

# Closing Remarks of NuMass 2022

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NuMass 2022

Determination of the Absolute Electron (Anti)-Neutrino Mass

6–10 June 2022





## Caveats

- ▶ This is a summary of some topics presented and discussed in the workshop.
- ▶ The selection of the topics is based on my limited understanding.
- ▶ I apologize to the many presenters of wonderful technical achievements that I am not able to report.

## The standard 3ν framework: parameters

Mixing matrix: **CKM** → **PMNS** (Pontecorvo-Maki-Nakagawa-Sakata)

$$U_{\alpha i} = \left[ \begin{array}{ccc} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{array} \right] \left[ \begin{array}{ccc} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{array} \right] \left[ \begin{array}{ccc} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{array} \right] \left[ \begin{array}{ccc} 1 & 0 & 0 \\ 0 & e^{i\alpha/2} & 0 \\ 0 & 0 & e^{i\beta/2} \end{array} \right]$$

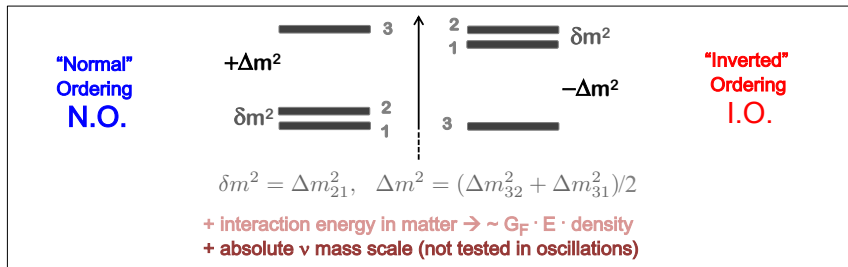
2-3 rotation
1-3 rotation
1-2 rotation
Extra CPV phases

+ CPV "Dirac" phase
[if Majorana]
not tested in oscillat.

$U(\nu) \rightarrow U^*(\bar{\nu})$

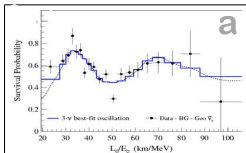
Mass [squared] spectrum

( $E \sim p + m^2/2E + \text{"interaction energy"}$ )

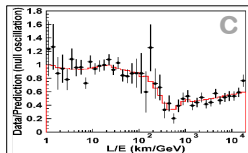


Beautiful  $\nu$  oscillation data have established this 3 $\nu$  framework...

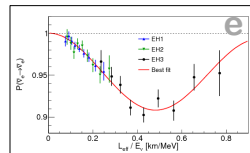
$e \rightarrow e$  (KamLAND, KL)



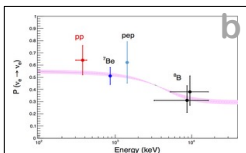
$\mu \rightarrow \mu$  (Atmospheric)



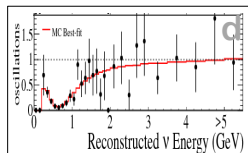
$e \rightarrow e$  (SBL React.)



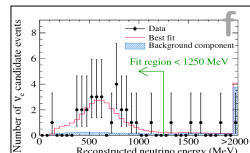
$e \rightarrow e$  (Solar)



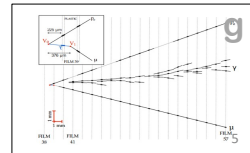
$\mu \rightarrow \mu$  (LBL Accel)



$\mu \rightarrow e$  (LBL Accel)



$\mu \rightarrow \tau$  (OPERA, SK, DC)



LBL = Long baseline (few x 100 km); SBL = short baseline (~1 km)

(a) KamLAND reactor [plot]; (b) Borexino [plot], Homestake, Super-K, SAGE, GALLEX/GNO, SNO; (c) Super-K atmosph. [plot], DeepCore, MACRO, MINOS etc.; (d) T2K [plot], NOvA, MINOS, K2K LBL accel.; (e) Daya Bay [plot], RENO, Double Chooz SBL reactor; (f) T2K [plot], MINOS, NOvA LBL accel.; (g) OPERA [plot] LBL accel., Super-K and IC-CD atmospheric.

### Sketchy 3ν picture (1 significant digit)

#### 5 knowns:

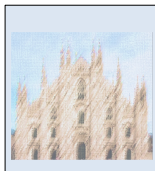
$$\delta m^2 \sim 8 \times 10^{-5} \text{ eV}^2$$

$$\Delta m^2 \sim 2 \times 10^{-3} \text{ eV}^2$$

$$\sin^2 \theta_{12} \sim 0.3$$

$$\sin^2 \theta_{23} \sim 0.5$$

$$\sin^2 \theta_{13} \sim 0.02$$



#### 5 unknowns:

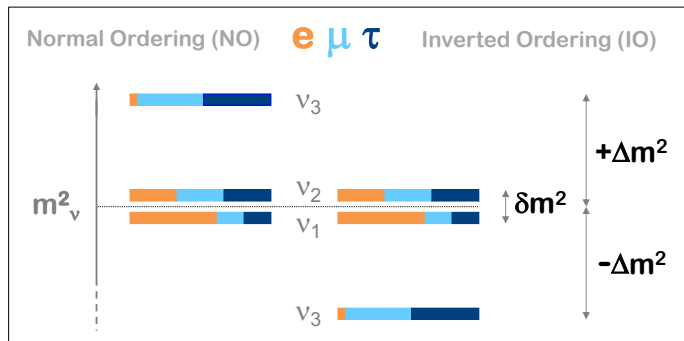
$\delta$  CPV Dirac phase

$\text{sign}(\Delta m^2) \rightarrow \text{NO/IO}$

$\theta_{23}$  octant degeneracy

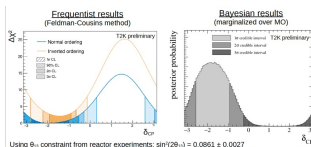
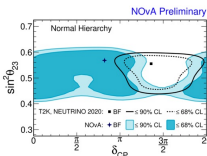
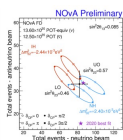
absolute mass scale

Dirac/Majorana nature

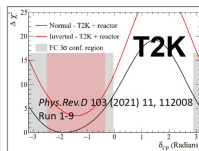
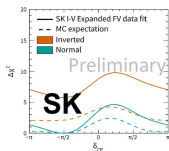


## NO/IO and $\delta_{CP}$ overview

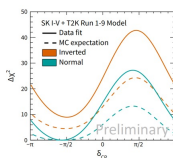
T2K and NOvA: still in tension; but no joint analysis yet. IC-DC: IO and  $\delta$  not discussed yet.



But SK atm. alone, and SK atm. +T2K → Increased preference for NO and for  $\sin\delta < 0$



→ SK+T2K



Comment #1: Separately revised cross sections have not shed light on T2K vs NOvA tension

Comment #2: ... but joint T2K+NOvA analysis with common interaction model still lacking

Comment #3: SK and T2K synergy strengthens current hints on NO/IO and  $\delta_{CP}$

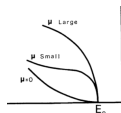
Comment #4: ... but SK speaker admits that "Results from both experiments exceed sensitivity"

## Absolute neutrino mass: main focus of NuMass talks!

Including the last  $3\nu$  unknowns & their observables ( $m_\beta, m_{\beta\beta}, \Sigma$ )

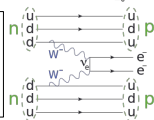
**$\beta$  decay**, sensitive to the “effective electron neutrino mass”:

$$m_\beta = \left[ c_{13}^2 c_{12}^2 m_1^2 + c_{13}^2 s_{12}^2 m_2^2 + s_{13}^2 m_3^2 \right]^{\frac{1}{2}}$$



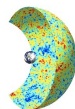
**$0\nu\beta\beta$  decay**: only if Majorana. “Effective Majorana mass”:

$$m_{\beta\beta} = \left| c_{13}^2 c_{12}^2 m_1 + c_{13}^2 s_{12}^2 m_2 e^{i\phi_2} + s_{13}^2 m_3 e^{i\phi_3} \right|$$



**Cosmology**: Dominantly sensitive to sum of neutrino masses:

$$\Sigma = m_1 + m_2 + m_3$$



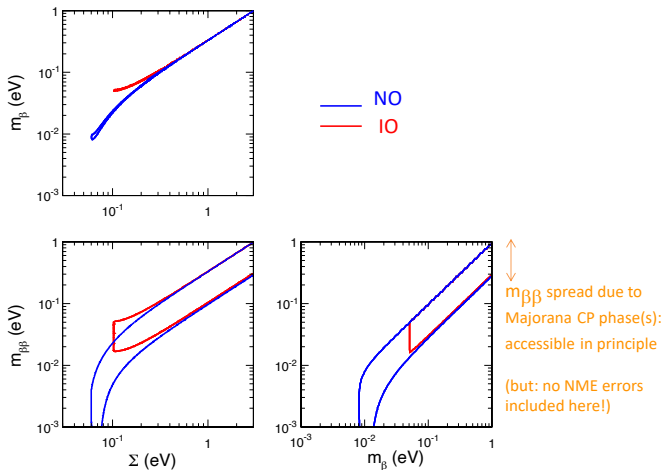
Note 1: These observables may provide handles to distinguish NO/IO.

Note 2: Majorana case gives a new source of CPV (unconstrained)

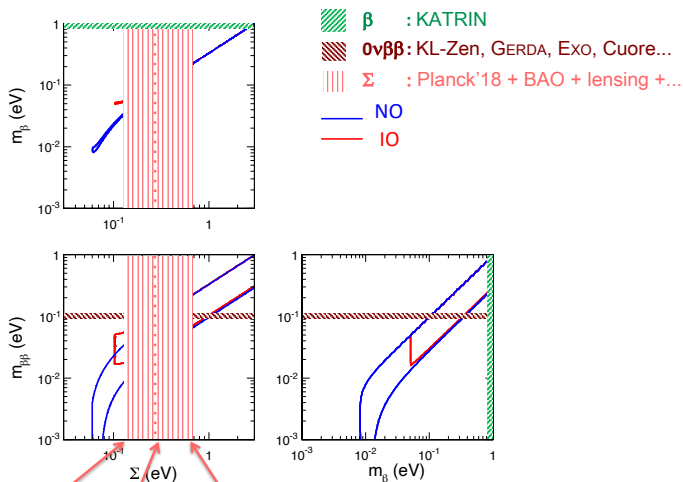
Note 2: The three observables are correlated by oscillation data →



## Impact of oscillations on non-oscillation parameter space ( $2\sigma$ )

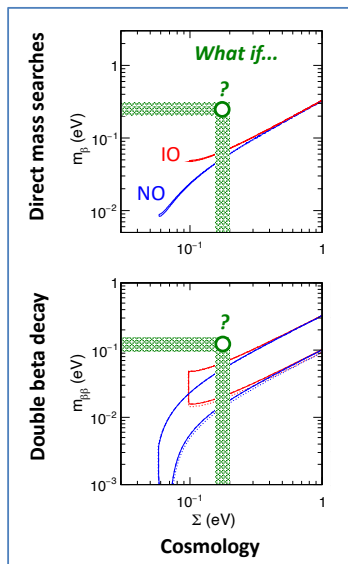


Cosmology: variety of upper bounds, with IO “under pressure”



“Aggressive” “Default” “Conservative” cosmological limits

Future data might also bring us beyond  $3\nu$  and re-shape the field...



**Lack of convergence** among data (barring expt mistakes) **might point towards new possibilities:**

- *Cosmology beyond  $\Lambda$ CDM*
  - *Alternative DBD mechanisms*
  - *New interactions (NSI)*
  - *New neutrino states ...*
- See talks by: de Gouvea, Tyagi

Main contender in current  $\nu$  physics: **Light sterile  $\nu$  at  $O(1 \text{ eV})$  scale** but... confusing/unconfirmed hints  
 → See talk by Ternes

**In any case:** generic expectations for new possible  $\nu$  mass state(s)

## Indirect probes of C $\nu$ B

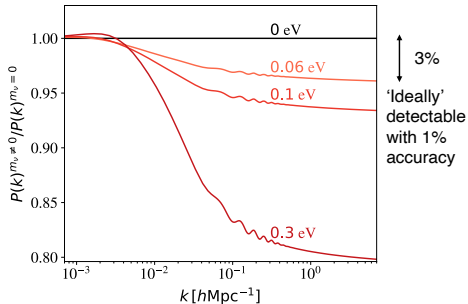
- Impact on power spectrum of matter density fluctuations

$$\delta_{\text{cdm}}^{m_\nu=0} \propto a$$

$$\delta_{\text{cdm}}^{m_\nu \neq 0} \propto a^{1-\frac{3}{5}\frac{\Omega_\nu}{\Omega_m}}$$

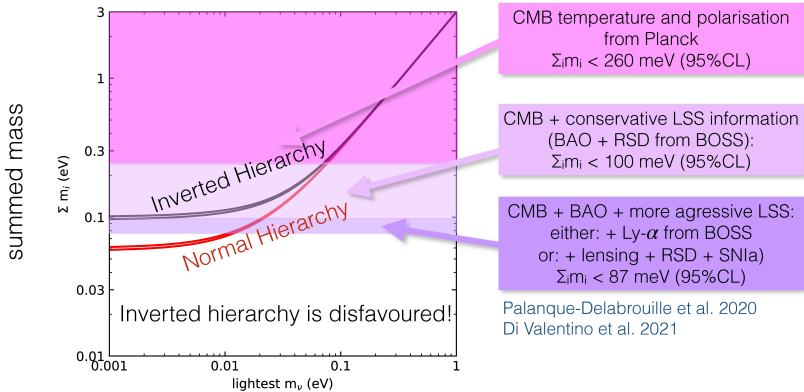
Planck TT,TE,EE + low E + lensing + BAO  
 $\Sigma m_\nu < 0.12$  eV (95%cl)

Planck TT,TE,EE + low E + lensing + BAO  
 +Lyman- $\alpha$   
 $\Sigma m_\nu < 0.09$  (95%cl) [Palanque-Delabrouille et al. 2020]



Bounds on  $\Sigma m_\nu$

95%CL upper bounds on  $\Sigma_i m_i$  for 7 parameters



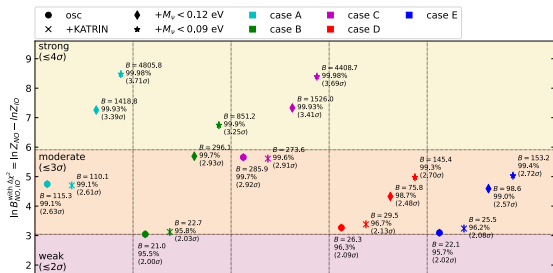
- ▶ The bound  $\sum_i m_i < 0.09 \text{ eV}$  (95% C.L.) was obtained assuming  $\sum_i m_i > 0$ .
- ▶ An analysis adopting the oscillation prior  $\sum_i m_i \gtrsim 0.05 \text{ eV}$  would return a looser bound.
- ▶ Thus, inverted hierarchy cannot be considered as disfavored at the  $2\sigma$  level.

[Palanque-Delabrouille, Yeche, Schoneberg, Lesgourgues, Walther, Chabanier, Armengaud, arXiv:1911.09073]

# The debate over the hierarchy

Degenerate hierarchy (DH) approximation:  $m_1 = m_2 = m_3$

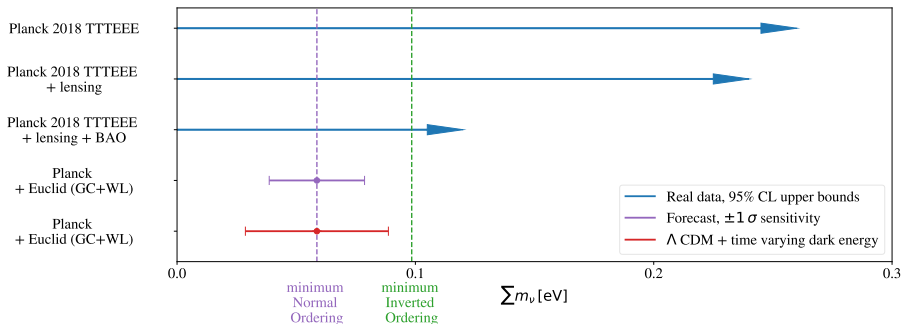
«Moderate evidence, mostly driven by neutrino oscillation data» [Gariazzo et al. 2022](#) (see also [Hergt et al. 2021](#))



case A is based on Jimenez et al.: a Gaussian prior on the logarithm of the three neutrino mass eigenstates

“The significance of the preference in favor of NO changes significantly when we consider different parameterizations.”

## Neutrino mass constraints: the future

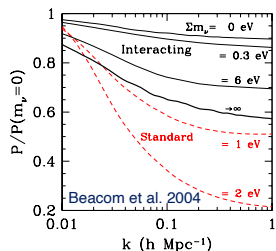




# The nightmare scenario

Cosmological data are more and more pointing towards  $\Sigma m_\nu < 0.06$  eV.

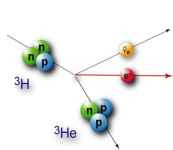
- Extended particle physics models (beyond SM)
  - Neutrino self-interactions and annihilation [“Neutrinoless Universe” Beacom et al. 2004, Esteban et al. 2021, Blinov et al. 2020, Kreisch et al. 2020, Archidiacono et al. 2020]



The interesting (exciting) scenario!

# Direct Neutrino Mass Measurement with Beta Decay

Use **Kinematics only**, look at the end-point shape

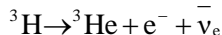


$$\frac{dN}{dE_e} = C \cdot F(E, Z) \cdot P_e \cdot (E_e + m_e c^2) \cdot (E_o - E_e) \sqrt{(E_o - E_e)^2 - m_{\nu_e}^2}$$

(some details/corrections not included)

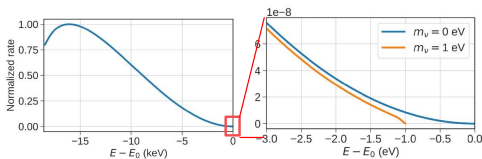
$$\sum_i |U_{ei}|^2 \cdot m_i^2 \sim m_i^2$$

in degenerated region



## Tritium as beta-source

- **low end-point** (18.6 keV)
  - relatively large deformation
  - electro-statically reachable
- **short life** (12.3 y):
  - small source amount
  - less scattering in source
- super-allowed transition
  - matrix element reliably calculable
- simplest molecular:
  - molecular states calculable



only  $2 \times 10^{-13}$  of all beta in last 1 eV

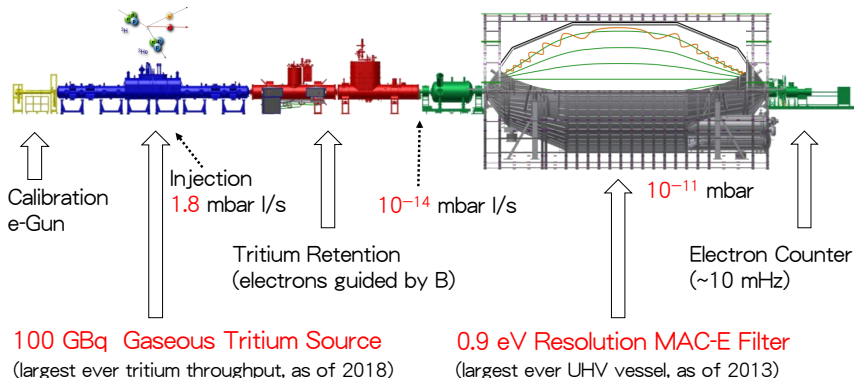
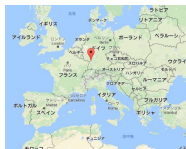
## Needs:

- strong stable source
- high precision spectroscopy

# KATRIN Experiment (Design)

## KARlsruhe TRItium Neutrino Experiment

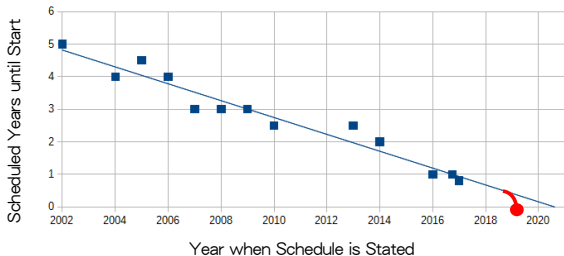
- located at Karlsruhe Institute of Technology, Karlsruhe, Germany
- design sensitivity:  $m(\nu_e) < 0.2 \text{ eV}$  (90%CL, 3 years)



All numbers are from KATRIN Design Report (2004)

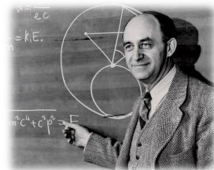
# KATRIN Construction History

- First Collaboration Meeting in 2001
- Design Report in 2004, planned to start in 4 years
- Start data taking in 2019



KATRIN has demonstrated end-points are deformed!!!

and that there can be one correct red point in a plot with many wrong blue points



Versuch einer Theorie der  $\beta$ -Strahlen. I<sup>1)</sup>.  
 Von E. Fermi in Rom.  
 Mit 3 Abbildungen. (Kingsgesungen am 16. Januar 1934.)

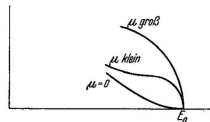
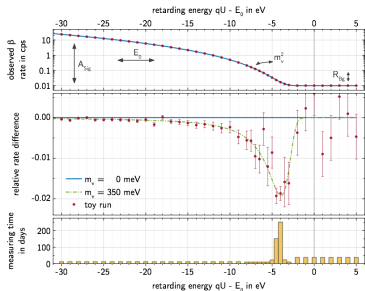


Fig. 1.

# Neutrino Mass Analysis and Results

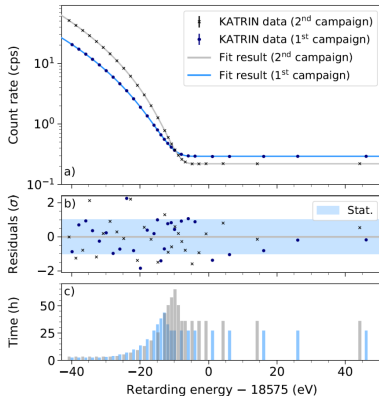
## Spectrum Model and Measurement Strategy (MC made before measurements)



$$N(qU) = \underbrace{A}_{\text{retarding energy}} \cdot \underbrace{\int_{qU}^{E_0} \frac{d\Gamma}{dE}(E; m_\nu^2, E_0)}_{\text{beta decay spectrum}} \cdot \underbrace{f(qU, E)}_{\text{apparatus response}} dE + \underbrace{B}_{\text{backgrounds}}$$

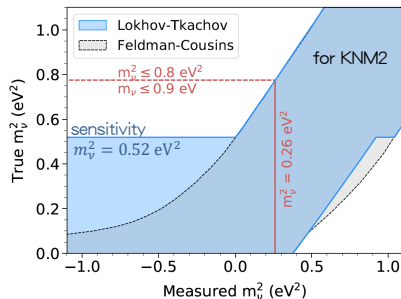
plus a number of nuisance parameters

## Measured Spectra and Fitting Results



$$m_\nu^2 = 0.26 \pm 0.34 \text{ eV}^2$$

# Mass Limit Setting



## Feldman-Cousins (FC)

- De facto standard
- Provided only for supplementary

## Lokhov-Tkachov (LT)

- No “tighter limit” from negative  $m^2$
- For negative  $m^2$ , stop at the sensitivity
- Coverage is still correct (no flip-flopping)

A. Lokhov is  
in this meeting!

## KNM 1

$$m_\nu^2 = -1.0^{+0.9}_{-1.1} \text{ eV}^2$$

$$m_\nu < 1.1 \text{ eV (LT 90\%CL)}$$

$$(m_\nu < 0.8 \text{ eV (FC 90\%CL)})$$

## KNM 2

$$m_\nu^2 = 0.26 \pm 0.34 \text{ eV}^2$$

$$m_\nu < 0.9 \text{ eV (90\%CL)}$$

## KNM 1 & 2 combined

$$m_\nu^2 = 0.1 \pm 0.3 \text{ eV}^2$$

$$m_\nu < 0.8 \text{ eV (90\%CL)}$$

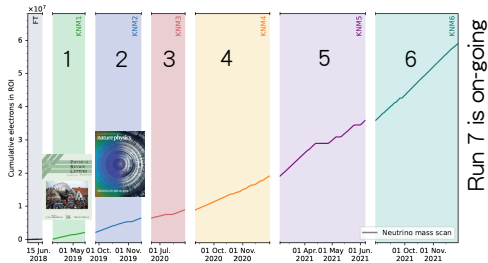
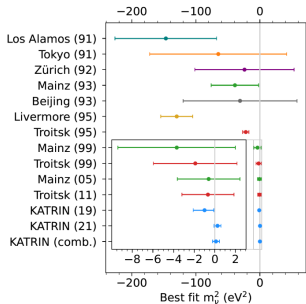
## Independent Consistency Check: Beta-decay Q-value

$$\text{KATRIN: } 18575.20 \pm 0.60 \text{ eV}$$

$$\Delta m (^3\text{He} - ^3\text{H}): 18575.72 \pm 0.07 \text{ eV (this information is not used for spectrum fitting)}$$

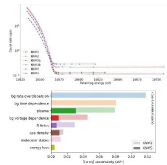
# Conclusions and Outlook

- First sub-eV result:  $m_\nu^2 = 0.1 \pm 0.3 \text{ eV}^2$ ,  $m_\nu < 0.8 \text{ eV}$  (90%CL)
- Final design sensitivity 0.2 eV (90%CL) in 5 years

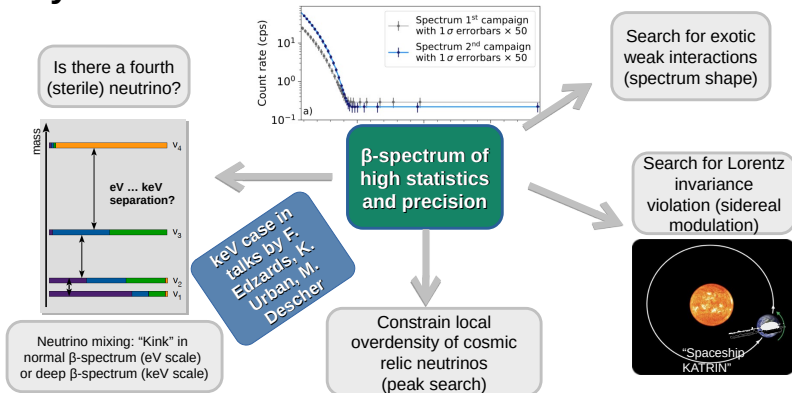


- Much more data already in hand, being analyzed (still blinded)
- After KNM 2: BG have been significantly reduced
- After KNM 2: Systematics have been significantly reduced

Preliminary estimated sensitivity up to KNM 5:  $<0.5 \text{ eV}$  (90%CL)  
 Unblinding of KNM 3-5 is planned in this summer



# “Beyond neutrino mass” in KATRIN





## Conclusions

RAA mostly resolved for some flux models

The Gallium anomaly is in strong tension with the analysis of reactor rate data

No indication from ratio analyses, Neutrino-4 result doubtful

First MicroBooNE data do not confirm the MiniBooNE excess (but can not rule it out either)

No (significant) signal in atmospheric or accelerator experiments

A global 3+1 fit is statistically not acceptable

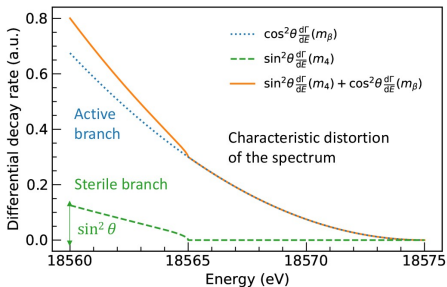
More data is needed to clarify open issues

Anomalous AND also Null results have to be checked

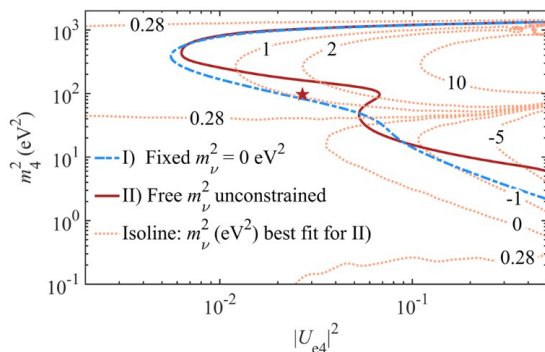
# Sterile neutrinos signature in $\beta$ -spectrum

- 3+1 sterile neutrino model
- Same data-set as for the neutrino mass
- Grid search in  $m_4, |U_{e4}|^2$  plane

$$\frac{d\Gamma}{dE} = \underbrace{(1 - |U_{e4}|^2) \frac{d\Gamma}{dE}(m_\beta^2)}_{\text{light neutrino}} + \underbrace{|U_{e4}|^2 \frac{d\Gamma}{dE}(m_4^2)}_{\text{heavy neutrino}}$$



## Interplay of $m_\nu^2$ and $m_4^2$ (2<sup>nd</sup> campaign)



Fixed  $m_\nu^2 = 0$

$$m_4^2 = 0.28 \text{ eV}^2, |U_{e4}| = 1.0$$

$$\Delta \chi_{null}^2 = 0.74$$

Free  $m_\nu^2$

$$m_4^2 = 98.3 \text{ eV}^2, |U_{e4}| = 0.027$$

$$\Delta \chi_{null}^2 = 2.49, m_\nu^2 = 1.1 \text{ eV}^2$$

KATRIN Collab., PRD 105, 072004 (2022)

13 08.06.2022

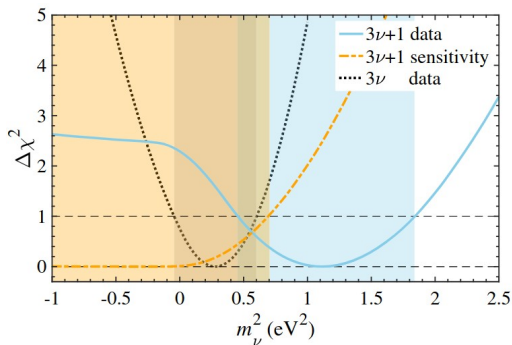
Alexey Lokhov – KATRIN beyond the neutrino mass

Institute of experimental particle physics

Please try to find the bound for free  $m_\nu$ :

- ▶ In the  $|U_{e4}|^2 - \Delta m_{\text{SBL}}^2$  plane, with  $\Delta m_{\text{SBL}}^2 = m_4^2 - m_\nu^2$  and  $m_4^2 > m_\nu^2$ .
- ▶ Considering only positive  $m_\nu^2$ .

## Interplay of $m_\nu^2$ and $m_4^2$ (2<sup>nd</sup> campaign)



- Sizable correlation of  $m_\nu^2$  and  $m_4^2$ 
  - reduction in  $m_\nu^2$  sensitivity
- Strong correlation for  $m_\nu^2 < 0 \text{ eV}^2$ 
  - flat  $\chi^2$  profile → loss of sensitivity
  - restored by external constraints
- $m_\nu^2 > 0 \text{ eV}^2$  → x2 uncertainty on  $m_\nu^2$
- Fully restore sensitivity using
  - $|U_{e4}|^2 < 10^{-4}$

KATRIN Collab., PRD 105, 072004 (2022)

$$m_\nu \lesssim \sqrt{2.3 \text{ eV}^2} \simeq 1.5 \text{ eV} \quad (90\% \text{ C.L.}) \quad \text{in } 3\nu+1$$

about 1.7 times larger than  $m_\nu \lesssim \sqrt{0.8 \text{ eV}^2} \simeq 0.9 \text{ eV} \quad (90\% \text{ C.L.}) \quad \text{in } 3\nu !$

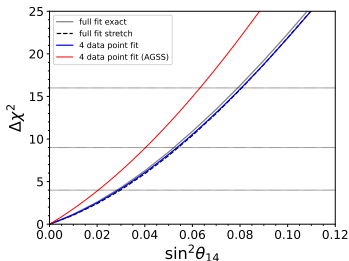
- ▶ I think that the interpretation depends on the values of  $|U_{e4}|^2$  corresponding to the marginalized  $\Delta\chi^2$  curve.
- ▶ The sensitivity to  $m_\nu$  obviously decreases for increasing  $|U_{e4}|^2$ , because

$$\frac{d\Gamma}{dE} = (1 - |U_{e4}|^2) \frac{d\Gamma}{dE}(m_\nu) + |U_{e4}|^2 \frac{d\Gamma}{dE}(m_4)$$

- ▶ A reliable bound on  $|U_{e4}|^2$  is the solar neutrino limit:

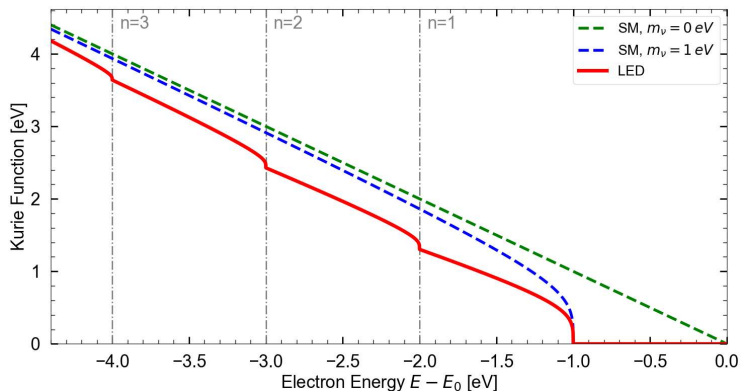
[Goldhagen, Maltoni, Reichard, Schwetz, arXiv:2109.14898]

$$\sin^2 \theta_{14} = |U_{e4}|^2$$

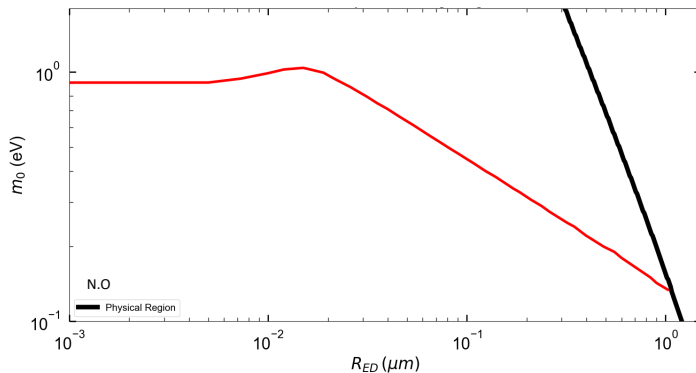


- ▶ I think that it would be useful to see:
  - ▶ The  $\Delta\chi^2$  contours in the  $m_\nu - |U_{e4}|^2$  plane.
  - ▶ The  $\Delta\chi^2$  for  $m_\nu$  with a solar  $\Delta\chi^2$  pull term.

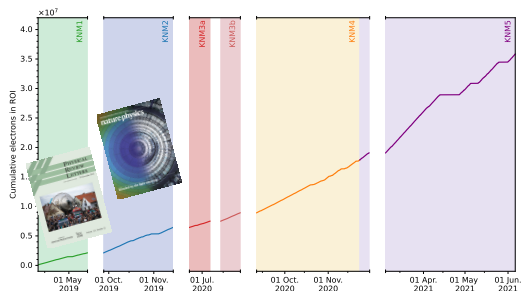
## Beta-decay Experiments



## Bound



## Measurement campaigns



KATRIN Neutrino mass Measurements

	Time (hrs)	$\rho d\sigma$ (m <sup>-2</sup> )	B <sub>g</sub> (mcps)
KNM1	522	$1.11 \times 10^{21}$	370
KNM2	294	$4.23 \times 10^{21}$	278
KNM3a	220	$2.08 \times 10^{21}$	137
KNM3b	224	$3.75 \times 10^{21}$	258
KNM4	1267	$3.77 \times 10^{21}$	150
KNM5	1232	$3.78 \times 10^{21}$	160

- Published results: KNM1 and KNM2  
Phys. Rev. Lett. 123, 221802  
Nat. Phys. 18, 160–166 (2022)
- Current analysis: KNM1 – KNM5
- Data-taking: KNM6, KNM7, ...

Campaigns

●●○○○

4/17 Wednesday 8<sup>th</sup> June 2022

Spectra fitting

○○

Leonard Köllenberger: (KATRIN) Details on the  $\nu$  mass analysis

Unbiased analysis

○

Systematics

○○○○○

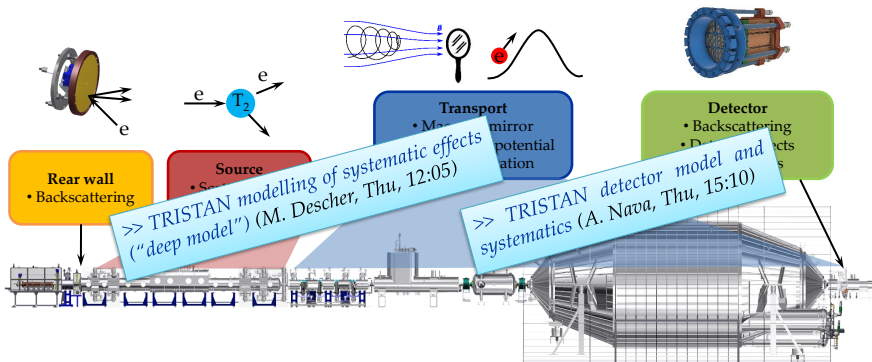
Outlook

○

Institute for Astroparticle Physics

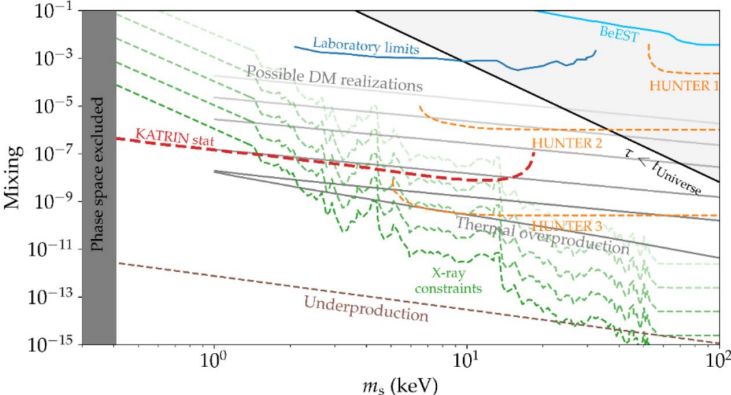


# KATRIN - Search for keV sterile $\nu$



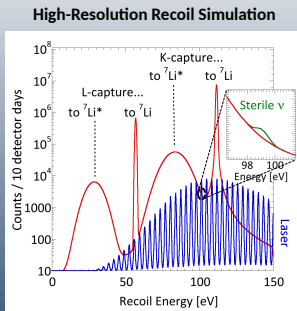
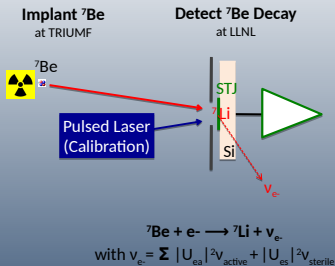
*>> Precise modeling of all relevant effects required*  
*>> Hardware modifications required*

### KATRIN keV sensitivity limit





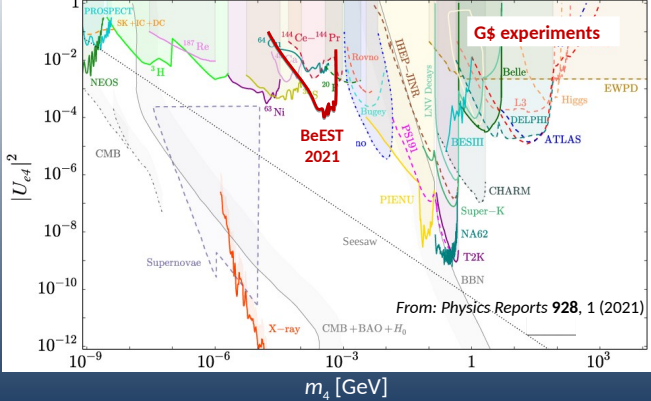
# The BeEST Sterile Neutrino Experiment



Calibrate STJ with pulsed laser.

Four peaks due to K- and L- capture into  ${}^7\text{Li}$  ground and excited state

# The BeEST in Context



## Cyclotron Radiation Emission Spectroscopy (CRES)



**PROJECT 8**

Frequency Approach



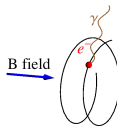
A. L. Schawlow

*"Never measure anything but frequency."*



O. Heaviside

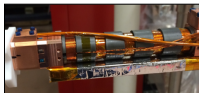
*Measure the cyclotron radiation  
from a single electron*



- Source transparent to microwave radiation
- No e- transport from source to detector
- Leverages precision inherent in frequency techniques

B. Monreal and JAF, Phys. Rev D80:051301

Cyclotron Radiation  
Emission  
Spectroscopy  
(CRES)

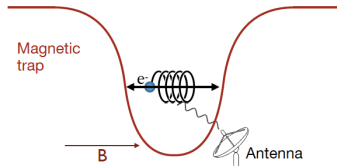


**PROJECT 8**

Frequency Approach



$$f_c = \frac{f_{c,0}}{\gamma} = \frac{1}{2\pi} \frac{eB}{m_e c^2 + E_{\text{kin}}}$$



$$f_{c,0} = 27.992\,491\,10(6) \text{ GHz T}^{-1}$$

- *Narrow band region of interest (@26 GHz).*
- *Small, but detectable power emitted.*

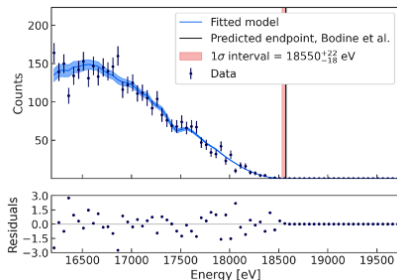
$$P(17.8 \text{ keV}, 90^\circ, 1 \text{ T}) = 1 \text{ fW}$$

$$P(30.2 \text{ keV}, 90^\circ, 1 \text{ T}) = 1.7 \text{ fW}$$

## Phase II Results

Phase II CRES instrument provides 1 mm<sup>3</sup> volume inside waveguide. Total of 3770 events observed over 3 months of data taking.

First endpoint CRES measurement conducted with no observed background in 81 days of data taking.



### T<sub>2</sub> endpoint

Frequentist:  $E_0 = (18550^{+22}_{-18})$  eV (1σ)

Bayesian:  $E_0 = (18553^{+17}_{-17})$  eV (1σ)

### Neutrino mass

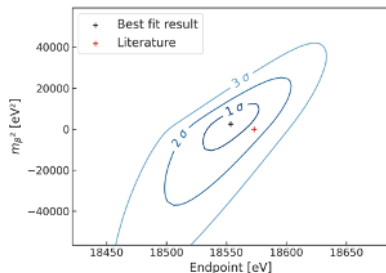
Frequentist:  $\leq 178$  eV/c<sup>2</sup> (90% C.L.)

Bayesian:  $\leq 169$  eV/c<sup>2</sup> (90% C.L.)

### Background rate

$\leq 3 \times 10^{-10}$  eV<sup>-1</sup>s<sup>-1</sup> (90% C.L.)

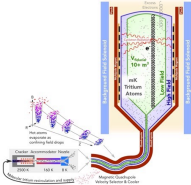
### First CRES Mass Limit



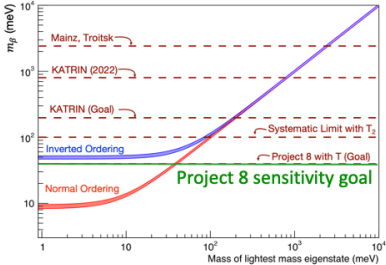
The two R&D efforts will combine into a demonstrator atomic cavity experiment at the end of Phase III, with a projected sensitivity of  $400 \text{ meV}/c^2$ .

A pathfinder for the final experimental goal of  $40 \text{ meV}/c^2$ .

CRES Phase III Pathfinder Experiment



Project 8  $\nu$  Mass Scale Sensitivity



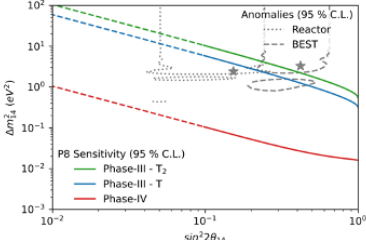


*A low background, high resolution tritium experiment can help resolve “kinks” in the spectrum.*

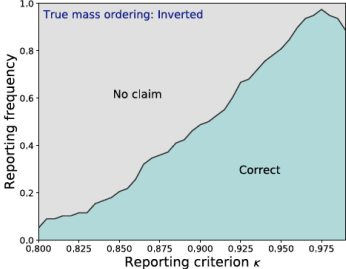
*Potentially sensitive to mass splittings or low mass sterile neutrinos.*

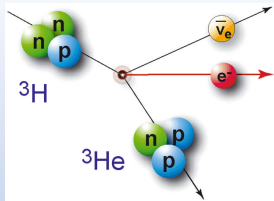
PHYSICAL REVIEW C 103, 065501 (2021)  
arXiv:2203.07349v1

Project 8 Sterile Neutrino Sensitivity



Sensitivity to Neutrino Mass Ordering





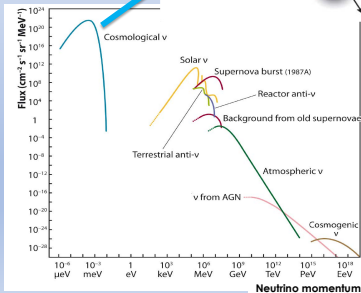
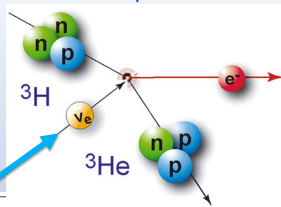
Tritium  $\beta$ -decay  
(12.3 yr half-life)

Neutrino momentum  $\sim 0.17$  meV

For  $m_\nu = 50$  meV,  
 $KE = p^2/2m$   
 $= 0.17$  meV  $(0.17 \text{ meV}/100 \text{ meV})$   
 $= 0.3 \mu\text{eV}$

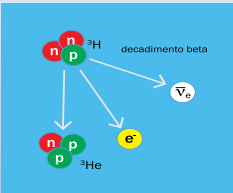
**Ultra-Cold!**

## Neutrino capture on Tritium

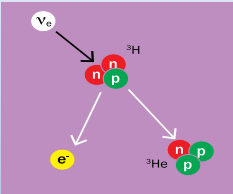
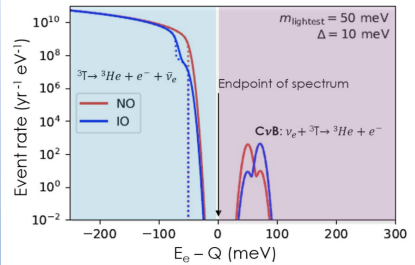


# Detection Concept: Neutrino Capture

- Basic concepts for relic neutrino detection were laid out in a paper by Steven Weinberg in **1962** [Phys. Rev. 128:3, 1457] applied for the first time to massive neutrinos in **2007** by Cocco, Mangano, Messina [DOI: 10.1088/1475-7516/2007/06/015] and revisited in **2021** by Cheipesh, Cheianov, Boyarsky [https://arxiv.org/abs/2101.10069]



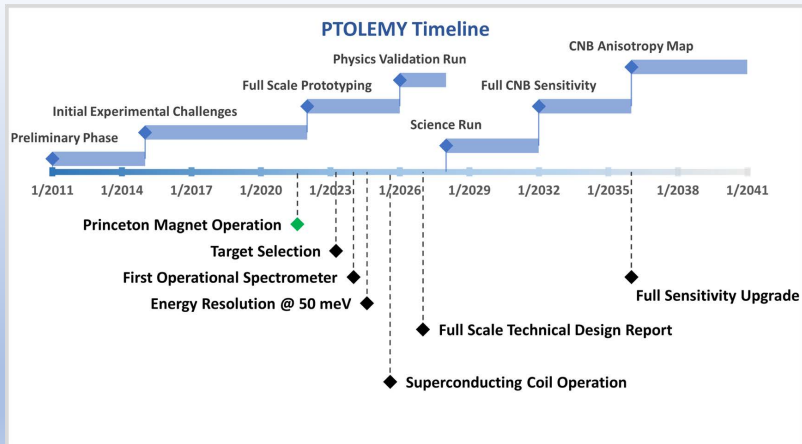
**What do we know?**  
 Gap (2m) constrained to  
 $m < \sim 200\text{meV}$   
 from precision cosmology  
 Electron flavor expected with  
 $m > \sim 50\text{meV}$   
 from neutrino oscillations



**CvB Detection Requires:**

few  $\times 10^{-6}$  energy resolution set by  $m_\nu$   
 KATRIN  $\sim 10^{-4}$  (current limitation)

**PTOLEMY:**  $10^{-4} \times 10^{-2}$   
 (compact filter)  $\times$  (microcalorimeter)



## Calorimetric approach as a viable alternative to spectrometers

■ **Pro:** Most of the unwanted source related effects are avoided.

■ **New way to probe sub-eV neutrino mass scale?**

■ A good isotope should have:

Low Q value

Proximity of a peak near the ROI

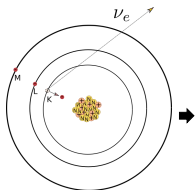
Short half life to reduce the experimental challenges

### Ideal calorimetric experiment

- The radioactive source is embedded in the detector(s)
- Only the neutrino energy escape detection.
- Important limits on the source intensity (statistics) that can be accumulated
- Activity also limited by the relation between energy resolution and detector size.

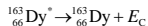
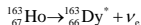
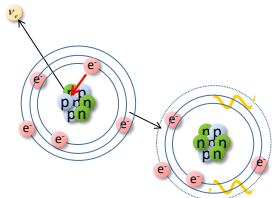
High number of events in the ROI

■ No convincing isotopes alternatives to  $^3\text{H}$  and  $^{163}\text{Ho}$  (yet).



Isotope	Q value [eV]	Half life [y]	Decay	B.R	Experiments
$^3\text{H}$	18592.01(7)	12	$\beta^-$	1	Simpson's
$^{187}\text{Re}$	2470.9(13)	$4.3 \times 10^{10}$	$\beta^-$	1	MANU, MIBETA
$^{163}\text{Ho}$	2833(30)	4570	EC	1	ECho, Holmes
$^{135}\text{Cs}$	440	$8.0 \times 10^{11}$	$\beta^-$	$1.6 \times 10^{-6}$	-
$^{115}\text{In}$	155	$4.3 \times 10^{20}$	$\beta^-$	$1.1 \times 10^{-6}$	-

# Electron Capture in $^{163}\text{Ho}$



- $\tau_{1/2} \cong 4570$  years ( $2 \cdot 10^{11}$  atoms for 1 Bq)
- $Q_{\text{EC}} = (2.833 \pm 0.030^{\text{stat}} \pm 0.015^{\text{syst}})$  keV  
S. Eliseev et al., *Phys. Rev. Lett.* **115** (2015) 062501

$$Q_{\text{EC}} = m(^{163}\text{Ho}) - m(^{163}\text{Dy})$$

Penning Trap Mass Spectroscopy

@TRIGA TRAP (Uni-Mainz) (♦)

@SHIPTRAP (GSI – Darmstadt) (♦♦)

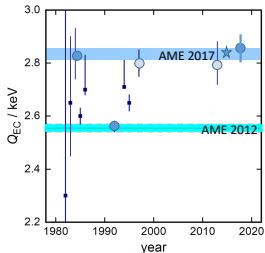
Future goal: 1 eV precision:

PENTATRAP @MPIK, Heidelberg (\*)

(♦) F. Schneider et al., *Eur. Phys. J. A* **51** (2015) 89

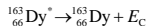
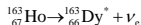
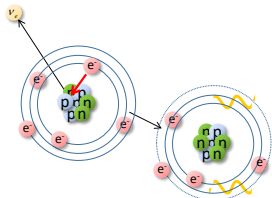
(♦♦) S. Eliseev et al., *Phys. Rev. Lett.* **115** (2015) 062501

(\*) J. Repp et al., *Appl. Phys. B* **107** (2012) 983, C. Roux et al., *Appl. Phys. B* **107** (2012) 997



## Electron Capture in $^{163}\text{Ho}$ – Spectrum

2



- $\tau_{1/2} \cong 4570$  years ( $2 \cdot 10^{11}$  atoms for 1 Bq)
- $Q_{\text{EC}} = (2.833 \pm 0.030^{\text{stat}} \pm 0.015^{\text{syst}})$  keV  
S. Eliseev et al., *Phys. Rev. Lett.* **115** (2015) 062501



Source = Detector

### Calorimetric measurement

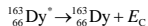
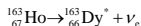
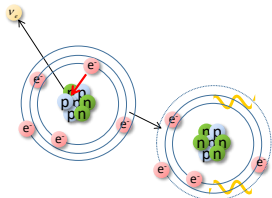
A. De Rujula and M. Lusignoli, *Phys. Lett.* **118B** (1982)

- |                |   |
|----------------|---|
| Advantages:    | Measured neutrino complementary spectrum<br>No final state problems |
| Disadvantages: | Unresolved pile-up  |



# Electron Capture in $^{163}\text{Ho}$ – Spectrum

3



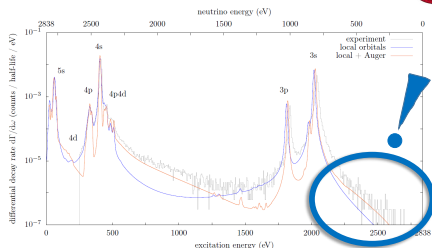
- $\tau_{1/2} \cong 4570$  years ( $2 \cdot 10^{11}$  atoms for 1 Bq)
- $Q_{\text{EC}} = (2.833 \pm 0.030^{\text{stat}} \pm 0.015^{\text{syst}})$  keV  
S. Eliseev et al., *Phys. Rev. Lett.* **115** (2015) 062501



Source = Detector

## Calorimetric measurement

A. De Rujula and M. Lusignoli, *Phys. Lett.* **118B** (1982)

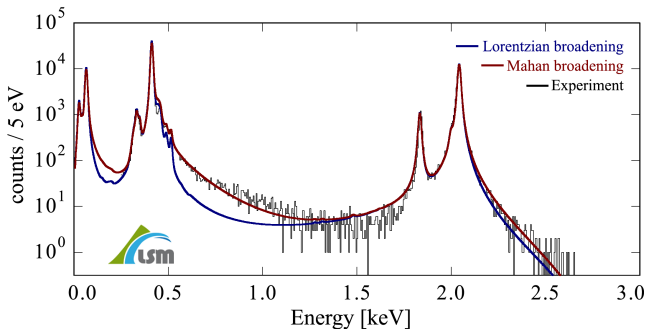


M. Braß and M. W. Haverkort, *New J. Phys.* **22** (2020) 093018



## Proof of concept

20



C. Velte et al., EPJC **79** (2019) 1026

Energy resolution

$$\Delta E_{\text{FWHM}} = 9.2 \text{ eV}$$

Background level

$$b < 1.6 \times 10^{-4} \text{ events/eV/pixel/day}$$

4 day measurement with 4 pixels loaded with  $\sim 0.2 \text{ Bq } ^{163}\text{Ho}$

- measurement performed underground
- test for data reduction and spectral shape analysis

- $Q_{\text{EC}} = (2838 \pm 14) \text{ eV}$
- $m(\nu_e) < 150 \text{ eV}$  (95% C.L.)

## ECHo-1k high statistics spectrum

### ECHo-1k chip-Au

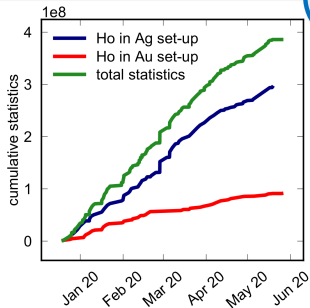
15 channels  
2 temperature channels  
23 pixel with implanted  $^{163}\text{Ho}$   
3 background pixels  
average activity = 0.94 Bq  
total activity of 28.1 Bq

### ECHo-1k chip-Ag

22 channels  
2 temperature channels  
34 pixel with implanted  $^{163}\text{Ho}$   
6 background pixels  
average activity = 0.71 Bq  
total activity of 25.9 Bq

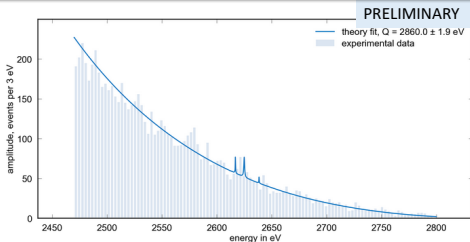
A number of  $^{163}\text{Ho}$  events larger than  $10^8$  has been acquired in the first months of 2020

This statistics allow for investigating the value of the electron neutrino effective mass down to 20 eV



## 2.5 keV < E < 2.8 keV

26



Determination of  $Q_{EC}$  by fitting the spectrum using:

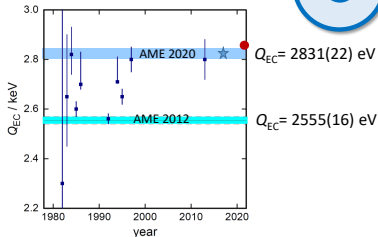
- Brass & Haverkort theory
- Flat background

$$Q_{EC} = (2860 \pm 2_{stat} \pm 5_{syst}) \text{ eV}$$



Systematic uncertainties related to theoretical spectral shape

...still too large for analysis of smaller endpoint region



$$\star Q_{EC} = (2.833 \pm 0.030^{stat} \pm 0.015^{syst}) \text{ keV}$$

S. Eliseev et al., *Phys. Rev. Lett.* **115** (2015) 062501

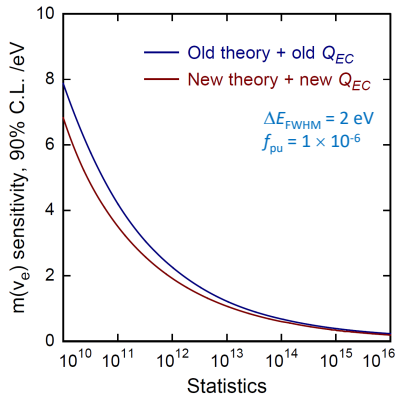
Waiting for new PENTATRAP\* results

(\*) J. Repp et al., *Appl. Phys. B* **107** (2012) 983

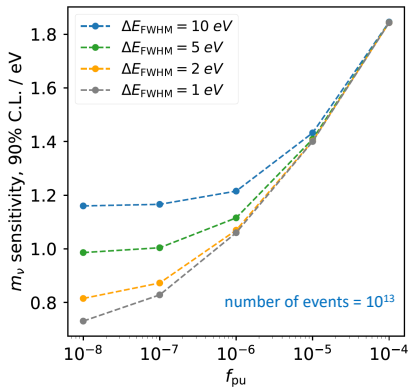
C. Roux et al., *Appl. Phys. B* **107** (2012) 997

# Updated sensitivity

Brass & Haverkort theoretical model + new  $Q_{EC}$ -value



Sensitivity for the coming phase of ECHO



## Conclusions

31

- ✓ The results obtained with  $^{163}\text{Ho}$  loaded MMCs paved the way to large scale neutrino mass experiments based on  $^{163}\text{Ho}$
- ✓ The ECHo collaboration has already contributed to a [more precise description](#) of the  $^{163}\text{Ho}$  spectrum
- ✓ A first improvement on the effective electron neutrino mass limit has been obtained in a [proof of concept](#) measurement
- ✓ More than  $10^8$   $^{163}\text{Ho}$  events have been acquired within the ECHo-1k phase → A new limit at the level of [20 eV on the effective electron neutrino mass](#) is coming soon
- ✓ Important steps [towards ECHo-100k](#) have been demonstrated: new ECHo-100k array + multiplexed readout

ECHo

... not only cool because of mK temperature



## Holmes

- Holmes is an ambitious project that aims to verify the feasibility of the calorimetric approach to the neutrino mass determination.
- High performing detectors are needed, in terms of energy resolution  $\Delta E$  and time resolution  $\tau_R$ : LTD



$\Delta E = \Delta E (A)$   
 $\tau_R = \tau_R (A)$   
{ Custom ion implanter



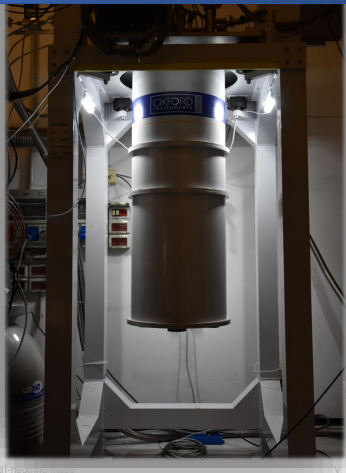
$^{163}\text{Ho}$



- Holmes has adopted a **high-risk/high-gain approach**.

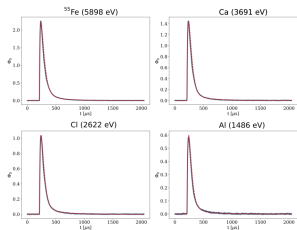
### Holmes in a nutshell

- Transition Edge Sensors (TES)  
 $\Delta E \approx 1 \text{ eV}$ ,  $\tau_R < 3 \mu\text{s}$
- Microwave multiplexing readout!
- Target activity (A) of 300 Bq/det
- $6 \times 10^5$  nuclei of  $^{163}\text{Ho}$
- $3 \times 10^{13}$  events recorded in three years
- $m_{\nu_e}$  sensitivity  $O(1) \text{ eV}$

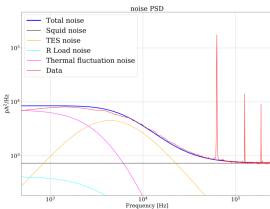


## Pile-up rejection

- The pile-up fraction  $f_{pp}$  is proportional to the time resolution  $\tau_R$ . The latter depends on the detector and readout characteristics and on the algorithms used to discriminate the signals.
- Requirements: high discrimination efficiency and **near zero energy dependence**.
- We have studied an application of the Wiener Filter while developing a new discrimination technique called DSVP.
- To test the algorithm, we wrote a tool to simulate the detector response (signal shape and noise spectrum).

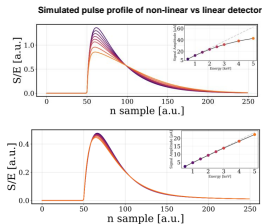


Real pulse Simulated pulse



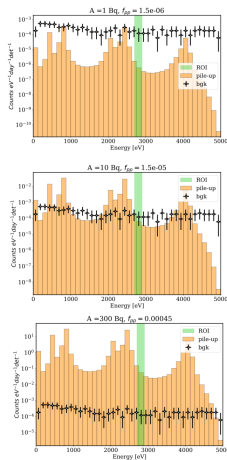
## Simulation program in a nutshell:

- Goal: create pseudo-real dataset
- Energy taken from the first order Ho spectrum
- 4-th order Runge-Kutta method to solve the  $n$  differential equations
- ARMA(p,q) to properly simulate the noise spectrum



## Conclusions

- Tested and tuned the final array fabrication processes.
  - | These did not spoil the detectors' performances.
- The software for analysis and signal processing of microcalorimeters events is up and running!
- The expected background contributions were assessed, both with simulations and dedicate measurements.
  - | A further reduction of a factor roughly 25% could be achieved with a similar setup studied in this work (muon veto).
  - | Pile-up reduction results equivalent to increase the measurement time by a factor 4: from 3 to 12 years.
- Ion implanter is working as expected. The production of a proper sputter target is almost ready!
- The first phase of the HOLMES experiment is expected on the last quarter of 2022: a low dose implantation of a  $2 \times 32$  pixel array.
  - Influence of the Ho on the detector response will be assessed
  - A high resolution Ho calorimetric spectrum will help to discriminate between the different theoretical models
  - A first limit on the neutrino mass  $O(10)$  eV will be reached
- These results will contribute to clarify if the calorimetric approach can still be considered a feasible way to reach the required sub-eV sensitivity on the neutrino mass.





## Double beta decay (DBD)

Rare nuclear decay between isobars with  $|\Delta Q|=2$

Even-even nuclei: favorable experimental condition

Decay modes:

- $(A,Z) \rightarrow (A,Z+2) + 2e^- + 2\bar{\nu}$  ( $2\nu\beta\beta$ )
  - 2<sup>nd</sup> order transition (Standard Model)
  - Very long half-lives ( $10^{18-24}$  yr)

- **$(A,Z) \rightarrow (A,Z+2) + 2e^-$  ( $0\nu\beta\beta$ )**

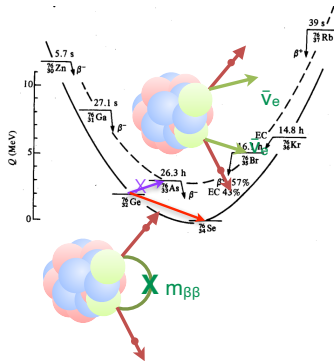
- Lepton number violation (LNV): BSM physics

Possible ONLY IF neutrinos are Majorana fermions

- LNV ( $\Delta L=2$ , B-L violation)
  - absolute neutrino mass scale
  - Majorana phases

SM extensions

- $(A,Z) \rightarrow (A,Z+2) + 2e^- + \text{exotics (e.g. } n\chi)$



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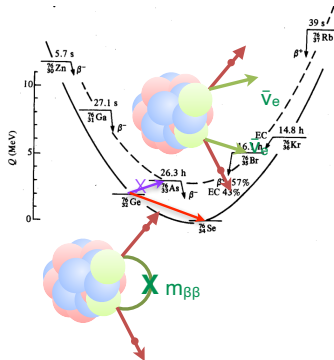
→ LNV ( $\Delta L=2$ , B-L violation)

→ absolute neutrino mass scale

→ Majorana phases

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## Process description

With nontrivial approximations, it is possible to separate atomic nuclear and particle contributions and factor the transition amplitude as

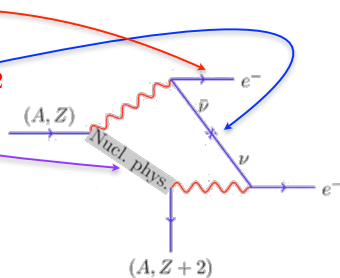
$$\Gamma^{0\nu} = G_x(Q, Z) |M_x(A, Z)|^2 |\eta_x|^2$$

$G_x(Q, Z)$  = phase space factor → calculable precisely

$M_x(A, Z)$  = nuclear matrix element (NME) → problematic

$\eta_x$  = particle physics parameter → depends on model

- massive Majorana neutrinos
- GUT's
- SUSY
- ...



## Process description

With nontrivial approximations, it is possible to separate **atomici nuclear** and **particle** contributions and factor the transition amplitude as

$$T_{1/2}^{-1} = G_{01} g_A^4 (M_{\text{light}}^{0\nu})^2 \frac{m_{\beta\beta}^2}{m_e^2} + \frac{m_N^2}{m_e^2} \tilde{G} \tilde{g}^4 \tilde{M}^2 \left(\frac{v}{\tilde{\Lambda}}\right)^6 + \frac{m_N^4}{m_e^2 v^2} \tilde{G}' \tilde{g}'^4 \tilde{M}'^2 \left(\frac{v}{\tilde{\Lambda}'}\right)^{10} + \dots$$

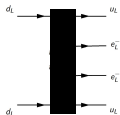
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- ...

# Lepton Number Violation (LNV) and $0\nu\beta\beta$

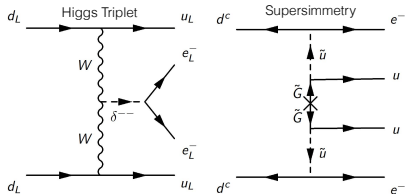
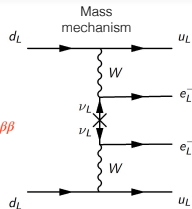


- A simple 'black box' with lots of possible modeling
- Great interest in the possibility of probing new physics

- Possibility of interference

$$T_{1/2}^{-1} = G_\nu \left| \frac{m_{ee}}{m_e} M_\nu + \epsilon M_e \right|^2$$

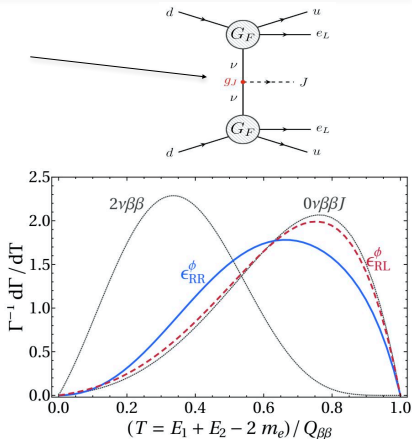
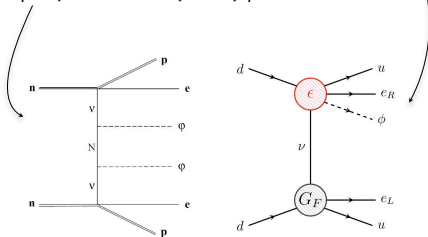
→ measurements on different isotopes



$A_{\beta\beta} \sim \text{LNV parameters}$

## Exotic channels

- Lepton number not necessarily violated: Dirac neutrinos
- J-type Majorons with coupling to neutrinos:  $g_{ij}\bar{\nu}_i\gamma_5\nu_j J$
- Light scalars associated with Weinberg-type operators
- $\phi$  Majorons with LR symmetry  $\phi$



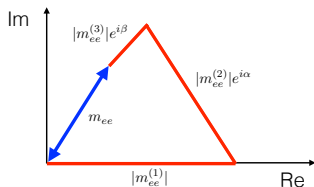
## The standard interpretation: light Majorana neutrino

- $\beta\beta_{0\nu}$  is assumed to be mediated by the exchange of light Majorana neutrinos and all other mechanisms make zero or negligible contributions
- $m_{\beta\beta}$  depends on a total of seven parameters ( $\theta_{12}$ ,  $\theta_{13}$ , three masses and two Majorana phases)
- neutrino oscillation experiments are sensitive only to the two mixing angles, two neutrino mass squared differences, and the mass ordering  
 $\Rightarrow$  only four out of seven degrees of freedom, can be bounded

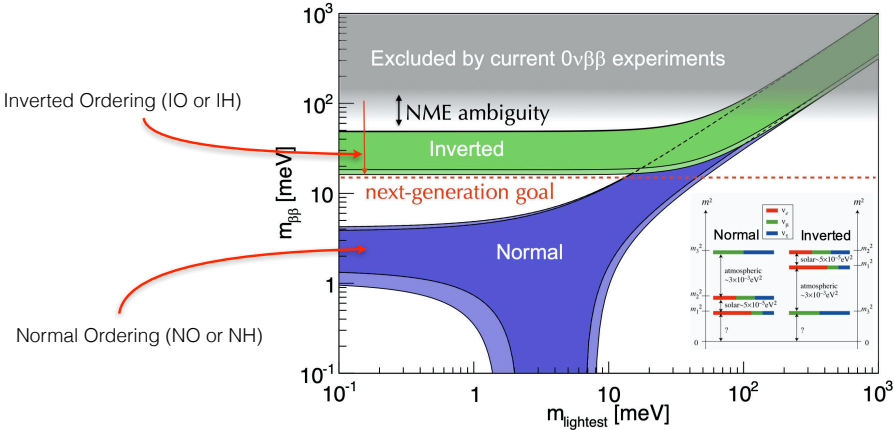
$$\eta_x = \langle m_{ee} \rangle = \sum_k U_{ek}^2 m_k$$

$$= c_{12}^2 c_{13}^2 m_1 + s_{12}^2 c_{13}^2 e^{i\alpha} m_2 + s_{13}^2 e^{i\beta} m_3$$

- Transition amplitude is proportional to the coherent sum of neutrino masses
- Majorana phases play a crucial role: possible cancellations

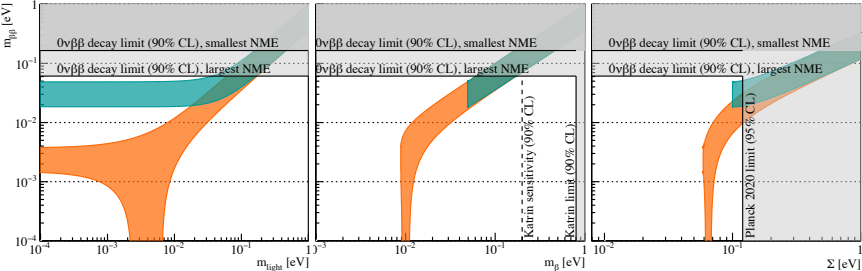
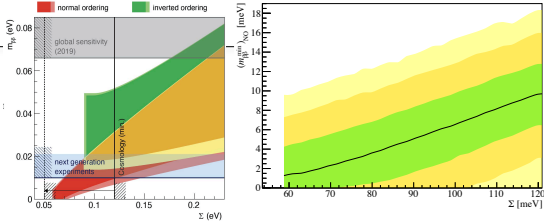


# Light Majorana neutrinos





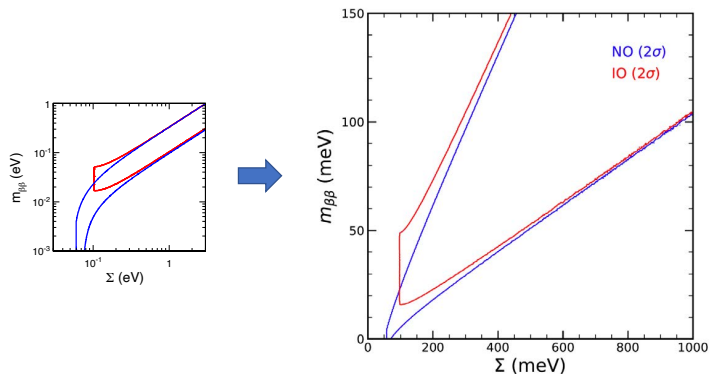
# $m_{\beta\beta}$ and the $m_\nu$ quest



## (1) Time to shift from log to linear scales!

$$m_{\beta\beta} = |U_{e1}^2 m_1 + U_{e2}^2 m_2 + U_{e3}^2 m_3|$$

$$\Sigma = m_1 + m_2 + m_3$$



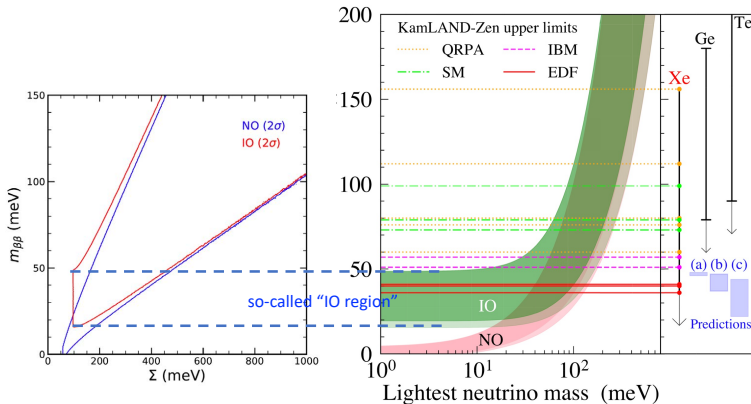
Regions allowed by current oscillation data for NO and IO separately.

[Width of regions dominated by unknown Majorana phases]

Similarly for  $m_{\beta\beta}$  vs  $\Sigma$ , without Majorana phase uncertainties

## (2) Language bonus: avoid log-inspired misleading jargon

The (improperly named) "IO region" is compatible with NO as well...

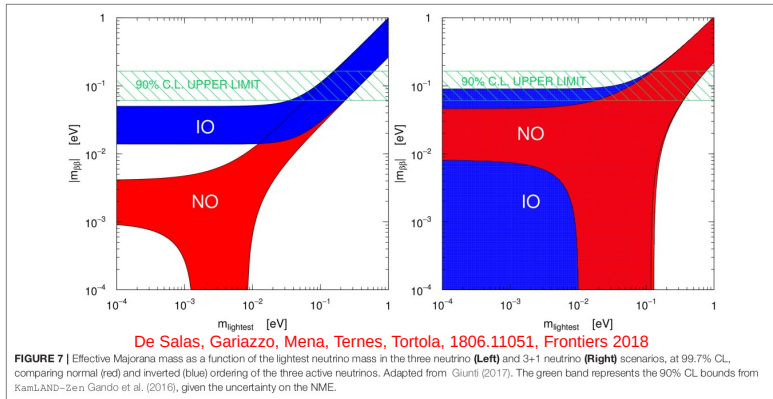


**Important to note that  $\Sigma$  is observable while  $m_{\text{lightest}}$  is not.**  
**Elongated stripes in  $\log(m_{\text{lightest}})$  have no real significance.**

**Stat. bonus: If Bayesian results depend sensitively on log vs linear priors etc.  $\rightarrow$  unreliable, IMHO**

## Neutrinoless $\beta\beta$ decay

More details in talk by  
Oliviero Cremonesi!



# $m_{\beta\beta}$ extraction from data: a hard task

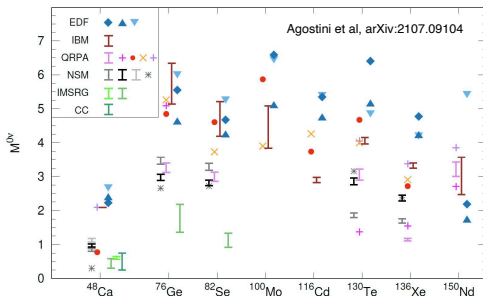
$M^{0\nu}$  values estimated from different calculation methods

$$\Gamma^{0\nu} = G_{0\nu}(Q, Z) |M_{0\nu}(A, Z)|^2 |m_{ee}|^2$$

- Factor 2-3 "uncertainty" among nuclear models.
- 'Ab initio' calculations for lighter isotopes give lower values

NME calculations:

- EDF: large
- QRPA: larger spread
- NSM: small
- IMSRG  $^{48}\text{Ca}$  ab initio: very small

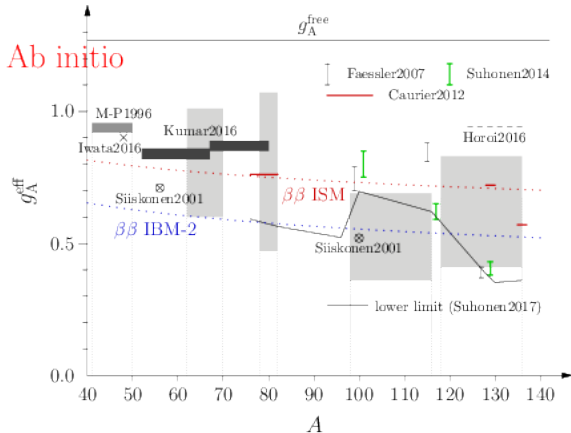


Much more problematic: nuclear matrix elements

- Significant differences
- Results inaccuracy

**g<sub>A</sub> quenching?**

Results extracted from the GT  $\beta^\pm/EC$  and  $2\nu\beta\beta$  calculations

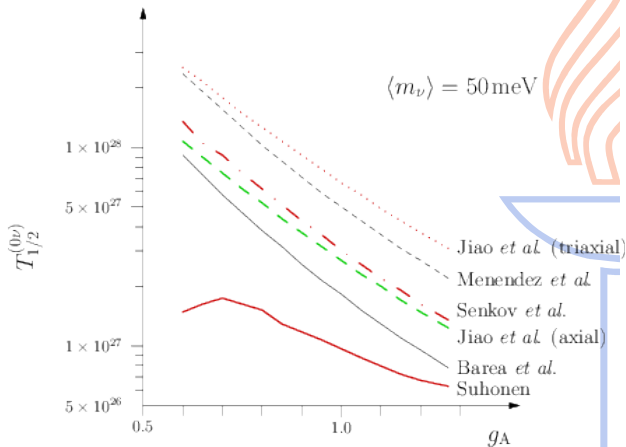


Ab initio: P. Gysbers *et al.*, Nature Physics 15 (2019) 428

- Faessler2007: pnQRPA A. Faessler *et al.*, arXiv 0711.3996v1 [Nucl-th]
- Suhonen2014: pnQRPA J. Suhonen *et al.*, Nucl. Phys. A 924 (2014) 1
- Suhonen2017: pnQRPA J. Suhonen, Phys. Rev. C 96 (2017) 055501
- Caurier2012: ISM E. Caurier *et al.*, Phys. Lett. B 711 (2012) 62
- Horoi2016: ISM M. Horoi *et al.*, Phys. Rev. C 93 (2016) 024308
- M-P1996: ISM G. Martínez-Pinedo *et al.*, Phys. Rev. C 53 (1996) R2602
- Iwata2016: ISM Y. Iwata *et al.*, Phys. Rev. Lett. 116 (2016) 112502
- Kumar2016: ISM V. Kumar *et al.*, J. Phys. G 43 (2016) 105104 Phys. Lett. B 711 (2012) 62
- Siiskonen2001: ISM T. Siiskonen *et al.*, Phys. Rev. C 63 (2001) 055501
- $\beta\beta$  ISM and IBM-2: J. Barea *et al.*, Phys. Rev. C 87 (2013) 014315
- Light hatched regions: pnQRPA H. Ejiri *et al.*, J. Phys. G 42 (2015) 055201 ; P. Pirinen *et al.*, Phys. Rev. C 91 (2015) 054309 ; F. Deppisch *et al.*, Phys. Rev. C 94 (2016) 055501

## Example: $0\nu\beta\beta$ NMEs of $^{76}\text{Ge}$ , effect on the half-life

- **Jiao *et al.*:** Phys. Rev. C 96 (2017) 054310 (GCM+ISM)
- **Menendez *et al.*:** Nucl. Phys. A 818 (2009) 139 (ISM)
- **Senkov *et al.*:** Phys. Rev. C 93 (2016) 044334 (ISM)
- **Barea *et al.*:** Phys. Rev. C 91 (2015) 034304 (IBM-2)
- **Suhonen:** Phys. Rev. C 96 (2017) 055501 (pnQRPA +  $g_{pp}$  + isospin restoration + **data on  $2\nu\beta\beta$** )



## How do we extract information on the value of $g_A$ ?

These methods are now available:

**For low momentum exchanges ( $g_A$ ):**

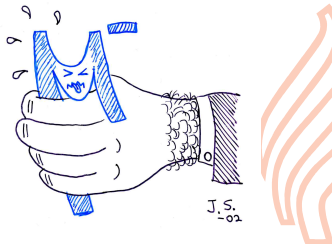
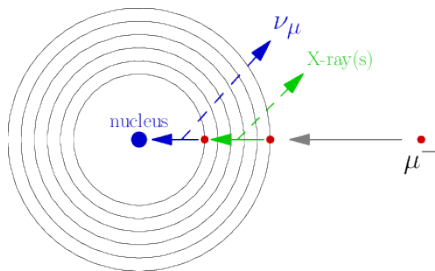
- study half-lives of  $\beta$  decays ( $1^+$  and  $2^-$  states)
- study half-lives of  $2\nu\beta\beta$  decays ( $1^+$  states)
- Study electron spectral shapes of  $\beta$  decays ( $J^\pi$  states)

**For high momentum exchanges like  $0\nu\beta\beta$  decay ( $g_{A,0\nu}$ ):**

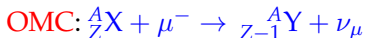
- Study nuclear muon capture ( $J^\pi$  states)



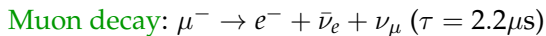
## Ordinary Muon Capture (OMC)



Nuclear muon capture:



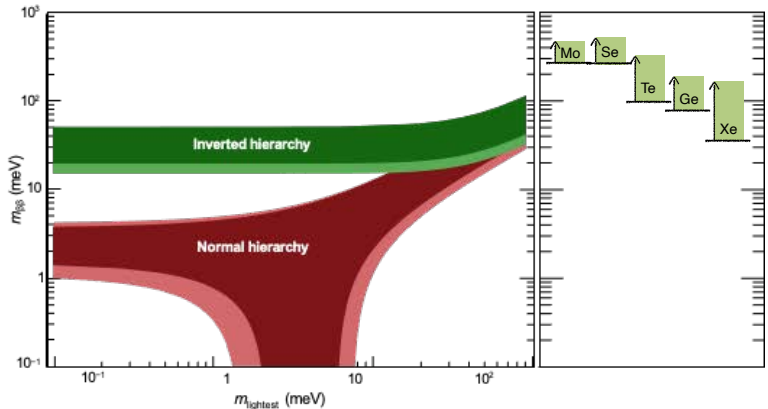
Also:



OMC probability  $\sim Z^4$   
(in Fe 91% are captured,  
breakeven at  $Z \sim 11$ )

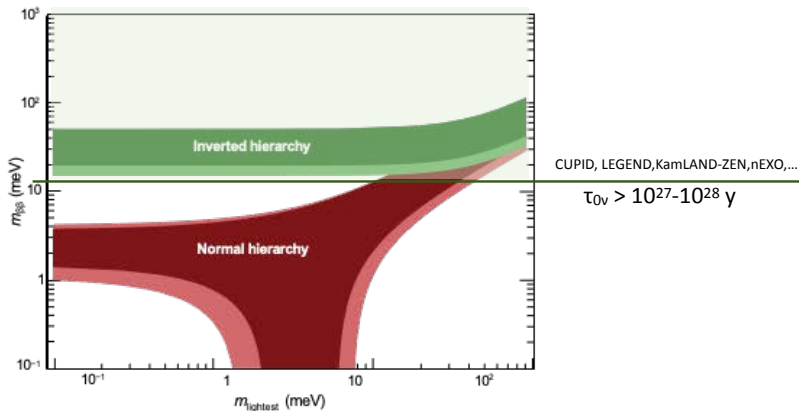
# Current results: $m_{\beta\beta}$

no  $g_A$  quenching



## Future sensitivities

no  $g_A$  quenching



## Conclusions

- ▶ The future is bright for:
  - ▶ Neutrino Physics.
  - ▶ The Determination of the Absolute Electron (Anti)-Neutrino Mass.
  - ▶ The Detection of the Cosmic Neutrino Background.
  - ▶ Finding Lepton Number Violation and Majorana Neutrino Masses.
  - ▶ The NuMass Workshop Series.
- ▶ Thanks to Angelo and all the local organizers for the perfect organization!
- ▶ See you at NuMass 2024!