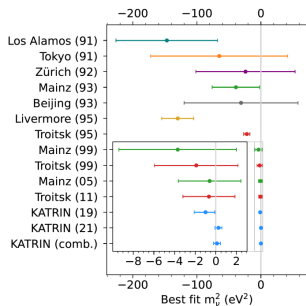
Independent Consistency Check: Beta-decay Q-value

$$\text{KATRIN: } 18575.20 \pm 0.60 \text{ eV}$$

$$\Delta m (^3\text{He} - ^3\text{H}): 18575.72 \pm 0.07 \text{ eV (this information is not used for spectrum fitting)}$$

Conclusions and Outlook

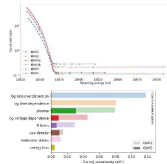
- First sub-eV result: $m_\nu^2 = 0.1 \pm 0.3 \text{ eV}^2$, $m_\nu < 0.8 \text{ eV}$ (90%CL)
- Final design sensitivity 0.2 eV (90%CL) in 5 years



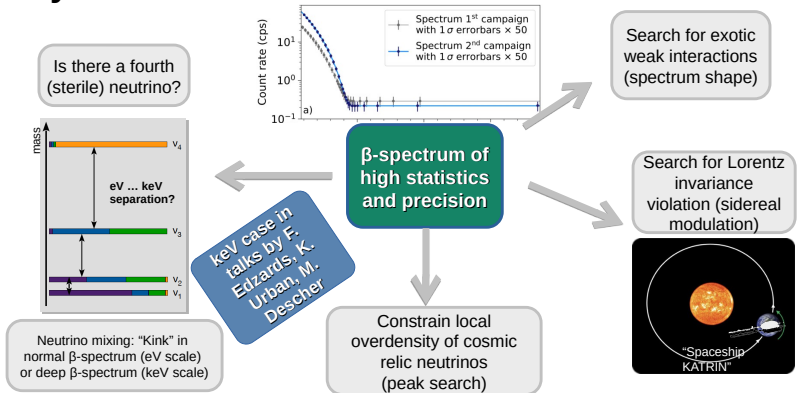
- Much more data already in hand, being analyzed (still blinded)
- After KNM 2: BG have been significantly reduced
- After KNM 2: Systematics have been significantly reduced

Preliminary estimated sensitivity up to KNM 5: $<0.5 \text{ eV}$ (90%CL)

Unblinding of KNM 3-5 is planned in this summer



“Beyond neutrino mass” in KATRIN



Conclusions

RAA mostly resolved for some flux models

The Gallium anomaly is in strong tension with the analysis of reactor rate data

No indication from ratio analyses, Neutrino-4 result doubtful

First MicroBooNE data do not confirm the MiniBooNE excess (but can not rule it out either)

No (significant) signal in atmospheric or accelerator experiments

A global 3+1 fit is statistically not acceptable

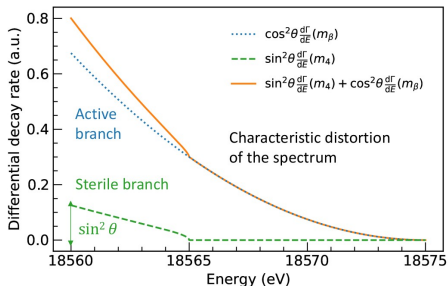
More data is needed to clarify open issues

Anomalous AND also Null results have to be checked

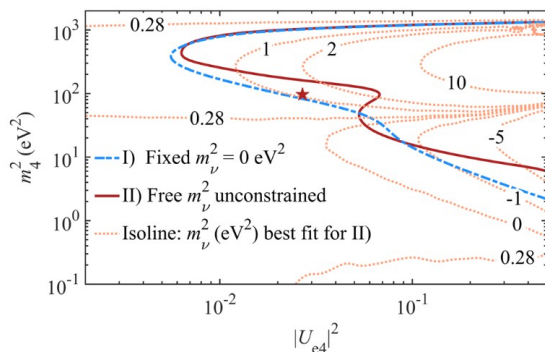
Sterile neutrinos signature in β -spectrum

- 3+1 sterile neutrino model
- Same data-set as for the neutrino mass
- Grid search in $m_4, |U_{e4}|^2$ plane

$$\frac{d\Gamma}{dE} = \underbrace{(1 - |U_{e4}|^2) \frac{d\Gamma}{dE}(m_\beta^2)}_{\text{light neutrino}} + \underbrace{|U_{e4}|^2 \frac{d\Gamma}{dE}(m_4^2)}_{\text{heavy neutrino}}$$



Interplay of m_ν^2 and m_4^2 (2nd campaign)



Fixed $m_\nu^2 = 0$

$$m_4^2 = 0.28 \text{ eV}^2, |U_{e4}| = 1.0$$

$$\Delta \chi_{null}^2 = 0.74$$

Free m_ν^2

$$m_4^2 = 98.3 \text{ eV}^2, |U_{e4}| = 0.027$$

$$\Delta \chi_{null}^2 = 2.49, m_\nu^2 = 1.1 \text{ eV}^2$$

KATRIN Collab., PRD 105, 072004 (2022)

13

08.06.2022

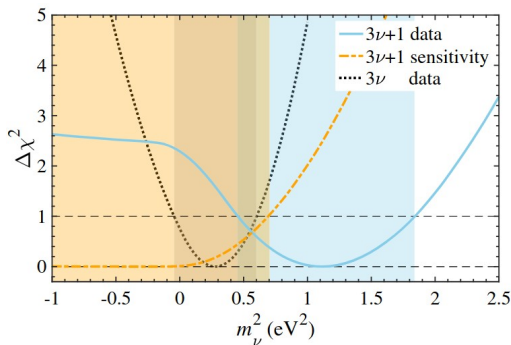
Alexey Lokhov – KATRIN beyond the neutrino mass

Institute of experimental particle physics

Please try to find the bound for free m_ν :

- ▶ In the $|U_{e4}|^2 - \Delta m_{\text{SBL}}^2$ plane, with $\Delta m_{\text{SBL}}^2 = m_4^2 - m_\nu^2$ and $m_4^2 > m_\nu^2$.
- ▶ Considering only positive m_ν^2 .

Interplay of m_ν^2 and m_4^2 (2nd campaign)



- Sizable correlation of m_ν^2 and m_4^2
 - reduction in m_ν^2 sensitivity
- Strong correlation for $m_\nu^2 < 0 \text{ eV}^2$
 - flat χ^2 profile → loss of sensitivity
 - restored by external constraints
- $m_\nu^2 > 0 \text{ eV}^2$ → x2 uncertainty on m_ν^2
- Fully restore sensitivity using
 - $|U_{e4}|^2 < 10^{-4}$

KATRIN Collab., PRD 105, 072004 (2022)

$$m_\nu \lesssim \sqrt{2.3 \text{ eV}^2} \simeq 1.5 \text{ eV} \quad (90\% \text{ C.L.}) \quad \text{in } 3\nu+1$$

about 1.7 times larger than $m_\nu \lesssim \sqrt{0.8 \text{ eV}^2} \simeq 0.9 \text{ eV} \quad (90\% \text{ C.L.}) \quad \text{in } 3\nu !$

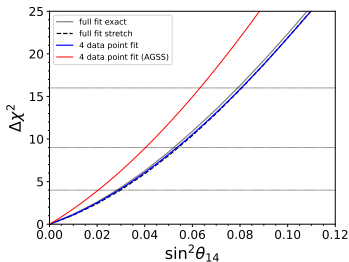
- ▶ I think that the interpretation depends on the values of $|U_{e4}|^2$ corresponding to the marginalized $\Delta\chi^2$ curve.
- ▶ The sensitivity to m_ν obviously decreases for increasing $|U_{e4}|^2$, because

$$\frac{d\Gamma}{dE} = (1 - |U_{e4}|^2) \frac{d\Gamma}{dE}(m_\nu) + |U_{e4}|^2 \frac{d\Gamma}{dE}(m_4)$$

- ▶ A reliable bound on $|U_{e4}|^2$ is the solar neutrino limit:

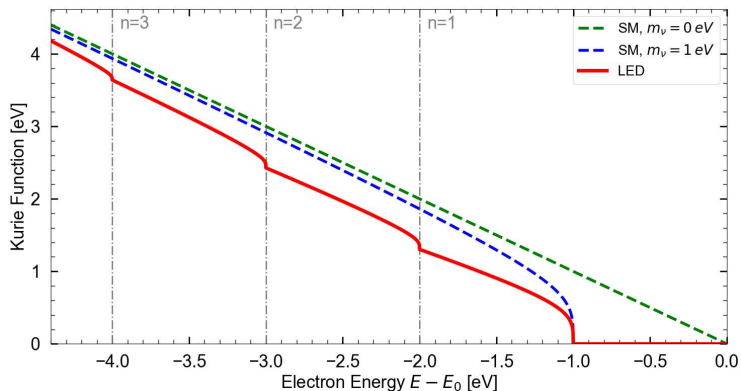
[Goldhagen, Maltoni, Reichard, Schwetz, arXiv:2109.14898]

$$\sin^2 \theta_{14} = |U_{e4}|^2$$

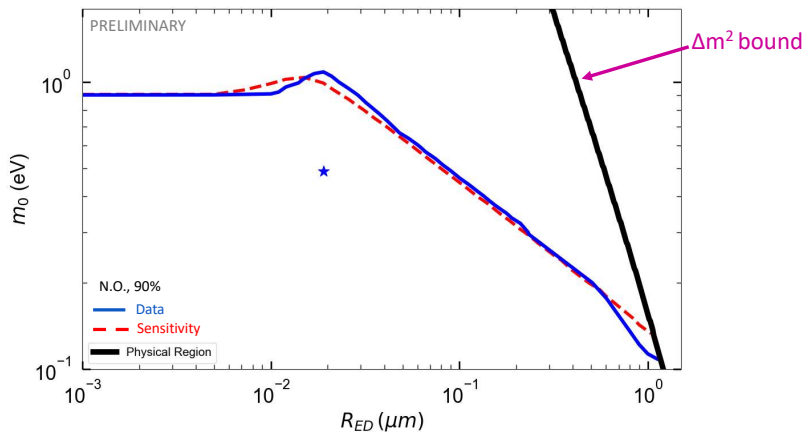


- ▶ I think that it would be useful to see:
 - ▶ The $\Delta\chi^2$ contours in the $m_\nu - |U_{e4}|^2$ plane.
 - ▶ The $\Delta\chi^2$ for m_ν with a solar $\Delta\chi^2$ pull term.

Beta-decay Experiments



Bound - KATRIN

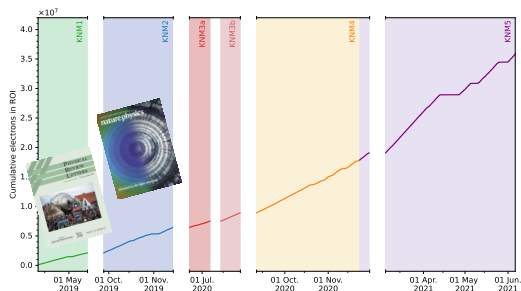


10 June 2022

NuMass 2022 Milano

20

Measurement campaigns



KATRIN Neutrino mass Measurements

	Time (hrs)	$\rho d\sigma$ (m^{-2})	Bg (mcps)
KNM1	522	1.11×10^{21}	370
KNM2	294	4.23×10^{21}	278
KNM3a	220	2.08×10^{21}	137
KNM3b	224	3.75×10^{21}	258
KNM4	1267	3.77×10^{21}	150
KNM5	1232	3.78×10^{21}	160

- Published results: KNM1 and KNM2
Phys. Rev. Lett. 123, 221802
Nat. Phys. 18, 160–166 (2022)
- Current analysis: KNM1 – KNM5
- Data-taking: KNM6, KNM7, ...

Campaigns

●●○○○

4/17 Wednesday 8th June 2022

Spectra fitting

○○

Leonard Köllenberger: (KATRIN) Details on the ν mass analysis

Unbiased analysis

○

Systematics

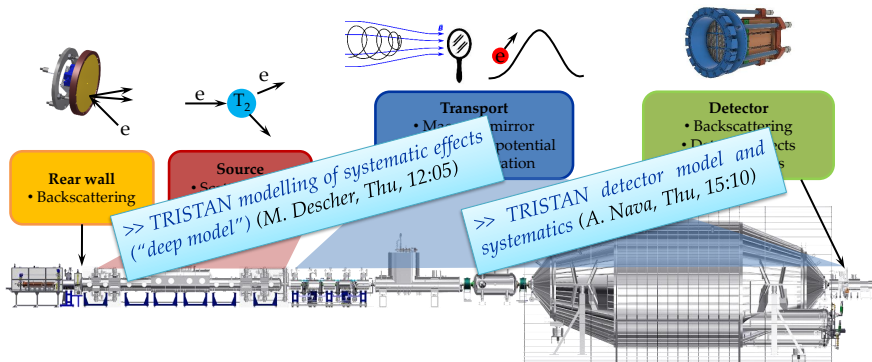
○○○○○

Outlook

○

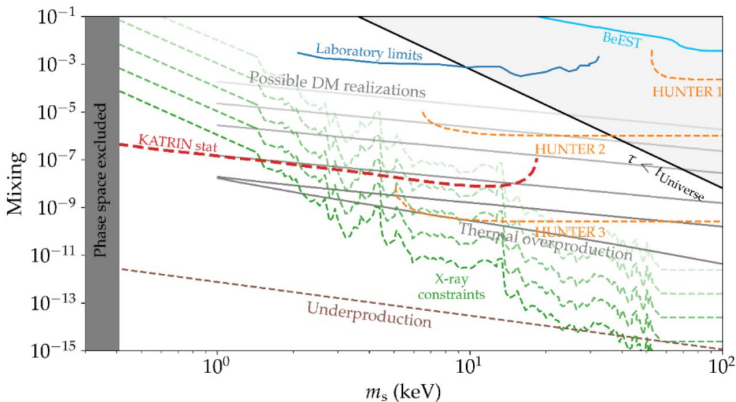
Institute for Astroparticle Physics

KATRIN - Search for keV sterile ν



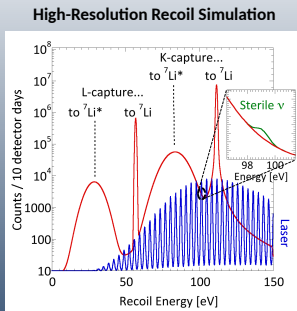
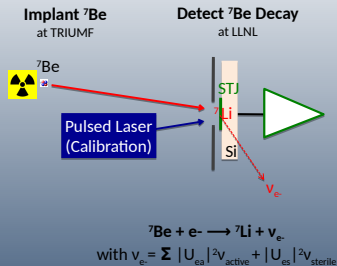
>> Precise modeling of all relevant effects required
>> Hardware modifications required

KATRIN keV sensitivity limit





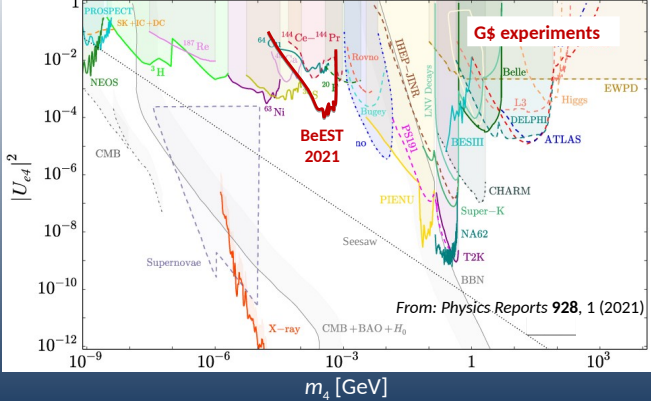
The BeEST Sterile Neutrino Experiment



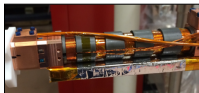
Calibrate STJ with pulsed laser.

Four peaks due to K- and L- capture into ^7Li ground and excited state

The BeEST in Context



Cyclotron Radiation Emission Spectroscopy (CRES)



PROJECT 8

Frequency Approach



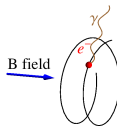
A. L. Schawlow

*"Never measure
anything but
frequency."*



O. Heaviside

*Measure the cyclotron radiation
from a single electron*



- Source transparent to microwave radiation
- No e- transport from source to detector
- Leverages precision inherent in frequency techniques

B. Monreal and JAF, Phys. Rev D80:051301

Cyclotron Radiation
Emission
Spectroscopy
(CRES)

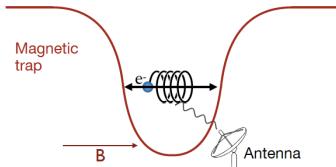


PROJECT 8

Frequency Approach



$$f_c = \frac{f_{c,0}}{\gamma} = \frac{1}{2\pi} \frac{eB}{m_e c^2 + E_{\text{kin}}}$$



$$f_{c,0} = 27.992\,491\,10(6) \text{ GHz T}^{-1}$$

- *Narrow band region of interest (@26 GHz).*
- *Small, but detectable power emitted.*

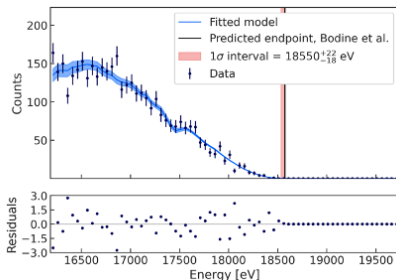
$$P(17.8 \text{ keV}, 90^\circ, 1 \text{ T}) = 1 \text{ fW}$$

$$P(30.2 \text{ keV}, 90^\circ, 1 \text{ T}) = 1.7 \text{ fW}$$

Phase II Results

Phase II CRES instrument provides 1 mm³ volume inside waveguide. Total of 3770 events observed over 3 months of data taking.

First endpoint CRES measurement conducted with no observed background in 81 days of data taking.



T₂ endpoint

Frequentist: $E_0 = (18550^{+22}_{-18})$ eV (1σ)

Bayesian: $E_0 = (18553^{+17}_{-17})$ eV (1σ)

Neutrino mass

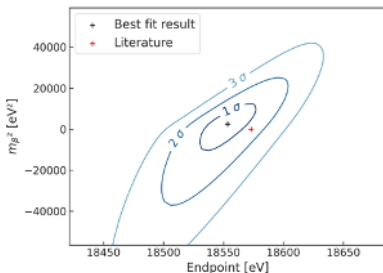
Frequentist: ≤ 178 eV/c² (90% C.L.)

Bayesian: ≤ 169 eV/c² (90% C.L.)

Background rate

$\leq 3 \times 10^{-10}$ eV⁻¹s⁻¹ (90% C.L.)

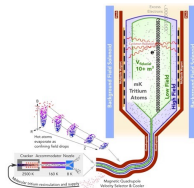
First CRES Mass Limit



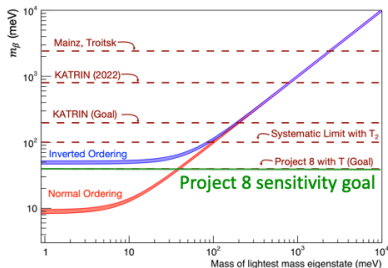
The two R&D efforts will combine into a demonstrator atomic cavity experiment at the end of Phase III, with a projected sensitivity of $400 \text{ meV}/c^2$.

A pathfinder for the final experimental goal of $40 \text{ meV}/c^2$.

CRES Phase III Pathfinder Experiment



Project 8 ν Mass Scale Sensitivity

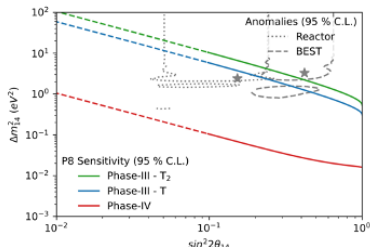


A low background, high resolution tritium experiment can help resolve “kinks” in the spectrum.

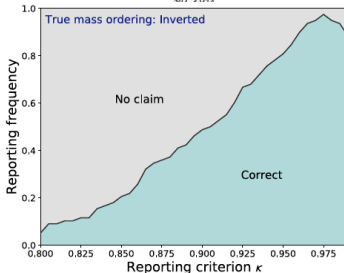
Potentially sensitive to mass splittings or low mass sterile neutrinos.

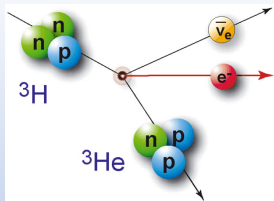
PHYSICAL REVIEW C 103, 065501 (2021)
arXiv:2203.07349v1

Project 8 Sterile Neutrino Sensitivity



Sensitivity to Neutrino Mass Ordering





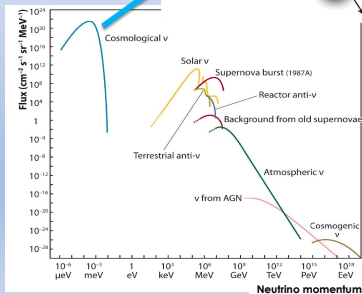
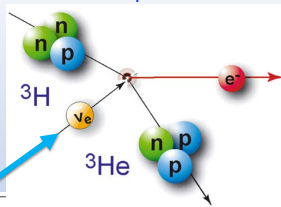
Tritium β -decay
(12.3 yr half-life)

Neutrino momentum ~ 0.17 meV

For $m_\nu = 50$ meV,
 $KE = p^2/2m$
 $= 0.17$ meV $(0.17 \text{ meV}/100 \text{ meV})$
 $= 0.3 \mu\text{eV}$

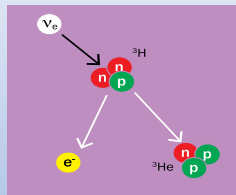
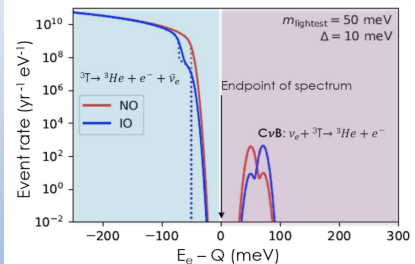
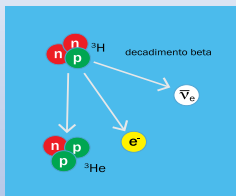
Ultra-Cold!

Neutrino capture on Tritium



Detection Concept: Neutrino Capture

- Basic concepts for relic neutrino detection were laid out in a paper by Steven Weinberg in **1962** [*Phys. Rev.* 128:3, 1457] applied for the first time to massive neutrinos in **2007** by Cocco, Mangano, Messina [[DOI: 10.1088/1475-7516/2007/06/015](https://doi.org/10.1088/1475-7516/2007/06/015)] and revisited in **2021** by Cheipesh, Cheianov, Boyarsky [<https://arxiv.org/abs/2101.10069>]



What do we know?

Gap (2m) constrained to

$m < \sim 200 \text{ meV}$

from precision cosmology

Electron flavor expected with

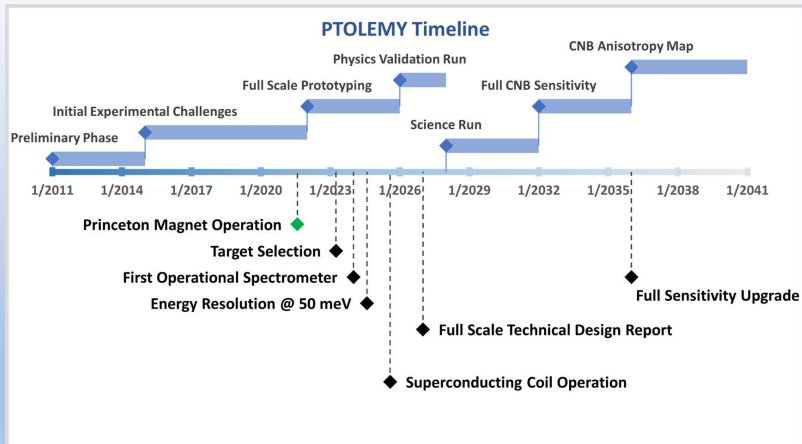
$m > \sim 50 \text{ meV}$

from neutrino oscillations

CvB Detection Requires:

few $\times 10^{-6}$ energy resolution set by m_ν
 KATRIN $\sim 10^{-4}$ (current limitation)

PTOLEMY: $10^{-4} \times 10^{-2}$
 (compact filter) \times (microcalorimeter)



Calorimetric approach as a viable alternative to spectrometers

■ **Pro:** Most of the unwanted source related effects are avoided.

■ **New way to probe sub-eV neutrino mass scale?**

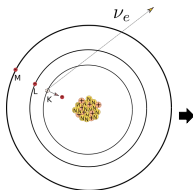
■ A good isotope should have:

Low Q value
Proximity of a peak near the ROI
Short half life to reduce the experimental challenges

Ideal calorimetric experiment

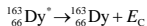
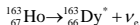
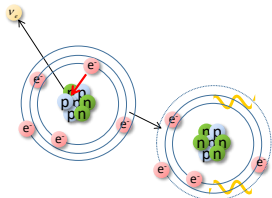
- The radioactive source is embedded in the detector(s)
- Only the neutrino energy escape detection.
- Important limits on the source intensity (statistics) that can be accumulated
- Activity also limited by the relation between energy resolution and detector size.

■ No convincing isotopes alternatives to ^3H and ^{163}Ho (yet).



Isotope	Q value [eV]	Half life [y]	Decay	B.R	Experiments
^3H	18592.01(7)	12	β^-	1	Simpson's
^{187}Re	2470.9(13)	4.3×10^{10}	β^-	1	MANU, MIBETA
^{163}Ho	2833(30)	4570	EC	1	ECho, Holmes
^{135}Cs	440	8.0×10^{11}	β^-	1.6×10^{-6}	-
^{115}In	155	4.3×10^{20}	β^-	1.1×10^{-6}	-

Electron Capture in ^{163}Ho



- $\tau_{1/2} \cong 4570$ years ($2 \cdot 10^{11}$ atoms for 1 Bq)
- $Q_{\text{EC}} = (2.833 \pm 0.030^{\text{stat}} \pm 0.015^{\text{syst}})$ keV
S. Eliseev et al., *Phys. Rev. Lett.* **115** (2015) 062501

$$Q_{\text{EC}} = m(^{163}\text{Ho}) - m(^{163}\text{Dy})$$

Penning Trap Mass Spectroscopy

@TRIGA TRAP (Uni-Mainz) (♦)

@SHIPTRAP (GSI – Darmstadt) (♦♦)

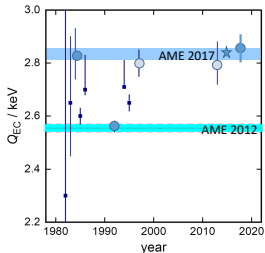
Future goal: 1 eV precision:

PENTATRAP @MPIK, Heidelberg (*)

(♦) F. Schneider et al., *Eur. Phys. J. A* **51** (2015) 89

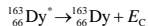
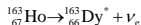
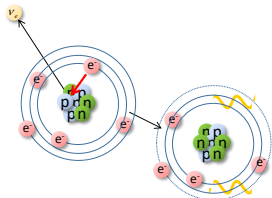
(♦♦) S. Eliseev et al., *Phys. Rev. Lett.* **115** (2015) 062501

(*) J. Repp et al., *Appl. Phys. B* **107** (2012) 983, C. Roux et al., *Appl. Phys. B* **107** (2012) 997



Electron Capture in ^{163}Ho – Spectrum

2



- $\tau_{1/2} \cong 4570$ years ($2 \cdot 10^{11}$ atoms for 1 Bq)
- $Q_{\text{EC}} = (2.833 \pm 0.030^{\text{stat}} \pm 0.015^{\text{syst}})$ keV
S. Eliseev et al., *Phys. Rev. Lett.* **115** (2015) 062501



Source = Detector

Calorimetric measurement

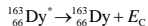
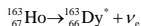
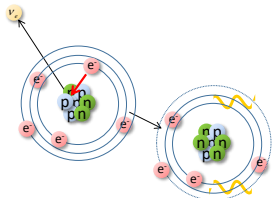
A. De Rujula and M. Lusignoli, *Phys. Lett.* **118B** (1982)

- | | |
|----------------|---|
| Advantages: | Measured neutrino complementary spectrum
No final state problems |
| Disadvantages: | Unresolved pile-up |



Electron Capture in ^{163}Ho – Spectrum

3



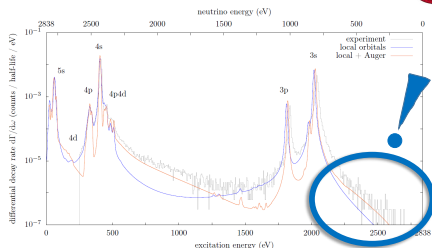
- $\tau_{1/2} \cong 4570$ years ($2 \cdot 10^{11}$ atoms for 1 Bq)
- $Q_{\text{EC}} = (2.833 \pm 0.030^{\text{stat}} \pm 0.015^{\text{syst}})$ keV
S. Eliseev et al., *Phys. Rev. Lett.* **115** (2015) 062501



Source = Detector

Calorimetric measurement

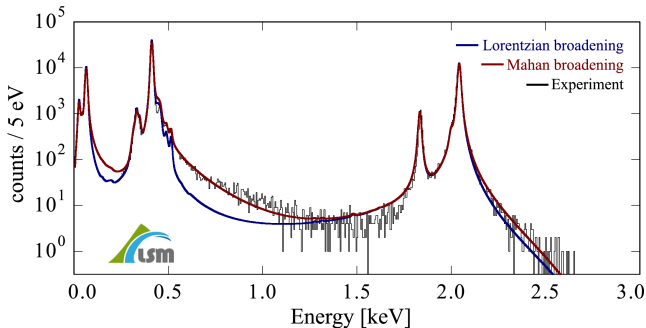
A. De Rujula and M. Lusignoli, *Phys. Lett.* **118B** (1982)



M. Braß and M. W. Haverkort, *New J. Phys.* **22** (2020) 093018

Proof of concept

20



C. Velte et al., EPJC **79** (2019) 1026

Energy resolution

$$\Delta E_{\text{FWHM}} = 9.2 \text{ eV}$$

Background level

$$b < 1.6 \times 10^{-4} \text{ events/eV/pixel/day}$$

4 day measurement with 4 pixels loaded with $\sim 0.2 \text{ Bq } ^{163}\text{Ho}$

- measurement performed underground
- test for data reduction and spectral shape analysis

- $Q_{\text{EC}} = (2838 \pm 14) \text{ eV}$
- $m(\nu_e) < 150 \text{ eV}$ (95% C.L.)

ECHo-1k high statistics spectrum

ECHo-1k chip-Au

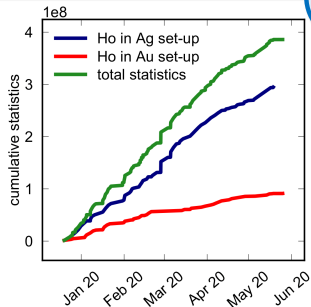
15 channels
2 temperature channels
23 pixel with implanted ^{163}Ho
3 background pixels
average activity = 0.94 Bq
total activity of 28.1 Bq

ECHo-1k chip-Ag

22 channels
2 temperature channels
34 pixel with implanted ^{163}Ho
6 background pixels
average activity = 0.71 Bq
total activity of 25.9 Bq

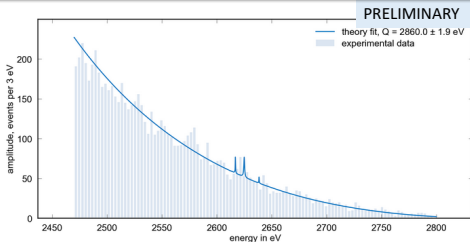
A number of ^{163}Ho events larger than 10^8 has been acquired in the first months of 2020

This statistics allow for investigating the value of the electron neutrino effective mass down to 20 eV



2.5 keV < E < 2.8 keV

26



Determination of Q_{EC} by fitting the spectrum using:

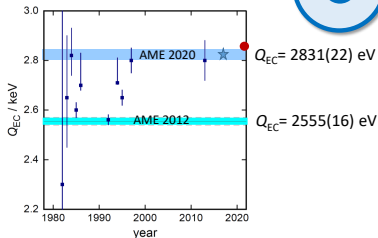
- Brass & Haverkort theory
- Flat background

$$Q_{EC} = (2860 \pm 2_{\text{stat}} \pm 5_{\text{syst}}) \text{ eV}$$



Systematic uncertainties related to theoretical spectral shape

...still too large for analysis of smaller endpoint region



$$\star Q_{EC} = (2.833 \pm 0.030^{\text{stat}} \pm 0.015^{\text{syst}}) \text{ keV}$$

S. Eliseev et al., *Phys. Rev. Lett.* **115** (2015) 062501

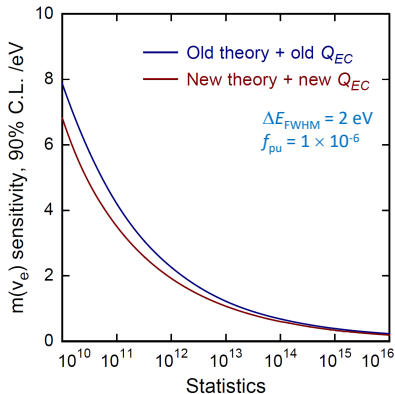
Waiting for new PENTATRAP* results

(*) J. Repp et al., *Appl. Phys. B* **107** (2012) 983

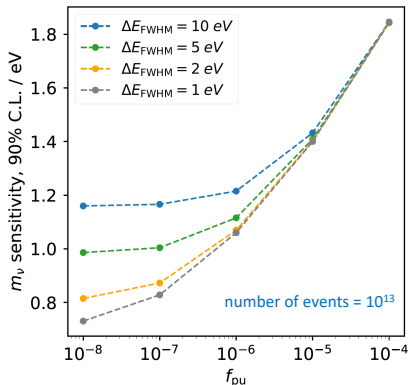
C. Roux et al., *Appl. Phys. B* **107** (2012) 997

Updated sensitivity

27

Brass & Haverkort theoretical model + new Q_{EC} -value

Sensitivity for the coming phase of ECHO



Conclusions

31

- ✓ The results obtained with ^{163}Ho loaded MMCs paved the way to large scale neutrino mass experiments based on ^{163}Ho
- ✓ The ECHo collaboration has already contributed to a [more precise description](#) of the ^{163}Ho spectrum
- ✓ A first improvement on the effective electron neutrino mass limit has been obtained in a [proof of concept](#) measurement
- ✓ More than 10^8 ^{163}Ho events have been acquired within the ECHo-1k phase → A new limit at the level of [20 eV on the effective electron neutrino mass](#) is coming soon
- ✓ Important steps [towards ECHo-100k](#) have been demonstrated: new ECHo-100k array + multiplexed readout

ECHo

... not only cool because of mK temperature

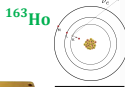


Holmes

- Holmes is an ambitious project that aims to verify the feasibility of the calorimetric approach to the neutrino mass determination.
- High performing detectors are needed, in terms of energy resolution ΔE and time resolution τ_R : LTD



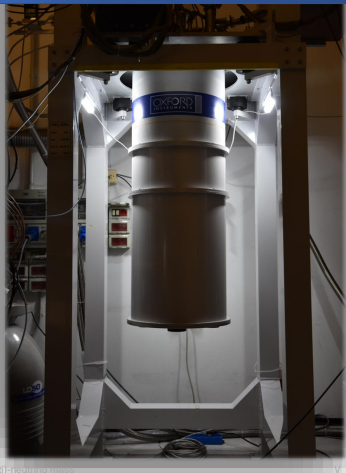
$\Delta E = \Delta E (A)$
 $\tau_R = \tau_R (A)$
(Custom ion implanter)



- Holmes has adopted a **high-risk/high-gain approach**.

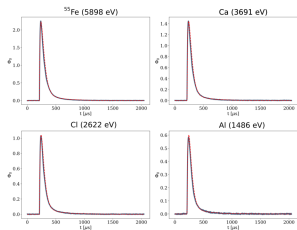
Holmes in a nutshell

- Transition Edge Sensors (TES)
 $\Delta E \approx 1 \text{ eV}$, $\tau_R < 3 \mu\text{s}$
- Microwave multiplexing readout!
- Target activity (A) of 300 Bq/det
- 6×10^5 nuclei of ^{163}Ho
- 3×10^{13} events recorded in three years
- m_{ν_e} sensitivity $O(1) \text{ eV}$

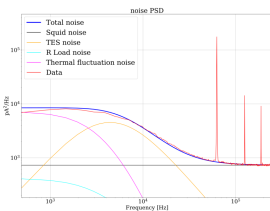


Pile-up rejection

- The pile-up fraction f_{pp} is proportional to the time resolution τ_R . The latter depends on the detector and readout characteristics and on the algorithms used to discriminate the signals.
- Requirements: high discrimination efficiency and **near zero energy dependence**.
- We have studied an application of the Wiener Filter while developing a new discrimination technique called DSVP.
- To test the algorithm, we wrote a tool to simulate the detector response (signal shape and noise spectrum).

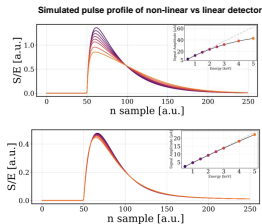


Real pulse Simulated pulse



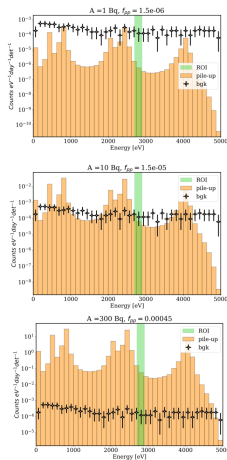
Simulation program in a nutshell:

- Goal: create pseudo-real dataset
- Energy taken from the first order Ho spectrum
- 4-th order Runge-Kutta method to solve the n differential equations
- ARMA(p,q) to properly simulate the noise spectrum



Conclusions

- Tested and tuned the final array fabrication processes.
 - | These did not spoil the detectors' performances.
- The software for analysis and signal processing of microcalorimeters events is up and running!
- The expected background contributions were assessed, both with simulations and dedicate measurements.
 - | A further reduction of a factor roughly 25% could be achieved with a similar setup studied in this work (muon veto).
 - | Pile-up reduction results equivalent to increase the measurement time by a factor 4: from 3 to 12 years.
- Ion implanter is working as expected. The production of a proper sputter target is almost ready!
- The first phase of the HOLMES experiment is expected on the last quarter of 2022: a low dose implantation of a 2×32 pixel array.
 - Influence of the Ho on the detector response will be assessed
 - A high resolution Ho calorimetric spectrum will help to discriminate between the different theoretical models
 - A first limit on the neutrino mass $O(10)$ eV will be reached
- These results will contribute to clarify if the calorimetric approach can still be considered a feasible way to reach the required sub-eV sensitivity on the neutrino mass.



Double beta decay (DBD)

Rare nuclear decay between isobars with $|\Delta Q|=2$

Even-even nuclei: favorable experimental condition

Decay modes:

- $(A,Z) \rightarrow (A,Z+2) + 2e^- + 2\bar{\nu}$ ($2\nu\beta\beta$)
 - 2nd order transition (Standard Model)
 - Very long half-lives (10^{18-24} yr)

- **$(A,Z) \rightarrow (A,Z+2) + 2e^-$ ($0\nu\beta\beta$)**

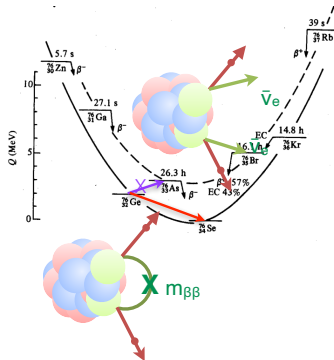
- Lepton number violation (LNV): BSM physics

Possible ONLY IF neutrinos are Majorana fermions

- LNV ($\Delta L=2$, B-L violation)
 - absolute neutrino mass scale
 - Majorana phases

SM extensions

- $(A,Z) \rightarrow (A,Z+2) + 2e^- + \text{exotics (e.g. } n\chi)$



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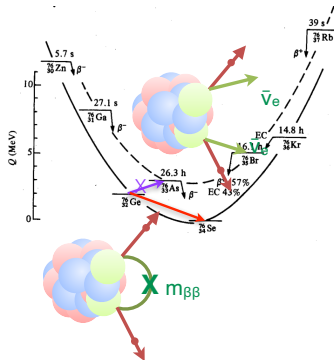
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Process description

With nontrivial approximations, it is possible to separate **atomici nuclear** and **particle** contributions and factor the transition amplitude as

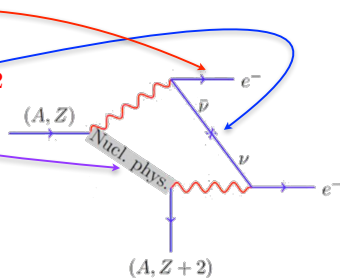
$$\Gamma^{0\nu} = G_x(Q, Z) |M_x(A, Z)|^2 |\eta_x|^2$$

$G_x(Q, Z)$ = phase space factor → calculable precisely

$M_x(A, Z)$ = nuclear matrix element (NME) → problematic

η_x = particle physics parametri → depends on model

- massive Majorana neutrinos
- GUT's
- SUSY
- ...



Process description

With nontrivial approximations, it is possible to separate **atomic nuclear** and **particle** contributions and factor the transition amplitude as

$$T_{1/2}^{-1} = G_{01} g_A^4 (M_{\text{light}}^{0\nu})^2 \frac{m_{\beta\beta}^2}{m_e^2} + \frac{m_N^2}{m_e^2} \tilde{G} \tilde{g}^4 \tilde{M}^2 \left(\frac{v}{\tilde{\Lambda}}\right)^6 + \frac{m_N^4}{m_e^2 v^2} \tilde{G}' \tilde{g}'^4 \tilde{M}'^2 \left(\frac{v}{\tilde{\Lambda}'}\right)^{10} + \dots$$

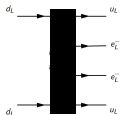
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Lepton Number Violation (LNV) and $0\nu\beta\beta$

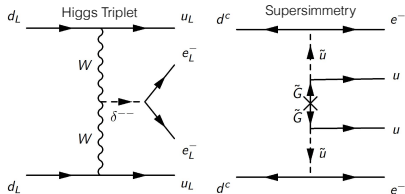
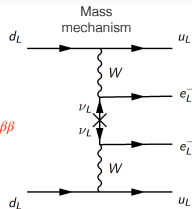


- A simple 'black box' with lots of possible modeling
- Great interest in the possibility of probing new physics

- Possibility of interference

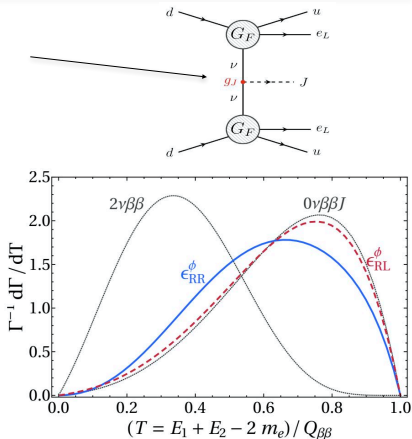
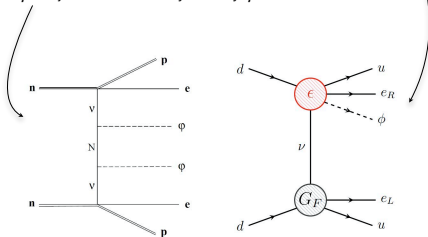
$$T_{1/2}^{-1} = G_\nu \left| \frac{m_{ee}}{m_e} M_\nu + \epsilon M_e \right|^2$$

→ measurements on different isotopes



Exotic channels

- Lepton number not necessarily violated: Dirac neutrinos
- J-type Majorons with coupling to neutrinos: $g_{ij}\bar{\nu}_i\gamma_5\nu_j J$
- Light scalars associated with Weinberg-type operators
- ϕ Majorons with LR symmetry ϕ



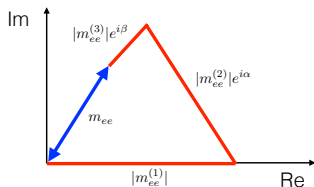
The standard interpretation: light Majorana neutrino

- $\beta\beta_{0\nu}$ is assumed to be mediated by the exchange of light Majorana neutrinos and all other mechanisms make zero or negligible contributions
- $m_{\beta\beta}$ depends on a total of seven parameters (θ_{12} , θ_{13} , three masses and two Majorana phases)
- neutrino oscillation experiments are sensitive only to the two mixing angles, two neutrino mass squared differences, and the mass ordering
 \Rightarrow only four out of seven degrees of freedom, can be bounded

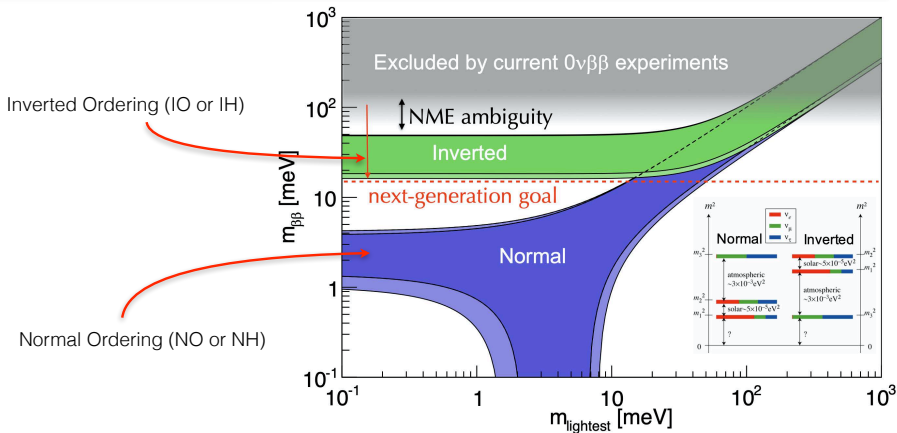
$$\eta_x = \langle m_{ee} \rangle = \sum_k U_{ek}^2 m_k$$

$$= c_{12}^2 c_{13}^2 m_1 + s_{12}^2 c_{13}^2 e^{i\alpha} m_2 + s_{13}^2 e^{i\beta} m_3$$

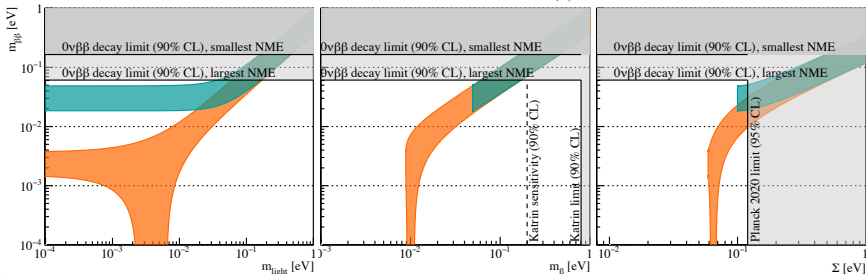
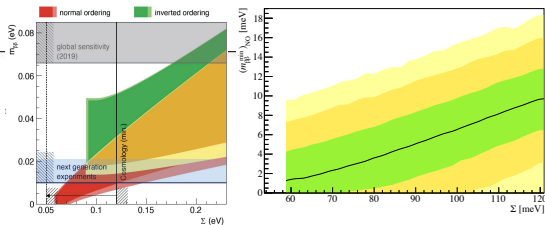
- Transition amplitude is proportional to the coherent sum of neutrino masses
- Majorana phases play a crucial role: possible cancellations



Light Majorana neutrinos



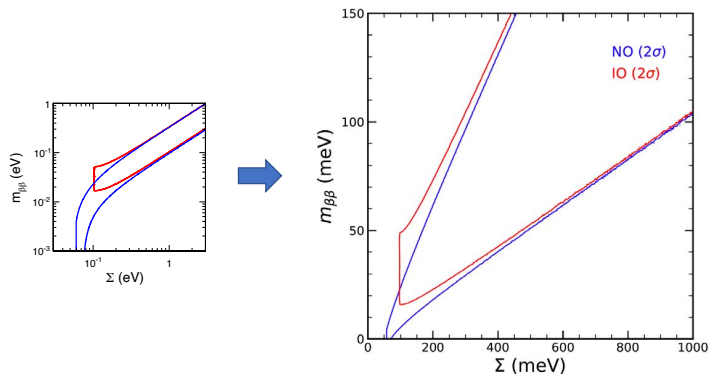
$m_{\beta\beta}$ and the m_ν quest



(1) Time to shift from log to linear scales!

$$m_{\beta\beta} = |U_{e1}^2 m_1 + U_{e2}^2 m_2 + U_{e3}^2 m_3|$$

$$\Sigma = m_1 + m_2 + m_3$$



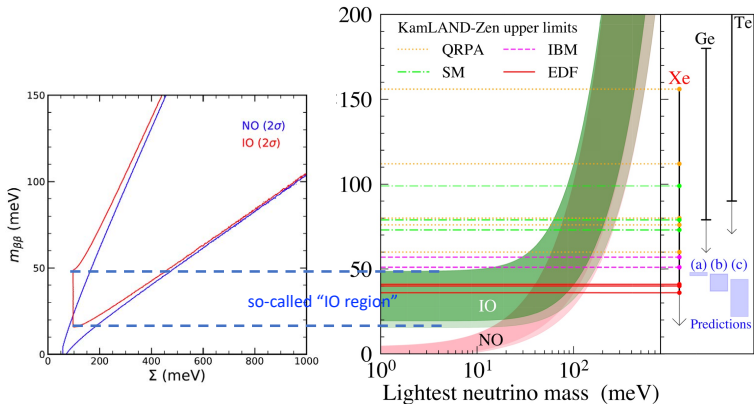
Regions allowed by current oscillation data for NO and IO separately.

[Width of regions dominated by unknown Majorana phases]

Similarly for $m_{\beta\beta}$ vs Σ , without Majorana phase uncertainties

(2) Language bonus: avoid log-inspired misleading jargon

The (improperly named) "IO region" is compatible with NO as well...

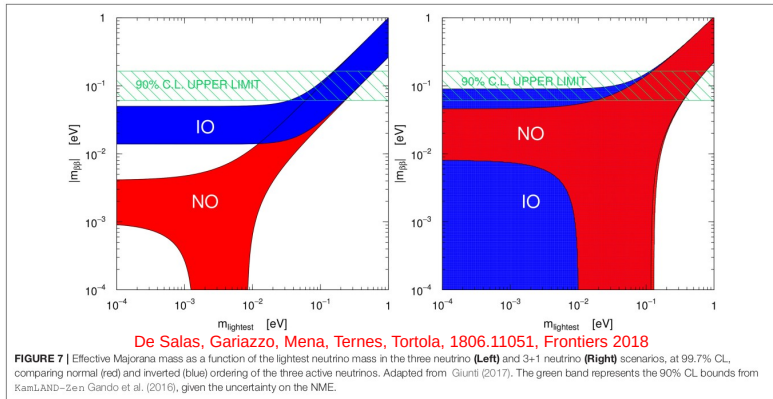


Important to note that Σ is observable while m_{lightest} is not.
Elongated stripes in $\log(m_{\text{lightest}})$ have no real significance.

Stat. bonus: If Bayesian results depend sensitively on log vs linear priors etc. \rightarrow unreliable, IMHO

Neutrinoless $\beta\beta$ decay

More details in talk by
Oliviero Cremonesi!



$m_{\beta\beta}$ extraction from data: a hard task

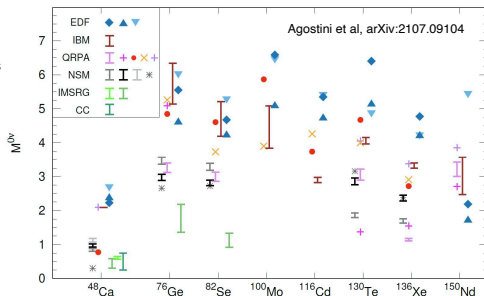
$M^{0\nu}$ values estimated from different calculation methods

$$\Gamma^{0\nu} = G_{0\nu}(Q, Z) |M_{0\nu}(A, Z)|^2 |m_{ee}|^2$$

- Factor 2-3 "uncertainty" among nuclear models.
- 'Ab initio' calculations for lighter isotopes give lower values

NME calculations:

- EDF: large
- QRPA: larger spread
- NSM: small
- IMSRG ^{48}Ca ab initio: very small

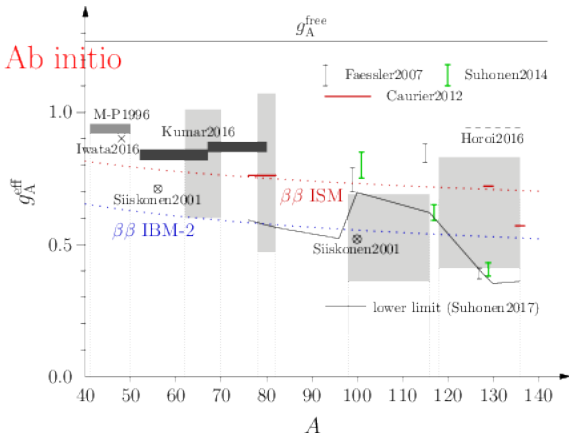


Much more problematic:
nuclear matrix elements

- Significant differences
- Results inaccuracy

g_A quenching?

Results extracted from the GT β^\pm /EC and $2\nu\beta\beta$ calculations



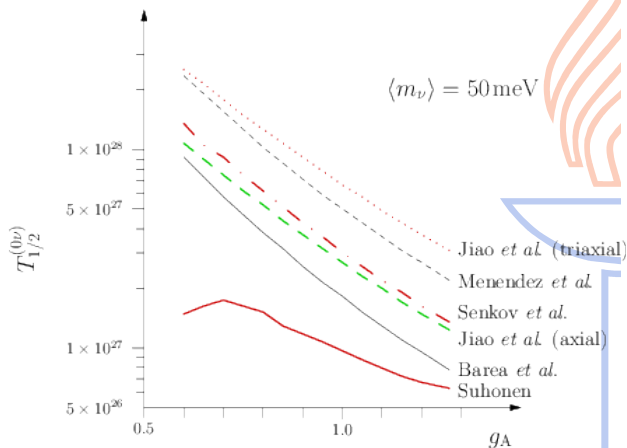
Ab initio

Ab initio: P. Gysbers *et al.*, Nature Physics 15 (2019) 428

- Faessler2007: **pnQRPA** A. Faessler *et al.*, arXiv 0711.3996v1 [Nucl-th]
- Suhonen2014: **pnQRPA** J. Suhonen *et al.*, Nucl. Phys. A 924 (2014) 1
- Suhonen2017: **pnQRPA** J. Suhonen, Phys. Rev. C 96 (2017) 055501
- Caurier2012: **ISM** E. Caurier *et al.*, Phys. Lett. B 711 (2012) 62
- Horoi2016: **ISM** M. Horoi *et al.*, Phys. Rev. C 93 (2016) 024308
- M-P1996: **ISM** G. Martínez-Pinedo *et al.*, Phys. Rev. C 53 (1996) R2602
- Iwata2016: **ISM** Y. Iwata *et al.*, Phys. Rev. Lett. 116 (2016) 112502
- Kumar2016: **ISM** V. Kumar *et al.*, J. Phys. G 43 (2016) 105104 Phys. Lett. B 711 (2012) 62
- Siiskonen2001: **ISM** T. Siiskonen *et al.*, Phys. Rev. C 63 (2001) 055501
- **beta beta ISM and IBM-2**: J. Barea *et al.*, Phys. Rev. C 87 (2013) 014315
- Light hatched regions: **pnQRPA** H. Ejiri *et al.*, J. Phys. G 42 (2015) 055201 ; P. Pirinen *et al.*, Phys. Rev. C 91 (2015) 054309 ; F. Deppisch *et al.*, Phys. Rev. C 94 (2016) 055501

Example: $0\nu\beta\beta$ NMEs of ^{76}Ge , effect on the half-life

- **Jiao *et al.*:** Phys. Rev. C 96 (2017) 054310 (GCM+ISM)
- **Menendez *et al.*:** Nucl. Phys. A 818 (2009) 139 (ISM)
- **Senkov *et al.*:** Phys. Rev. C 93 (2016) 044334 (ISM)
- **Barea *et al.*:** Phys. Rev. C 91 (2015) 034304 (IBM-2)
- **Suhonen:** Phys. Rev. C 96 (2017) 055501 (pnQRPA + g_{pp} + isospin restoration + **data on $2\nu\beta\beta$**)



How do we extract information on the value of g_A ?

These methods are now available:

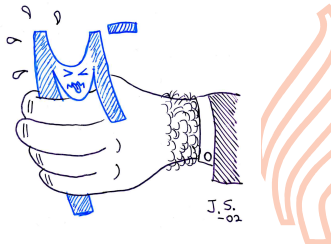
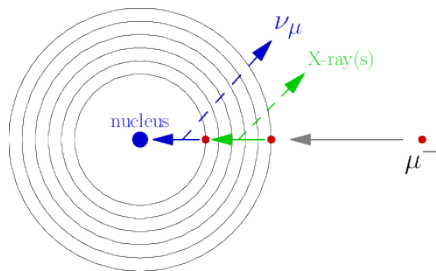
For low momentum exchanges (g_A):

- study half-lives of β decays (1^+ and 2^- states)
- study half-lives of $2\nu\beta\beta$ decays (1^+ states)
- Study electron spectral shapes of β decays (J^π states)

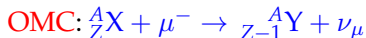
For high momentum exchanges like $0\nu\beta\beta$ decay ($g_{A,0\nu}$):

- Study nuclear muon capture (J^π states)

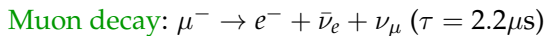
Ordinary Muon Capture (OMC)



Nuclear muon capture:



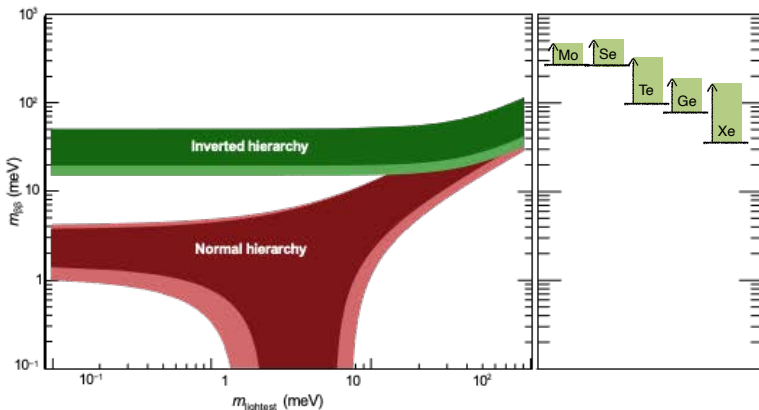
Also:



OMC probability $\sim Z^4$
(in Fe 91% are captured,
breakeven at $Z \sim 11$)

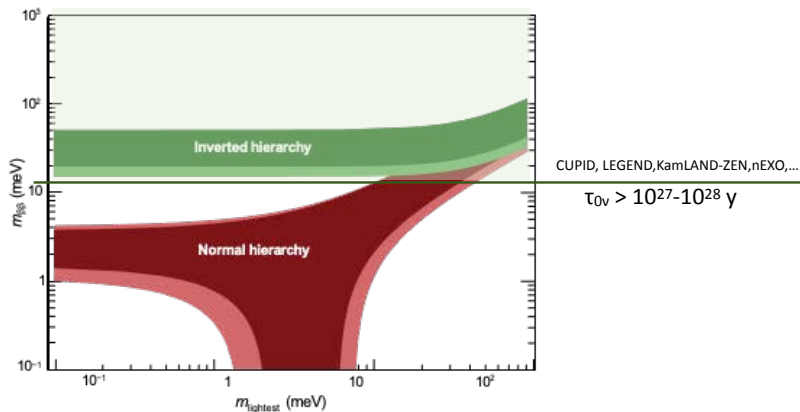
Current results: $m_{\beta\beta}$

no g_A quenching



Future sensitivities

no g_A quenching



Conclusions

- ▶ The future is bright for:
 - ▶ Neutrino Physics.
 - ▶ The Determination of the Absolute Electron (Anti)-Neutrino Mass.
 - ▶ The Detection of the Cosmic Neutrino Background.
 - ▶ Finding Lepton Number Violation and Majorana Neutrino Masses.
 - ▶ The NuMass Workshop Series.
- ▶ Thanks to Angelo and all the local organizers for the perfect organization!
- ▶ See you at NuMass 2024!