Suppression of charged particle production at large transverse momentum in central Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV

ALICE Collaboration

Abstract

Inclusive transverse momentum spectra of primary charged particles in Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV have been measured by the ALICE Collaboration at the LHC. The data are presented for central and peripheral collisions, corresponding to 0–5% and 70–80% of the hadronic Pb–Pb cross section. The measured charged particle spectra in $|\eta| < 0.8$ and $0.3 < p_T < 20$ GeV/c are compared to the expectations in pp collisions at the same $\sqrt{s_{NN}}$, scaled by the number of underlying nucleon–nucleon collisions. The comparison is expressed in terms of the nuclear modification factor $R_{AA}$. The result indicates only weak medium effects ($R_{AA} \approx 0.7$) in peripheral collisions. In central collisions, $R_{AA}$ reaches a minimum of about 0.14 at $p_T = 6–7$ GeV/c and increases significantly at larger $p_T$. The measured suppression of high-$p_T$ particles is stronger than that observed at lower collision energies, indicating that a very dense medium is formed in central Pb–Pb collisions at the LHC.

High energy-ion collisions enable the study of strongly interacting matter under extreme conditions. At sufficiently high collision energies Quantum-Chromodynamics (QCD) predicts that hot and dense deconfined matter, commonly referred to as the Quark–Gluon Plasma (QGP), is formed. With the advent of a new generation of experiments at the CERN Large Hadron Collider (LHC) [1] a new energy domain is accessible to study the properties of this state.

Previous experiments at the Relativistic Heavy Ion Collider (RHIC) reported that hadron production at high transverse momentum ($p_T$) in central (head-on) Au–Au collisions at a centre-of-mass energy per nucleon pair $\sqrt{s_{NN}}$ of 200 GeV is suppressed by a factor 4–5 compared to expectations from an independent superposition of nuclear–nuclear (NN) collisions [2–5]. The dominant production mechanism for high-$p_T$ hadrons is the fragmentation of high-$p_T$ partons that originate in hard scatterings in the early stage of the nuclear collision. The observed suppression at RHIC is generally attributed to energy loss of the partons as they propagate through the hot and dense QCD medium [6–10].

To quantify nuclear medium effects at high $p_T$, the so-called nuclear modification factor $R_{AA}$ is used. $R_{AA}$ is defined as the ratio of the charged particle yield in Pb–Pb to that in pp, scaled by the number of binary nucleon–nucleon collisions $\langle N_{coll} \rangle$

$$R_{AA}(p_T) = \frac{(1/N_{coll}^{pp})d^2N_{AA}^{tr}/d\eta dp_T}{(N_{coll}/1/N_{coll}^{pp})d^2N_{NN}^{tr}/d\eta dp_T}.$$

where $\eta = -\ln(\tan(\theta/2)$ is the pseudo-rapidity and $\theta$ is the polar angle between the charged particle direction and the beam axis. The charged particle yields in Pb–Pb and pp are normalized to the number of underlying nucleon–nucleon collisions. The nuclear modification factor $R_{AA}$ is given by the product of $\langle N_{coll} \rangle$ and the inelastic NN cross section $\sigma_{NN}^{inel}$. If no nuclear modification is present, $R_{AA}$ is unity at high $p_T$.

At the larger LHC energy the density of the medium is expected to be higher than at RHIC, leading to a larger energy loss of high-$p_T$ partons. On the other hand, the less steeply falling spectrum at the higher energy will lead to a smaller suppression in the $p_T$ spectrum of charged particles, for a given magnitude of partonic energy loss [9,10]. Both the value of $R_{AA}$ in central collisions as well as its $p_T$ dependence may also in part be influenced by gluon shadowing and saturation effects [12–15] which in general decrease with increasing $x$ and $Q^2$.

This Letter reports the measurement of the inclusive primary charged particle transverse momentum distributions at mid-rapidity in central and peripheral Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV by the ALICE experiment [16]. Primary particles are defined as prompt particles produced in the collision, including decay products, except those from weak decays of strange particles. The data were collected in the first heavy-ion collision period at the LHC. A detailed description of the experiment can be found in [16].

For the present analysis, charged particle tracking utilizes the Inner Tracking System (ITS) and the Time Projection Chamber (TPC) [17], both of which cover the central region in the pseudorapidity range $|\eta| < 0.9$. The ITS and TPC detectors are located in...
the ALICE central barrel and operate in the 0.5 T magnetic field of a large solenoidal magnet. The TPC is a cylindrical drift detector with two readout planes on the endcaps. The active volume covers 85 < r < 247 cm and −250 < z < 250 cm in the radial and longitudinal directions, respectively. A high voltage membrane at z = 0 divides the active volume into two halves and provides the electric drift field of 400 V/cm, resulting in a maximum drift time of 94 μs.

The ITS is used for charged particle tracking and trigger purposes. It is composed of six cylindrical layers of high resolution silicon tracking detectors with radial distances to the beam line from 3.9 to 43 cm. The two innermost layers are the Silicon Pixel Detectors (SPD) with 9.8 million pixels, read out by 1200 chips. Each chip provides a fast signal if at least one of its pixels is hit. The signals from the 1200 chips are combined in a programmable logic unit which supplies a trigger signal. The SPD contributes to the minimum-bias trigger, if hits are detected on at least two chips on the outer layer. The SPD is followed by two layers of Silicon Drift Detectors (SDD) consisting of double-sided silicon micro-strip sensors, for a total of 2.6 million readout channels.

The two forward scintillator hodoscopes (VZERO-A and VZERO-C) cover the pseudo-rapidity ranges 2.8 < η < 5.1 and −3.7 < η < −1.7. The sum of the amplitudes of the signals in the VZERO scintillators is used as a measure for the event centrality. The VZERO detectors also provide a fast trigger signal if at least one particle hit was detected.

During the heavy-ion data-taking period, up to 114 bunches, each containing about 7 × 10^5 ions of 208Pb, were collided at √sNN = 2.76 TeV in the ALICE interaction region. The rate of hadronic events was about 100 Hz, corresponding to an estimated luminosity of 1.3 × 10^30 cm^-2 s^-1. The detector readout was triggered by the LHC bunch-crossing signal and a minimum-bias interaction trigger based on trigger signals from VZERO-A, VZERO-C, and SPD. The present analysis combines runs taken with two different minimum-bias conditions. In the first set of runs, two out of the three trigger signals were required, while in the second set a coincidence between VZERO-A and VZERO-C was used. Both trigger conditions have similar efficiency for hadronic interactions, but the latter suppresses a large fraction of electromagnetic reactions.

The following analysis is based on 2.3 × 10^8 minimum-bias Pb–Pb events, which passed the offline event selection. This selection is based on VZERO timing information and the correlation between TPC tracks and hits in the SPD to reject background events coming from parasitic beam interactions. Additionally, a minimal energy deposit in the Zero Degree Calorimeters (ZDC) is required to further suppress electromagnetic interactions. Only events with reconstructed vertex at |z_{c,exl}| < 10 cm were used. The definition of the event centrality is based on the sum of the amplitudes measured in the VZERO detectors as described in [18]. Alternative centrality measures utilize the cluster multiplicity in the outer layer of the SPD or the multiplicity of reconstructed tracks. The correlation between the VZERO amplitude and the uncorrected TPC track multiplicity in |y| < 0.8 is illustrated in Fig. 1. The VZERO amplitude distribution is fitted using a Glauber model [11] to determine percentage intervals of the hadronic cross section, as described in [18]. We used a Glauber model Monte Carlo simulation assuming ΔN_NN = 64 mb, a Woods–Saxon nuclear density with radius 6.62 ± 0.06 fm and surface diffuseness 0.546 ± 0.010 fm [19]. A minimum inter-nucleon distance of 0.4 ± 0.4 fm is assumed. The Glauber Monte Carlo allows one to relate the event classes to the mean numbers of participating nucleons (N_{part}) and binary collisions (N_{coll}) (see Table 1) by geometrically ordering events according to the impact parameter distribution. The errors include the experimental uncertainties in the parameters used in the Glauber simulation and an uncertainty of ±5 mb in σ_{NN}^{inel}. The TPC multiplicity distributions for the central and peripheral event samples selected for this analysis, corresponding to the 0–5% and 70–80% most central fraction of the hadronic Pb–Pb cross section, are shown in the lower panel of Fig. 1. Charged particle tracks are reconstructed in the ITS and TPC detectors. Track candidates in the TPC are selected in the pseudo-rapidity range |y| < 0.8. Track quality cuts in the TPC are based on the number of reconstructed space points (at least 70 out of a maximum of 159) and the χ^2 per space point of the momentum fit (lower than 4). The TPC track candidates are projected to the ITS and used for further analysis, if at least two matching hits in the ITS are found, including at least one in the SPD. The average number of associated hits in the ITS is 4.7 for the selected tracks. The event vertex is reconstructed by extrapolating the particle tracks to the interaction region. The event vertex reconstruction is fully efficient in both the peripheral and the central event sample. Tracks are rejected from the final sample if their distance of closest approach to the reconstructed vertex in longitudinal and radial direction, d_L and d_R, satisfies d_L > 2 cm or d_R > 0.18 cm + 0.035 cm · p_T^{-1.01}, with p_T in GeV/c.

The efficiency and purity of primary charged particles using these cuts are estimated using a Monte Carlo simulation including HIJING [20] events and a GEANT3 [21] model of the detector...
response [22]. We used a HIJING tune which reproduces approximately the measured charged particle density in central collisions [18]. In central events, the overall primary charged particle approximately the measured charged particle density in central collisions is 65% at \( p_T = 0.6 \) GeV/c and above. In peripheral events, the efficiency is larger by about 2–3%. The contamination from secondaries is 6% at \( p_T = 0.3 \) GeV/c and decreases to about 2% at \( p_T > 1 \) GeV/c, with no significant centrality dependence. This contribution was estimated using the \( d_N/dy \) distributions of data and HIJING and is consistent with a first estimate of the strangeness to charged particle ratio from the reconstruction of \( K^\circ \), \( \Lambda \) and \( \Lambda \) invariant mass peaks.

The momentum of charged particles is reconstructed from the track curvature measured in the ITS and TPC. The momentum resolution can be parametrized as \( \sigma(p_T)/p_T = a^2 + b \cdot (p_T)^2 \). It is estimated from the track residuals to the momentum fit and verified by cosmic muon events and the width of the invariant mass peaks of \( \Lambda \), \( \Lambda \) and \( K^0 \). While \( a = 0.01 \) for both centrality bins, there is a weak centrality dependence of \( b \), i.e. \( b = 0.0045 \) (GeV/c)^{-1} in peripheral events and \( b = 0.0056 \) (GeV/c)^{-1} in central events. This is related to a slight decrease for more central events of the average number of space points in the TPC. The modification of the spectra arising from the finite momentum resolution is estimated by Monte Carlo. It results in an overestimate of the yield by up to 8% at \( p_T = 20 \) GeV/c in central events. This was accounted for by introducing a \( p_T \) dependent correction factor to the \( p_T \) spectra. From the mass difference between \( \Lambda \) and \( \Lambda \) and the ratio of positive over negative charged tracks, assuming charge symmetry at high \( p_T \), the upper limit of the systematic uncertainty of the momentum scale is estimated to be \( |\Delta p_T/p_T| < 0.002 \). This has negligible effect on the measured spectra.

Table 2 shows the systematic uncertainties obtained by a comparison of different centrality measures (using the SPD instead of VZERO), and by varying the track and event quality cuts and the Monte Carlo assumptions. In particular, we studied a variation of the most abundant charged particle species (p, \( \pi \), K) by \( \pm 30\% \), the material budget by \( \pm 7\% \), and the secondary yield from strangeness decays in the Monte Carlo by \( \pm 30\% \). We have used the differences between the standard analysis and one based only on the use of TPC tracks to estimate the uncertainty on the track efficiency corrections, included in the systematic errors. The total systematic uncertainties on the corrected \( p_T \) spectra depend on \( p_T \) and are 8–10% and 5–7% for the peripheral and central event samples, respectively.

The determination of \( R_{AA} \) requires a pp reference at \( \sqrt{s} = 2.76 \) TeV, where no pp measurement exists. Different approaches are at hand which allow a prediction of the \( p_T \) spectrum at a given \( \sqrt{s} \) by scaling existing data at different energies. Such approaches assume general scaling properties of perturbative QCD (pQCD) or rely on next-to-leading order (NLO) pQCD calculations. The present analysis follows a data-driven approach with minimal theoretical assumptions where, in order to minimize systematic uncertainties, only measurements by ALICE are considered. In this approach, the pp reference spectrum is obtained by interpolating the differential yields \( d^2N_{pp}/dp_T^2 \) of charged particles measured in inelastic pp collisions at \( \sqrt{s} = 0.9 \) and 7 TeV by ALICE [23,24].

Fig. 2. The \( p_T \) distributions of primary charged particles at mid-rapidity (|\( \eta \) | < 0.8) in central (0–5%) and peripheral (70–80%) Pb–Pb collisions at \( \sqrt{s_{NN}} = 2.76 \) TeV. Error bars are statistical only. The systematic data errors are smaller than the symbols. The scaled pp references are shown as the two curves, the upper for 0–5% centrality and the lower for 70–80%. The systematic uncertainties of the pp reference spectra are contained within the thickness of the line.

Alternative approaches to derive the pp reference spectrum are investigated to study the sensitivity of \( R_{AA} \) to the specific choice of our method. Replacing in the interpolation the \( p_T \) spectrum at 0.9 TeV by the one measured in pp at \( \sqrt{s} = 1.96 \) TeV in |\( \eta \) | < 1 by the CDF Collaboration [30] results in a pp reference spectrum which is 5–15% lower than the reference spectrum described above. A different procedure to obtain a pp reference is based on a scaling of the \( p_T \) spectra at 0.9 or 7 TeV to 2.76 TeV by the relative \( \sqrt{s} \) dependence predicted by NLO pQCD calculations [31] (referred to as “NLO scaling”). Using the 7 TeV spectrum as a starting point,
good agreement with the reference obtained from interpolation is found. Starting instead from 0.9 TeV results in a spectrum which is 30–50% higher than the interpolation reference. The pp reference spectra derived from the use of the CDF data in the interpolation and from NLO scaling of the 0.9 TeV data are used in the following to illustrate the dependence of $R_{AA}$ at high $p_T$ on the choice of the reference spectrum.

The $p_T$ distributions of primary charged particles in central and peripheral Pb–Pb collisions at 2.76 TeV are shown in Fig. 2, together with the binary-scaled yields from pp collisions. The $p_T$-dependence is similar for the pp reference and for peripheral Pb–Pb collisions, exhibiting a power law behaviour at $p_T > 3$ GeV/c, which is characteristic of perturbative parton scattering and vacuum fragmentation. In contrast, the spectral shape in central collisions clearly deviates from the scaled pp reference and is closer to an exponential in the $p_T$ range below 5 GeV/c.

Fig. 3 shows the nuclear modification factor $R_{AA}$ for central and peripheral Pb–Pb collisions. The nuclear modification factor deviates from one in both samples. At high $p_T$, where production from hard processes is expected to dominate, there is a marked difference between peripheral and central events. In peripheral collisions, the nuclear modification factor reaches about 0.7 and shows no pronounced $p_T$ dependence for $p_T > 2$ GeV/c. In central collisions, $R_{AA}$ is again significantly different from one, reaching a minimum of $R_{AA} \approx 0.14$ at $p_T = 6–7$ GeV/c. In the intermediate region there is a strong dependence on $p_T$ with a maximum at $p_T = 2$ GeV/c. This may reflect a variation of the particle composition in heavy-ion collisions with respect to pp, as observed at RHIC [32,33]. A significant rise of $R_{AA}$ by about a factor of two is observed for $7 < p_T < 20$ GeV/c. Shown as histograms in Fig. 3, for central events only, are the results for $R_{AA}$ at high $p_T$, using alternative procedures for the computation of the pp reference, as described above. For such scenarios, the overall value for $R_{AA}$ is shifted, but a significant increase of $R_{AA}$ in central collisions for $p_T > 7$ GeV/c persists.

In Fig. 4 the ALICE result in central Pb–Pb collisions at the LHC is compared to measurements of $R_{AA}$ of charged hadrons ($\sqrt{s_{NN}} = 200$ GeV) by the PHENIX and STAR experiments [34, 35] at RHIC. At 1 GeV/c the measured value of $R_{AA}$ is similar to those from RHIC. The position and shape of the maximum at $p_T \sim 2$ GeV/c and the subsequent decrease are similar at RHIC and LHC, contrary to expectations from a recombination model [36]. Despite the much flatter $p_T$ spectrum in pp at the LHC, the nuclear modification factor at $p_T = 6–7$ GeV/c is smaller than at RHIC. This suggests an enhanced energy loss at LHC and therefore a denser medium. A quantitative determination of the energy loss and medium density will require further investigation of gluon shadowing and saturation in the present energy range and detailed theoretical modeling.

In summary, we have measured the primary charged particle $p_T$ spectra and nuclear modification factors $R_{AA}$ in central (0–5%) and peripheral (70–80%) Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV with the ALICE experiment. The nuclear modification factor in peripheral collisions is large and independent of $p_T$ for $p_T > 2$ GeV/c, indicating only weak parton energy loss. For central collisions, the value for $R_{AA}$ is found to be $\approx 0.14$ at $p_T = 6–7$ GeV/c, which is smaller than at lower energies, despite the much less steeply falling $p_T$ spectrum at the LHC. Above 7 GeV/c, $R_{AA}$ increases significantly. The observed suppression of high-$p_T$ particles provides evidence for strong parton energy loss and large medium density at the LHC.
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ALICE Collaboration

K. Aamodt a, A. Abrahantes Quintana b, D. Adamov c, A.M. Adare d, M.M. Aggarwal e, G. Aglieri Rinella f, A.G. Agocs g, S. Aguilar Salazar h, Z. Ahamed i, N. Ahmad j, A. Ahmad Masoodi j, S.U. Ahr k l, A. Akhendinov l, D. Aleksandrov m, B. Alessandro n, R. Alfaro Molina h, A. Alici o, A. Alkin p, E. Almaráv Aviña q, T. Alt v, V. Altini r, S. Altipinara e, I. Altsybeev r, C. Andrei u, A. Andronic s, V. Anguslov v, C. Anson w, T. Anticic x, F. Antinori y, P. Antonioliz, L. Aphelache t, H. Appelschäuser, N. Arbor, S. Arcelli, A. Aren d, N. Armosto a, R. Arnaldi a, T. Aronsson d, I.C. Arsene a, A. Asryan t, A. Augustins t, R. Averbeck t, T.C. Awes ae, J. Åystö af, M.D. Azmi, M. Bach q, A. Badalà à, Y.W. Baek k l, S. Bagnasco n, R. Bailhache ab, R. Bala ab, R. Baldini Ferroli a, A. Baldissieriai,
9 Now at Centro de Investigación y de Estudios Avanzados (CINVESTAV), Mexico City and Mérida, Mexico.
10 Now at Benemérita Universidad Autónoma de Puebla, Puebla, Mexico.
11 Now at Laboratoire de Physique Subatomique et de Cosmologie (LPSC), Université Joseph Fourier, CNRS–IN2P3, Institut Polytechnique de Grenoble, Grenoble, France.
12 Now at Institut Pluridisciplinaire Hubert Curien (IPHC), Université de Strasbourg, CNRS–IN2P3, Strasbourg, France.
13 Now at Sezione INFN, Padova, Italy.
14 Deceased.
15 Also at Division of Experimental High Energy Physics, University of Lund, Lund, Sweden.
16 Also at University of Technology and Austrian Academy of Sciences, Vienna, Austria.
17 Also at European Organization for Nuclear Research (CERN), Geneva, Switzerland.
18 Now at Oak Ridge National Laboratory, Oak Ridge, TN, United States.
19 Also at European Organization for Nuclear Research (CERN), Geneva, Switzerland.
20 Also at Wayne State University, Detroit, MI, United States.
21 Also at Frankfurt Institute for Advanced Studies, Johann Wolfgang Goethe-Universität Frankfurt, Frankfurt, Germany.
22 Now at Research Division and ExtreMe Matter Institute EMMI, GSI Helmholtzzentrum für Schwerionenforschung, Darmstadt, Germany.
23 Also at Fachhochschule Köln, Köln, Germany.
24 Also at Institute of Experimental Physics, Slovak Academy of Sciences, Košice, Slovakia.
25 Now at Instituto de Ciencias Nucleares, Universidad Nacional Autónoma de México, Mexico City, Mexico.
26 Also at Laboratoire de Physique Subatomique et de Cosmologie (LPSC), Université Joseph Fourier, CNRS–IN2P3, Institut Polytechnique de Grenoble, Grenoble, France.
27 Also at “Vinča” Institute of Nuclear Sciences, Belgrade, Serbia.
28 Also at University of Houston, Houston, TX, United States.
29 Also at Department of Physics, University of Oslo, Oslo, Norway.
30 Also at Variable Energy Cyclotron Centre, Kolkata, India.
31 Now at Department of Physics, University of Oslo, Oslo, Norway.
32 Also at Dipartimento Interateneo di Fisica ‘M. Merlin’ and Sezione INFN, Bari, Italy.
33 Now at Nikhef, National Institute for Subatomic Physics and Institute for Subatomic Physics of Utrecht University, Utrecht, Netherlands.
34 Also at Hua-Zhong Normal University, Wuhan, China.
35 Also at Centro Fermi – Centro Studi e Ricerche e Museo Storico della Fisica “Enrico Fermi”, Rome, Italy.